

Special Issue Reprint

Photonics for Bioapplications

Sensors and Technology

Edited by Nélia Jordão Alberto, Maria de Fátima Domingues, Nunzio Cennamo and Adriana Borriello

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Guest Editors

Nélia Jordão Alberto Maria de Fátima Domingues Nunzio Cennamo Adriana Borriello



Guest Editors

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About the Editors

Nélia Jordão Alberto

Nélia Jordão Alberto completed her PhD in Physics in 2011, from the University of Aveiro (Portugal). Currently, she is an Assistant Researcher at the Instituto de Telecomunicações (Portugal), and a lecturer at the Physics Department of the Aveiro University. She is the author/co-author of more than 120 peer-reviewed papers in journals and conferences, several book chapters, and co-editor of two books. Dr. Alberto was listed in the top 2% of the most-cited scientists according to Stanford University and Elsevier rankings in the years 2020, 2022, 2023 and 2024. Her research interests include the development of photonic-based solutions and their application, with a special focus on the biomedical context. Dr. Alberto has participated as technical program committee member of distinct international events (AOP2024, Globecom2024 SAC-EH, Globecom2025 SAC-EH, IEEE Healthcom 2024 and 2025, IEEE ICC'25 and ICC'26, . . .).

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Maria de Fátima Domingues is an Assistant Professor at the Biomedical Engineering and Biotechnology Department at Khalifa University, Abu Dhabi, UAE and a Research Collaborator at Instituto de Telecomunicações, Aveiro. Maria de Fátima Domingues received her PhD in Physics Engineering from University of Aveiro, Portugal, in December 2014, and in 2015 started a joint Research position at the Instituto de Telecomunicações – Aveiro and the Consejo Superior de Investigaciones Científicas (CSIC)-Madrid, Spain. She is the lead researcher in biomedical applied photonics at the BMEBD-KU, with +130 co-authored peer-reviewed publications in this field. In 2024 Dr. Domingues was listed in the top 2% of the most-cited scientists according to Stanford University and Elsevier rankings. She has been actively cooperating in project proposals both as a PI and co-PI, at national (UAE), and international (EU) levels. She is a project reviewer for entities such as the European Research Council (ERC) and Marie Skłodowska Curie Actions HORIZON MSCA DN. She has been part of the organizing committee of various international conferences and scientific meetings, including IEEE HealthCom 2022 and Healthcom 2025, IEEE WF-IoT2024, 2025 and 2026, IEEE ICC 2024 and Globecom 2025. She is an IEEE Senior Member and an active member of several IEEE Societies. Since January 2022, she is the Secretary for the IEEE EMBS Portuguese Chapter.

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Nunzio Cennamo (NC) was born in Italy in 1975. He received a master's degree in Electronic Engineering in 2002 and the PhD degree in Electronic Engineering in 2005, both from the Second University of Naples, Naples, Italy. He is a full Professor in Electronics at the University of Campania Luigi Vanvitelli, Naples (Italy). His research interests include the design and fabrication of optical sensors, chemical sensors, biosensors and optoelectronic devices. He is the author of more than 230 international journal and conference papers, and 15 patents. He is cofounder of the Spin Off "MORESENSE srl" in Milan (Fondazione Filarete, Milan, Italy). He is an Associate Editor of "Photonics Research" (OSA), of "Applied Sciences" (MDPI), and "Sensors" (MDPI). He is an Organizer and General Chair of several International Conferences, the IEEE "Sensor Applications Symposium" (IEEE SAS 2024), Naples (Italy) July 23-25, "7th International Symposium on Sensor Science" (I3S 2019) 9-11 May 2019, Naples (Italy), and of several editions of the "International Electronic Conference on Applied Sciences" (ASEC 2021, ASEC 2022, ASEC 2023, ASEC 2024, ASEC 2025). He is a Publicity Chair of the last two editions (the 6th and 7th) of the "IEEE International Symposium M&N" (18-20

July 2022 Padua - Italy and 3-5 July 2024 Rome - Italy), and from the 2025 Edition, a member of the IEEE SAS Steering Committee. He has been an invited speaker several times in Webinars, Summer Schools, and International Conferences. He is an Organizer/Chair of several Special Sessions in International Conferences. He is a Principal Investigator of several national projects and Coordinator of the Local Unit of national and international projects (total budget major than 2 M Euros).

Adriana Borriello

Adriana Borriello (AB) obtained her PhD in Cellular Biochemistry in 1999 (Second University of Naples). Since 2015, AB is Associate Professor of Biochemistry at the University of Campania "L. Vanvitelli" and since 2007, Biology Manager at AOU Vanvitelli. She has obtained the National Scientific Qualification for the role of Full Professor, SSD General Biochemistry (2017). AB is Academic Board Member of the "Biochemical and Biotechnological Sciences" PhD Program since 2008 and has intense teaching activities in Medicine and Surgery and Degree Courses for health professionals (Vanvitelli). She is senior member of several Scientific Societies and Associate Editor and Editorial Board Member for international journals, such as Frontiers in Endocrinology, Cells (MDPI), and Diseases (MDPI), as well as Guest Editor for special issues in Cells and Biosensors (MDPI). AB has been PI/participant in numerous national research projects. Her research interests are related to cell division control, post-translational regulation of the Cip/Kip inhibitors, and characterization of their new roles in controlling cytoskeleton dynamics. Interests also concern genetic/biochemical mechanisms of hematological diseases, including congenital polycythemias due to altered hypoxia response. Recent efforts are devoted to develop, through multidisciplinary collaborative studies, surface plasmon resonance-based point-of-care biosensors for analytes of interest in biomedical applications. She has been invited speaker/Committee member in numerous international conferences and author of more than 80 peer-reviewed papers in international, high-impact journals, including the New England Journal of Medicine, Leukemia, Cancer Research, and Free Radicals in Biology and Medicine, and several book chapters, with >3,000 total citations and an H-index of 33 (Scopus).





Editorial

Photonics for Bioapplications: Sensors and Technology

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Over the past decade, interest in advancing photonic systems for bioapplications has been steadily growing, and various key factors have driven this trend. One major factor is the rapid progress in photonic technology—effects like miniaturization, increased sensitivity, and improved resolution—which has broadened how these systems can be used.

Another driving force is the rising demand for more precise diagnostics. As the need for non-invasive, real-time, and highly sensitive diagnostic tools increases, photonic systems have become crucial, especially in areas like biosensors and point-of-care devices. Furthermore, new applications have emerged; photonic technologies are now being used in fields such as environmental monitoring, biochemical analysis, biomedical imaging, and personalized medicine, further expanding their reach and significance.

Additionally, the development of new materials and bio-receptors, smart sensing biosurfaces, innovative sensor designs, and related sensing instrumentation have significantly enhanced the performance of photonics systems. These advancements have made the technology more suitable and adaptable to complex bioapplications. The interdisciplinary nature of this field has also contributed to its development. By merging photonics with biology, chemistry, and medicine, researchers from different scientific backgrounds are finding new opportunities to collaborate and innovate. Overall, this increasing interest in photonic systems reflects their potential to address major challenges in healthcare, environmental science, and biotechnology.

This Special Issue was designed to showcase the latest innovations in photonic sensing and interrogation systems for bioapplications. It addressed recent technological advancements and novel materials to smart bio-surfaces, emerging applications, innovative sensor designs, and cutting-edge instrumentation. Contributions included original research articles, communications, and review papers, aiming to provide a snapshot of the current state of research and the possible future directions for the field. The goal was to create a high-quality collection of papers on the following, or related, key topics:

- Optical fiber biosensing/immunosensors;
- New bio/chemical probes for bioapplications;
- Wearable biomedical sensors;
- Point-of-care devices;
- Molecular diagnosis;
- Biomarker detection;

- Imaging sensors;
- Optical sensors in e-Health architectures;
- Optical chemical sensors;
- Energy-efficient eHealth architecture;
- Big data analysis for eHealth;
- Plasmonic sensors and interrogation systems;
- Molecularly imprinted polymers for sensing;
- Advanced signal processing techniques;
- Photonic integrated circuits (PICs) for bioapplications;
- Low-cost, miniaturized, selective, and multiparameter photonic devices.

The contributions to this Special Issue resulted in a collection of 11 manuscripts that report advances in photonics for bioapplications, emphasizing the relevance of the topic to the research community and showcasing the diverse applications where these innovative technological developments can enhance both performance and impact.

For instance, Yin et al. developed an enzyme-assisted fluorescent biosensor that utilizes circular single-stranded DNA for miRNA detection (Contribution 1). This device reaches a detection threshold as low as 0.36 nM, outperforming many previous miRNA sensors and demonstrating robust detection even in complex samples like serum, making it a valuable tool for early disease diagnosis.

Trypsin is recognized as a potential biomarker in cancers like colorectal, gastric, and pancreatic. A novel near-infrared sensitive dye–peptide conjugate (SQ-3 PC) for selectively detecting trypsin activity via fluorescence ON/OFF sensing is proposed (Contribution 2). The probe exhibits high sensitivity, a linear response up to 30 nM, and a limit of detection and limit of quantification of 1.07 nM and 3.25 nM, respectively. The device demonstrates strong selectivity for trypsin, even in the presence of other enzymes, making it suitable for complex biological applications.

The communication proposed by Xu et al. reports a new THz metasurface sensor for cinnamoylglycine detection, a urinary biomarker of gestational diabetes mellitus (GDM) (Contribution 3). The device, produced from lithium tantalate prism tetramers on quartz, achieves a high Q-factor of 231, enhancing molecular detection. Simulations show a detection limit of 1.23 μ g \times cm⁻², demonstrating excellent sensitivity. Given the risks associated with GDM, like preeclampsia and preterm birth, the sensor offers a faster and more accurate alternative to traditional tests, with potential for detecting other biomarkers and monitoring disease progression.

Park et al. developed an *Asterias forbesi*-inspired surface-enhanced Raman spectroscopy (SERS) with numerous hotspots (Contribution 4). The substrate showed strong SERS performance, minimized oxidation damage, and ensured repeatability for uric acid detection. Under laboratory conditions, uric acid was detected at concentrations as low as 1.16 nM, and selectivity was demonstrated against various metabolites. In a simulated real-world scenario, the sensor successfully detected a wide range of uric acid concentrations, encompassing both deficiency and excess levels, in serum and urine samples.

In the study presented by Jiao et al., the polarimetric aberrations caused by three different illumination schemes in backscattering Mueller Matrix imaging were analyzed (Contribution 5). The optimized schemes established key criteria for selecting spatial illumination patterns in non-collinear backscattering Mueller Matrix measurements, offering valuable insights for advancing quantitative tissue polarimetric imaging and biosensing.

In (Contribution 6), the viscoelastic properties of AuNP-assisted radiofrequency-heated and ablated tissues were analyzed, correlating them with temperature profiles using confocal Brillouin micro-spectroscopy and Mg-silicate-NP-doped temperature-sensing fibers. The results reveal that changes in the Brillouin peak were linked to thermal doses and

protein denaturation. The hybrid Brillouin–optical backscattering reflectometry technique shows potential for monitoring tissue deformations during thermal therapies, enabling the real-time tracking of viscoelastic and thermal properties in metal-nanoparticle-embedded tissues for hyperthermal and theranostic applications.

A red-light-driven photoelectrochemical (PEC) sensor based on a g- C_3N_4 /CuS/TiO₂ ternary heterojunction for the sensitive detection of the synergistic antioxidant effects of sesamol and other antioxidants is reported in (Contribution 7). Sesamol is a natural antioxidant that reduces free radicals and supports brain function, improving memory and cognition in Alzheimer's patients. By scavenging holes in the valence band, the antioxidants suppress charge carrier recombination, boosting the PEC photocurrent. The sensor exhibited high sensitivity, selectivity, and stability, when applied to detect sesamol in soybean and peanut oils. It holds potential for applications in nutrition analysis, food fraud detection, and medical research, contributing to improved consumer health and product quality.

Zhou et al. developed titanium metal–organic framework (Ti-MOF) microflowers for the detection and removal of Al(III) ions in water (Contribution 8). Ti-MOF offers high water stability, porosity, and strong luminescence, with a detection range of 0.4–15 μ M and a limit of 75 nM, well below the World Health Organization's aluminum safety limit. It also allows selective Al(III) detection by shielding interference from Fe(III) using ascorbic acid, making it an effective sensor and adsorbent for water purification.

A nano-slit array sensor based on temperature variation is proposed for the highly sensitive and specific detection of molecular fingerprints (Contribution 9). As the temperature changes, the dielectric properties of the temperature-sensitive semiconductor material InSb dynamically adjust, causing shifts in the transmission resonance angle and producing transmission envelope curves over a broad frequency range. Using this temperature scanning approach, the characteristic fingerprint spectra of hexogen, a high-impact military explosive, were successfully identified, allowing both qualitative and quantitative detection with a sensitivity limit of $1.61~\mu g/cm^2$.

A distributed optical fiber sensing network was proposed in (Contribution 10) for the two-dimensional in situ thermal mapping of advanced methods for radiofrequency thermal ablation. Radiofrequency ablation was conducted on a bovine phantom using a standard approach, while methods enhanced with agarose and gold nanoparticles were carried out to improve the ablation efficiency. Results showed that agarose-mediated thermal ablation treated the largest area, but with low repeatability, while AuNP-mediated ablation at a 4 mg/mL density offered the best balance between ablation efficacy and repeatability.

The Special Issue concludes with a review paper dedicated to cost-effective optical fiber-based biosensing platforms, exploring various geometries, interrogation techniques, and encapsulation methods (Contribution 11). The potential of these cost-effective optical fiber sensing techniques is highlighted through numerous successful applications, in three key areas: cancer, cardiovascular biomarker detection, and environmental monitoring. Some weaknesses of these systems are also highlighted, emphasizing the need to improve aspects related to repeatability and reproducibility, as well as the development of techniques for sensor reuse.

In conclusion, this Special Issue presents a selection of manuscripts that highlight recent advances in the field of photonics, with a particular focus on bioapplications. The works gathered here reflect the diversity and breadth that characterize both the research domain itself and the applications. Although they represent only a sample of the extensive ongoing developments, the published articles provide a meaningful overview of the current trends and pave the way for future opportunities for research and innovation in the

field of photonic sensors, emphasizing their potential contributions to biomedicine and related areas.

We sincerely thank the authors for their valuable contributions to this Special Issue, which brings together original research articles that provide a comprehensive and current overview of a diverse array of topics, all unified by a shared focus on "Photonics for bioapplications: Sensors and Technology". Additionally, we express our gratitude to the academic editors and reviewers for their dedicated efforts in ensuring the scientific rigor and quality of the published manuscripts. Their efforts have significantly enhanced the overall impact and relevance of this Special Issue.

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List of Contributions:

- Yin, X.; Liao, Y.; Li, F.; Li, J.; Du, J. Enzyme-assisted fluorescence biosensor based on circular single-stranded DNA without group modification for microRNA detection. *Biosensors* 2024, 14, 527.
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Article

Enzyme-Assisted Fluorescence Biosensor Based on Circular Single-Stranded DNA Without Group Modification for MicroRNA Detection

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Abstract: Fluorescent biosensor, which has the characteristics of high sensitivity, specificity, and low cost, can be directly detected in physiological fluids such as blood and serum. Therefore, the development of fluorescence sensor platforms for miRNA detection has a positive effect on the prevention and treatment of various diseases. In this paper, miR-34a was selected as a biological indicator of Alzheimer's disease (AD). We designed a circular single-stranded DNA (CSSD) biosensor, which uses two unmodified single-stranded DNA (ssDNA) with complementary ends, DNAa and DNAb, to form CSSD by DNA sequence pairing to improve thermal stability and achieve signal amplification. At the same time, CSSD can react with miR-34a, and then the DNA of the DNA-RNA chain is hydrolyzed by duplex-specific nuclease (DSN enzyme). Finally, miR-34a is released to partake in the subsequent step, thus realizing cycle amplification. By evaluating the change in fluorescence signal under the optimized conditions, we discovered that this approach exhibits impressive sensitivity, with a detection threshold reaching as low as 0.36 nM. This surpasses the performance of numerous preceding miRNA detection biosensors. Furthermore, the system displays excellent detection capabilities even in intricate settings like serum, showcasing a strong ability to differentiate and choose effectively. In summary, this is a signal-off fluorescent biosensor, which realizes the purpose of double amplification of biosensor signal by using CSSD and enzyme assistance so that it can be used as a valuable tool for early diagnosis of diseases.

Keywords: duplex-specific nuclease; fluorescence reduction; GelRed; miR-34a

1. Introduction

Alzheimer's disease (AD) is a secluded, age-related neurodegenerative disease, which is the most common form of dementia [1]. The elderly are the people mainly affected by AD, and the incidence of AD increases with the increase in life expectancy [2]. The incidence of AD is expected to double every five years after the age of 65 [3], and there is currently no cure [4]. In the absence of effective treatment, early diagnosis plays a vital role in improving the subsequent condition, and the diagnosis of Alzheimer's disease is expensive, so it is imperative to create a novel and cost-effective diagnostic approach [5]. MicroRNA (miRNA) are a length of 19 to 25 nucleotides of non-coding RNA [6], and they are involved in many critical biological processes. As a result of miRNA in the important role of gene expression [7], their disorders are associated with the pathology of many diseases [8], such as liver, skin, skeletal muscle, cardiovascular, immune, neurological and inflammatory system diseases, as well as cancer [9]. This also means that the miRNA biomarkers can be used in various diseases [10]. In the past, miR-34a was regarded as a biomarker for malignant tumors and cancers [11-13]. However, recent studies have shown that miR-34a is related to brain neurophysiology and pathology [14], especially Alzheimer's disease. Therefore, in addition to being a cancer marker [13], miR-34a is also expected to

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be a biomarker for AD [5,14]. Due to the extremely low abundance of biomarkers in cells, the development of sensitive platforms for the determination of biomarkers has broad research prospects [15]. Biosensors, including colorimetric, Raman scattering spectroscopy, electrochemical and fluorescent biosensors, etc., are characterized by high sensitivity, high specificity and high cost effectiveness, and can be directly detected in physiological fluids (blood, serum, etc.) [16]. Many researchers choose to use fluorescence detection biosensors to detect biomarkers [17], and their methods are widely used in medicine, food analysis, environmental and agricultural analysis, and industry [18-20]. DNA fluorescence biosensors transform the chemical energy of the tested biomaterial system into light energy. This categorization divides DNA fluorescence sensors into two groups: those with labeling and those without. Wang et al. designed a unique DNA detection method that involves the assembly of two functional nucleic acids to form a new nucleic acid structure upon exposure to the target DNA [21]. The nucleic acid structure represents a functional catalytic nucleic acid that can undergo substrate cleavage reactions in the presence of Mg²⁺. This leads to the cleavage of DNA labeled with fluorophore and cataplexy groups, resulting in a transition from low to high fluorescence intensity. Consequently, this method enables the quantitative detection of target DNA, with a sensor sensitivity of 0.01 pM and excellent selectivity.

Double-stranded specific nuclease (DSN) selectively degrades DNA in double-stranded DNA and DNA–RNA hybrids, has little activity against single-stranded nucleic acid molecules and double-stranded RNA, and can distinguish single-base mismatches well. The characteristics of DSN enzymes make them a promising option for signal amplification in DSN-based biosensors [22]. Liao et al. devised a nanochannel biosensor incorporating a DSN enzyme signal amplification technique [23]. In this system, the capture probe DNA is firmly attached to the surface of the nanochannel. Upon interaction with the target miRNA, the nanochannel undergoes hybridization, resulting in the formation of the DNA–RNA duplex. The DSN enzyme cleaved the probe DNA in the DNA–RNA double-stranded structure to extract the target miRNA. The target miRNA can then bind to other DNA probes, leading to signal amplification. GR serves as a secure and reliable intercalating dye, attaching to nucleic acid through electrostatic and charge interactions. It can directly replace the highly toxic ethidium bromide (EB), and its sensitivity is much higher than that of EB.

In this experiment, we designed an enzyme-assisted fluorescence biosensor without group modification. Based on the Watson–Crick base complementary pairing principle, the design had two single-stranded DNA (ssDNA) strands so that their terminals are complementary sequences, and there were two miR-34a binding sites on each DNA strand. The two ssDNA were hybridized to form circular single-stranded DNA (CSSD). GR intercalation dye was then added, which could select the DNA specifically and was fluorescent. The stronger the thermal stability of DNA, the stronger the fluorescence signal [24]. When the DSN enzyme is activated, it can cause the DNA strand in double-stranded DNA-RNA to be hydrolyzed, leading to the loss of the fluorescence signal if there is a specific miRNA that binds to CSSD. The target RNA chain and GR were released and they entered the next cycle. If no target miRNA was present, the fluorescence signal was always present. The sensitivity and specificity of the sensor were measured. This kind of biosensor provides a new idea for RNA detection.

2. Materials and Methods

2.1. Materials

The chemicals and reagents employed in the experiments are of analytical reagent grade and do not need any purification of the received pharmaceutical products. The DSN buffer was made up of a solution containing 50 mM Tris-HCl (pH 8.0), while the DSN master buffer was a mixture of 500 mM Tris-HCl (pH 8.0) and 50 mM MgCl $_2$. DNA annealing buffer (5×, sterile solution, 0.1 M Tris-HCl, 0.1 M NaCl, 0.05 M EDTA) and PCR amplification buffer (10×, pH 8.3, 100 mM Tris-HCl buffer, 15 mM Mg $^{2+}$, 500 mM KCl.) were purchased from Shanghai Yuanye Biotechnology Co., Ltd. (Shanghai, China); GR

nucleic acid dye (10,000×) aqueous solution was purchased from Sangong Bioengineering Co., Ltd. (Shanghai, China); Tris-HCI buffer (pH 8.0) was purchased from Shanghai MacLean Biochemical Technology Co., Ltd. (Shanghai, China); DEPC-treated water was purchased from Beijing Solarbio Science & Technology Co., Ltd. (Beijing, China); glycerol solution (100%, sterile) was purchased from Beijing Ranjeco Technology Co., Ltd. (Beijing, China); DSN enzyme was purchased from Shenzhen Newbond Biotechnology Co., Ltd. (Shenzhen, China); and Fetal bovine serum (FBS) was purchased from Zhejiang Tianhang Biotechnology Co., Ltd. (Deqing, China).

The DNA utilized in this experiment, as shown in Table 1, was supplied by Shanghai Sangong Bioengineering Co., Ltd. (Shanghai, China). The lengths of DNAa and DNAb are 66 nt, and they contain two fragments complementary to miR-34a with a length of 22 nt. In addition, they have a complementary 9 nt fragment at each end. The resulting CSSD has a double chain length of 18 nt and a single chain length of 48 nt.

Table 1. Oligonucleotide sequen	ces used in this experiment.
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Name	Sequence (from 5' to 3')		
miR-34a	5'-UGG CAG UGU CUU AGC UGG UUG U-3'		
DNAa	5'-GCTGACTGCACAACCAGCTAAGACACTGCCATT		
DNAa	TTACAACCAGCTAAGACACTGCCACGATCTCAC-3'		
DNIAL	5'-GCTAGAGTGACAACCAGCTAAGACACTGCCATT		
DNAb	TTACAACCAGCTAAGACACTGCCACGACTGACG-3'		
miR-34a	5'-UAGCUUAUCAGACUGAUGUUGA-3'		
miR-192	5'-CUG ACC UAUGAA UUG ACA GCC-3'		
miR-34a	5'-UGG AGU GUG ACA AUG GUG UUG-3'		
miR-15a	5'-UAG CAG CAC AUA AUG GUU UGU G-3'		
miR-660	5'-UAC CCA UUG CAU AUC GGA GGU GU G-3'		
miR-155	5'-UUA AUG CUA AUC GUG AUA GGG GU-3'		

2.2. Apparatus

Fluorescence signal values were detected using the Shimadzu RF-6000 fluorescence spectrophotometer (Shimadzu, Kyoto, Japan). The experiment uses GR, an intercalated fluorescent dye that selectively binds DNA and RNA strands to produce a fluorescent signal. GR is attached to nucleic acids through electrostatic and charge interactions, and the fluorescence intensity of free GR is negligible. In the presence of ssDNA, GR was embedded into ssDNA to form GelRed-ssDNA complex with significant fluorescence intensity. Strong electrostatic interactions could stabilize the GelRed-ssDNA complex.

2.3. Preparation of Reagents

DNA was diluted to 10 μ mol/L and miRNA was diluted to 20 μ mol/L with DEPC treated water (enzyme free, sterile). Our choice of buffer for the experiment was the DNA annealing buffer (1 \times), stored in a refrigerator at 4 °C. The DSN buffer and the DSN master buffer both were stored at a temperature of 4 °C. Lyophilized DSN enzyme was added to DSN buffer at 5 μ L per 10 U and an equal volume of 100% glycerol was added. The DSN enzyme, set at 1 U/ μ L, was stored at -20 °C for utilization.

2.4. Generation of CSSD

According to the previous work of Meng et al. [20], the ends of two ssDNA (DNAa, DNAb) are complementary sequences. First, 1 μM of DNAa and DNAb were added to the annealing buffer with 5 μL each, then heated at 95 °C for 2 min, and then cooled at room temperature for incubation. Finally, circular DNA was formed, and its concentration was 2.5 $\mu mol/L$.

2.5. Detection of Target miR-34a

According to the previous work of Sapia et al. [24,25], the experimental steps designed by us are as follows. Add 7.5 μ L GR dye (3×) nucleic acid solution to CSSD solution and incubate at 25 °C for 15 min. After adding 2 μ L of deionized water for 5 min at 95 °C, the mixture is cooled for 30 min at room temperature. Then, 0.5 μ L DSN enzyme (1 U/ μ L) is added, and the DSN enzyme is inactivated at 45 °C for 1 h and 75 °C for 5 min. The measured fluorescence value is represented by F₀. GR solution was added to CSSD solution and incubated at 25 °C for 15 min. Two μ L miR-34a are added, incubated at 95 °C for 5 min, cooled for 30 min, and then the DSN enzyme is added. Then, the fluorescence value F is detected, and finally the fluorescence comparison result (F₀-F) is obtained.

2.6. Testing miR-34a

The fluorescence signal value was measured by a fluorescence spectrophotometer with model RF-6000. The mixture used for the test consisted of 2 mL of ultra-pure water mixed with the experimental solution. The excitation wavelength used in this mixture was 270 nm, while the emission wavelength was 545 nm, and the target scan range was between 510 nm and 580 nm. The excitation frequency band was 5 nm and the emission frequency band was 10 nm.

3. Results and Discussion

3.1. Principle of miRNA Detection

The experimental design of the enzyme-assisted fluorescent biosensor detection principle of miR-34a is shown in Figure 1. In this study, we first designed two DNA single strands with complementary ends, DNAa and DNAb, which contain two complementary sequences of miR-34a, respectively. Therefore, hybridization of DNAa and DNAb produces CSSD with four miRNA binding sites. Since the fluorescence signal intensity of GR is enhanced with the increasing thermal stability of DNA, the fluorescence signal intensity can be improved after the two DNA single strands are combined to form a ring, and thus the fluorescence value F_0 is obtained. Starting from the presence or absence of target detection miRNA, two routes can be obtained. Route one is that when the detection target exists, miR-34a carries out the hybridization reaction with circular DNA. At this time, under the action of the DSN enzyme, the DNA strand in the DNA-RNA hybrid strand is selectively hydrolyzed, miR-34a and GR dye are released, and the fluorescence signal value decreases. The released miR-34a can react with other unreacted circular DNA to achieve cyclic amplification. After the above reaction, the fluorescence value F was obtained after high temperature treatment and the DSN enzyme was inactivated. In route two, no RNA hybridizes with DNA due to the deletion of miR-34a. Even in the presence of DSN enzymes, single-stranded DNA cannot be hydrolyzed, so the fluorescence signal value remains unchanged. Finally, the fluorescence difference F₀-F can be obtained by changing the fluorescence value, thereby determining the presence of miR-34a to be tested. Because the fluorescence values of the two lines are based on the GR dye, F₀ and F are essentially the same.

3.2. Feasibility Study of the miR-34a Assay

According to the excitation spectrum and emission spectrum of the sample to be tested, as shown in Figure 2A, we determined that the excitation wavelength for detection was 270 nm and the emission wavelength was 510–580 nm. In this experiment, in the presence of target miR-34a, two ssDNAs with complementary ends were combined to form CSSD. Oligo is one of the three most popular sequence design software in the world. It has powerful features, and its results are widely recognized. The structure of our designed single-stranded oligonucleotide was manually input into Oligo 7, and the Tm value was calculated using the nearest neighbor method. The Tm of DNAa and DNAb were measured at 77.1 $^{\circ}$ C and 76.1 $^{\circ}$ C, respectively. The Tm of PCR product CSSD was 96.9 $^{\circ}$ C, and the Tm of CSSD/RNA complex was 101.9 $^{\circ}$ C. The circular DNA strands were subjected

to GR, a nucleic acid dye, to generate a fluorescent signal. Because GR fluorescence intensity is related to DNA thermostability, CSSD fluorescence intensity was higher than ssDNA fluorescence intensity. After that, under the action of the DSN enzyme, which only hydrolyzes the DNA strand in the DNA-RNA hybrid strand, the CSSD hybridized with the target strand miR-34a was hydrolyzed. The fluorescence signal generated by the binding of the DNA strand to the dye also disappeared. To confirm the viability of the experiment, we introduced the miR-34a target substance.

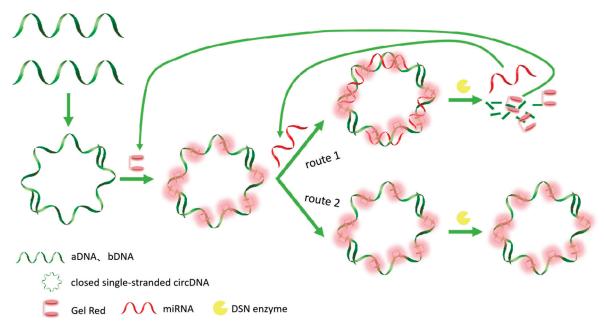


Figure 1. The principle of a biosensor based on CSSD amplification of enzyme-assisted fluorescent signals for miRNA detection.

First, to verify whether the CSSD produced by ssDNA hybridization could really improve fluorescence intensity, three sets of comparison experiments were conducted. The experimental findings depicted in Figure 2B were obtained under the same conditions as those of our previous tests. The blue curve represents the fluorescence intensity with the addition of 5 μL DNAa and 5 μL DNAb, the brown curve represents the fluorescence intensity with the addition of only 10 μL DNAa, and the yellow curve represents the fluorescence intensity with the addition of only 10 μL DNAb. After the base complementary pairing, the fluorescence intensity of the addition of two ssDNAs was significantly greater than that of the addition of only one DNA, indicating that the fluorescence intensity was indeed improved. In addition, when the DNA concentration was the same, the fluorescence signal was significantly enhanced, which also indicated the formation of CSSD. The difference in fluorescence intensity of DNAa & DNAb, DNAa, and DNAb after enzyme digestion reaction is shown in Figure 2B. The value of DNAa & DNAb was significantly higher than that of single DNA, indicating that under the same concentration of RNA, the sensitivity of single DNA to RNA was significantly lower than that of CSSD.

To assess the viability of the sensor, we carried out experiments in parallel settings and acquired data, as shown in Figure 2C. When DEPC-treated water was added, the green curve shows its fluorescence intensity. When miRNA was added, the fluorescence intensity curve decreased significantly, indicating that miRNA was indeed hybridizing with DNA, and DNA was hydrolyzed by DSN enzymes, resulting in a decrease in fluorescence intensity. Therefore, this strategy can be used to detect miRNA.

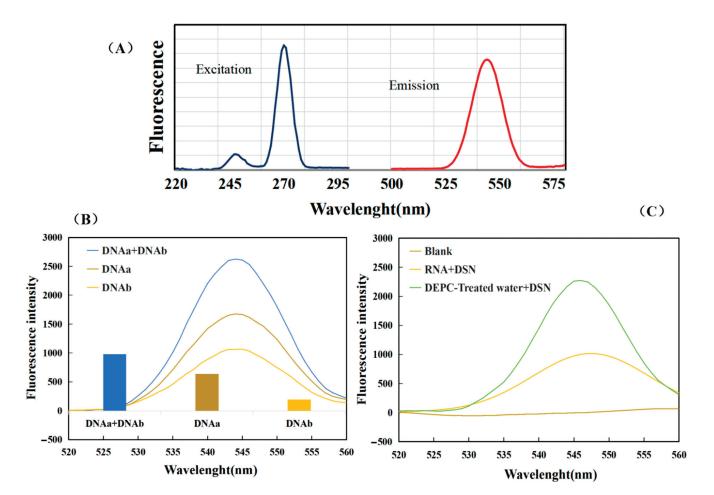


Figure 2. Feasibility analysis. (A) Excitation spectrum and emission spectrum of the sample to be measured. (B) The blue curve represents the fluorescence intensity with the addition of 5 μ L DNAa and 5 μ L DNAb, the brown curve represents the fluorescence intensity with the addition of only 10 μ L DNAa, and the yellow curve represents the fluorescence intensity with the addition of only 10 μ L DNAb. (C) The green curve shows the fluorescence signal curve with the addition of treated water and DSN enzyme, and the yellow curve shows the fluorescence signal curve with the addition of miRNA and DSN enzyme.

3.3. Optimization of Experimental Conditions

In this experiment, the generation of CSSD and the hydrolysis of DSN enzyme were both critical. Hence, it was imperative to fine-tune the key experimental parameters influencing fluorescence intensity in order to enhance the biosensor's efficacy. Here, we tested five key conditions, including the amount of GR dye, the incubation time of CSSD, the amount of DSN enzyme, the reaction temperature of the DSN enzyme, and the hydrolysis time of CSSD. Each group experiment had only one variable and conducted a series of three parallel experiments to reduce influence of random error. The other conditions were a DNAa (1 μ M) dose of 2.5 μ L, a DNAb (1 μ M) dose of 2.5 μ L, and a miR-34a (2 μ M) dose of 1.5 μ L.

3.3.1. Optimization of Incubation Time of CSSD

When the dosage of DNAa, DNAb, and miR-34a was determined, the amount of CSSD incubation was related to the level of fluorescence value F_0 , subsequently impacting the outcomes of successive experiments like reaction efficiency and signal output. The concentration of DNAa and DNAb configured and used in this experiment was 10 μ M, and the incubation times of DNAa and DNAb were different under the same experimental conditions. A total of five experiments were set up, and each experiment was repeated

three times. F₀ was the fluorescence signal value for adding only DNA and GR, and F was the fluorescence value for adding RNA and DSN enzymes. As shown in Figure 3A, with the increase in incubation time, the fluorescence difference (F₀-F) first increased and then decreased. The maximum value was reached at 45 min. The reason for it might have been the excessive time of incubation, which likely loosened some bound structures to gradually separate. Finally, the optimal time was chosen to be 45 min.

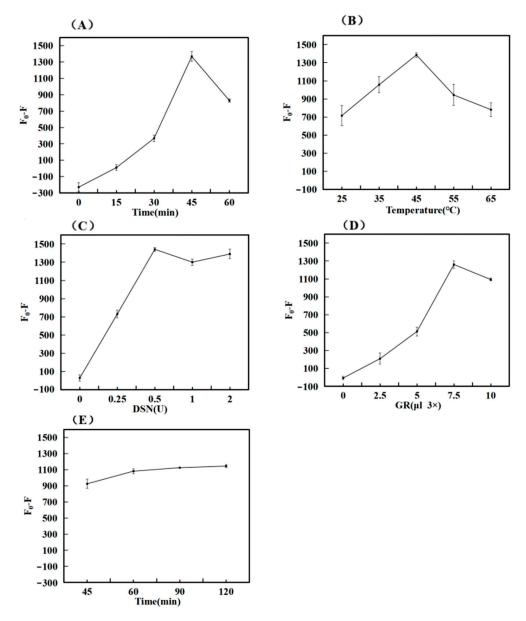


Figure 3. Optimization of experimental conditions. **(A)** Incubation Time of CSSD; **(B)** Reaction Temperature of DSN Enzyme; **(C)** Dosage of DSN Enzyme; **(D)** Dosage of GelRed Dye; **(E)** Time of DNA hydrolysis.

3.3.2. Optimization of DSN Enzyme Reaction Temperature

Temperature is one of the important factors affecting the fluorescence intensity of hybridization. When the temperature drops too low, it can cause the reaction time to stretch or lead to incomplete reactions. Conversely, excessively high temperatures might hinder the activity of the DSN enzyme, thereby impacting the outcomes of the experiment. To address this, we designed five experiments to be tested at different reaction temperatures, and each test was repeated three times. As illustrated in Figure 3B, the difference in fluorescence (F_0-F) shows an increasing trend between 25 °C and 45 °C, while it declines from 45 °C

to 65 $^{\circ}$ C. This latter decrease could be attributed to the elevated temperature impairing enzyme activity. Considering both the cost-effectiveness and practicality of the biosensor, we determined that 45 $^{\circ}$ C is optimal for the reaction.

3.3.3. Optimization of DSN Enzyme Dosage

The amount of DSN enzyme dosage was related to whether the DNA in the DNA–RNA hybridization chain could be completely hydrolyzed, thus affecting subsequent experiments. The concentration of DSN enzyme configured and used in this experiment was 1 $U/\mu L$. To establish the ideal concentration of the DSN enzyme, various amounts were introduced while maintaining consistent experimental conditions, ensuring that the volume of the test solution was standardized with the buffer solution. The outcomes of the tests are displayed in Figure 3C. As the quantity of DSN enzyme rose, the variation in fluorescence (F₀-F) also increased, peaking at 0.5 U. However, beyond this point, additional increases in DSN enzyme dosage did not significantly alter the fluorescence intensity. To optimize resources, the DSN enzyme concentration was therefore maintained at 0.5 U for the following experimental adjustments.

3.3.4. Optimization of GelRed Dye Dosage

The amount of GR affects the intensity of fluorescence signal detection. The concentration of GR dye used in the experiment was $3\times$. Under the same experimental conditions, GR dye with the same concentration but different amounts was added, and the volume of the tested solution was adjusted with ultra-pure water to be consistent. The experiment was repeated three times. After testing, the results obtained from the experiment are shown in Figure 3D. The difference in fluorescence (F₀-F) rose as the dosage of GR dye escalated, peaking at 7.5 μ L before tapering off with additional increments of the dye. This could be attributed to the possibility that at higher GR concentrations, there is less available binding space with DNA, which in turn leads to a reduction in the fluorescence signal. Therefore, the subsequent optimization experiment set the concentration of GR at 7.5 μ L.

3.3.5. Optimization of DNA Hydrolysis Time

The hydrolysis time of DNA is also the action time of the DSN enzyme. And the primary cause of the diminished fluorescence signal is the hydrolysis of DNA. When the hydrolysis time is too brief, it results in only a slight attenuation of the signal, ultimately compromising the sensor's sensitivity. Conversely, if the duration is excessively lengthy, it leads to an inefficient use of time. In order to determine the ideal hydrolysis time of DNA, we changed the reaction time of the DSN enzyme and inactivated the enzyme with high temperature after reaching the expected time. Upon analyzing the results, we illustrated our findings in Figure 3E. We observed that the discrepancy in fluorescence grew over time, although the rate of increase began to slow down. Ultimately, we decided that a 60 min hydrolysis period was optimal for DNA.

3.4. Analysis of Performance of miR-34a Biosensors

In this study, we designed a DSN enzyme-assisted fluorescent biosensor platform for miR-34a detection. To investigate the detection performance of the sensor, under optimal experimental conditions, different concentrations of miR-34a were detected. Based on the above, the optimized experimental conditions were as follows: 0.5 U DSN enzyme, 1 h DNA incubation time, 45 °C DSN enzyme reaction temperature, and 7.5 μL GR dosage, the fluorescence signal value was the best result. Ten sets of experiments were conducted to measure the fluorescence intensity across various concentrations of miR-34a, with each set repeated thrice for accuracy.

As shown in Figure 4A, this experiment is one in which the fluorescence value decreased with the increase in the target miRNA concentration. The standard detection curve in the sensitivity analysis was derived from the fluorescence value (F_0) of the light blue curve in Figure 4A minus the curve (F) of different concentrations. Ten sets of different

data (F_0 -F) were obtained, and the linear regression curves F_0 -F = 501.19C - 247.04 and R^2 = 0.9816 were obtained through simple fitting, indicating the high feasibility of this experiment. The lowest detection limit (LOD) was determined by tripling the standard deviation (3 σ) of the baseline signal. This was calculated using the formula 3 σ /S, where σ represents the baseline signal's standard deviation and S stands for the slope of the fitting line in Figure 4B. The detection limit was calculated to be 0.36 nM. The closer R^2 was to one, the higher the accuracy of detecting the target miRNA within the detection range, and the better photobleaching observed in this experiment. As shown in Table 2 and the sensitivity of this work, our work is superior to many previous works on miRNA detection systems.

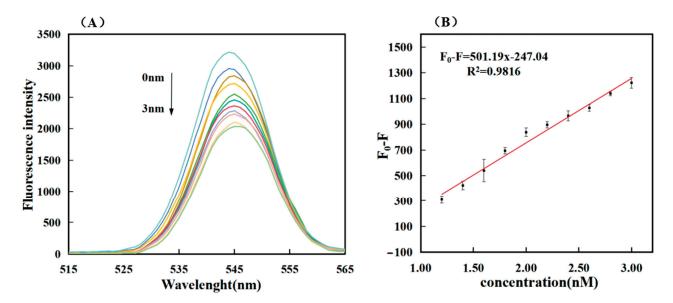


Figure 4. The fluorescence value of miR-34a was measured against the standard detection curve, and the concentrations of top to bottom substances in Figure (**A**) were 0 nM, 1.2 nM, 1.4 nM, 1.6 nM, 1.8 nM, 2.0 nM, 2.2 nM, 2.4 nM, 2.6 nM, 2.8 nM and 3.0 nM, respectively; Figure (**B**) Linear calibration curve of RNA and fluorescence intensity (correlation between RNA concentration and fluorescence signal strength). The error bars represent standard deviations obtained from from triplicate exper.

Table 2. Comparison of some other sensors for miR-34a detection.

Detection Technique	Linear Range	LOD	Reference
PAMAM Dendrimer Modified Electrodes	0–7.5 μg/mL	135 nM	[26]
Ionic Liquid Modified Single-use Electrodes	2–10 μg/mL	125 nM	[27]
Ionic Liquid-Modified Graphite Electrodes	2–10 μg/mL	109 nM	[28]
The Chemical Activation of PGE Surface	$0-2.5 \mu g/mL$	72.5 nM	[29]
Graphene Oxide Modified Graphite Electrodes	142–568 nM	41.2 nM	[30]
Electropolymerization	5–80 μg/mL	28.4 nM	[31]
Enzyme-assisted fluorescence biosensor	1.2–3 nM	0.36 nM	This work
Gold Nanoflower @Graphene Quantum Dots Probe	0.4–4 fM	0.1 fM	[32]
Catalytic Hairpin Assembly &DNA Walker	200 pM-10 nM	43.72 pM	[33]
Gold Nanostructures	0.1–1000 nM	39 pM	[5]

3.5. The Specificity of the Sensor

The experimental test of resistance to interference from other substances was also an important indicator of sensor performance. In order to see the influence of other miRNA sequences on the detection results, the signal values obtained from the detection of the target miR-34a, and other homologous miRNA were tested under the same optimized conditions. In theory, other non-targeted miRNAs cannot fully react with the designed DNA strand, and the efficiency of enzyme reactions will decrease. In order to verify

whether the sensor constructed in this experiment was selective for the target miR-34a, we performed specific detection experiments using different miRNAs. These included miR-155, miR-192, miR-15a, miRNA-122, miR-660, and miRNA-21. Among them, the concentration of miR-34a was 3 nM, and the concentration of other detected substances was 15 nM.

The findings from the experiment are illustrated in Figure 5, where (A) displays the fluorescence spectra for different miRNA detection. Upon examination of the figure, it is clear that the fluorescence values of the remaining miRNA are quite high, whereas the fluorescence level of miR-34a is significantly lower. This suggests that in the experiments, other types of miRNA had difficulty binding to circular DNA and performing enzymatic reactions. The DNA was not hydrolyzed, but other RNAs were stained with GR dye, resulting in enhanced fluorescence intensity. As can be seen from the bar graph in Figure 5B, after three tests, the fluorescence intensity of the target miR-34a decreased, resulting in a positive fluorescence intensity difference, while the staining of other miRNAs increased the fluorescence intensity, resulting in a negative fluorescence intensity difference. Thus, the target RNA was very different from other miRNAs. In summary, this approach exhibits strong discriminatory capacity towards other similar miRNAs, showcasing commendable specificity and selectivity.

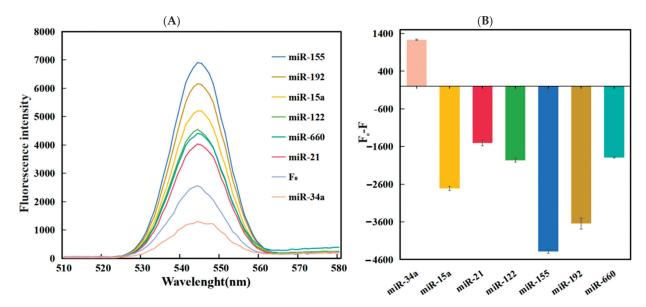


Figure 5. (**A**) Fluorescence spectrum plot of the relationship between miRNA and fluorescence intensity of different species; (**B**) Plot of different miRNAs of fluorescence intensity at the excitation wavelength of 545 nm.

3.6. Detection of miR-34a in Complex Environments

To further verify whether our method was suitable for detection in complex environments, we recovered three different concentrations of the target RNA (1.5 nM, 2.0 nM, 2.5 nM) by adding different concentrations of the target DNA to the reaction to study the reaction in the amplification system. As shown in Table 3, Added (nM) refers to the actual added RNA content. The estimated RNA content, marked as Found (nM), was obtained by calculating the obtained fluorescence value. The ratio between the two is marked as Recovery (%), and RSD is the standard deviation of the recovery rate. The calculated recoveries ranged from 92.8% to 105.6%, and the relative standard deviation recoveries were 5.2%, 6.3%, and 5.5%. The recovery fluctuated within a small range of 100%, meeting the experimental requirements. A lower standard deviation suggests that the test results were more consistent with the expected results, pointing to greater accuracy in the method. The experimental findings indicate that the proposed detection method holds promise for detecting miR-34a in real, intricate biological samples.

Table 3. Recovery of miR-34a in FBS (n = 3).

Sample	Added (nM)	Found (nM)	Recovery (%)	RSD (%)
		1.43	95.2	
1	1.5	1.53	102.0	5.2%
		1.58	105.6	
		2.10	105.0	
2	2.0	1.86	92.8	6.3%
		1.93	96.6	
		2.34	93.7	
3	2.5	2.58	103.0	5.5%
		2.35	94.0	

4. Conclusions

In conclusion, we have constructed an enzyme-assisted fluorescent biosensor, leveraging CSSD for the detection of miRNA. This sensor employs the DSN enzyme to break down the CSSD that is hybridized with the target miR-34a. As a result, miR-34a is freed to engage in the next phase, facilitating cycle amplification. During this process, the DNA's interaction with GR is disrupted, leading to a reduced fluorescence signal. Experimental findings demonstrated a clear distinction between the target miRNA and other RNA, illustrating strong selectivity and suitability for qualitative analysis. Notably, this method is capable of detecting miRNA at remarkably low concentrations, with a limit of detection as low as 0.36 nM and a linear detection range extending from 1.2 nM to 3 nM. Impressive outcomes were achieved even in intricate settings. This approach can also be adapted to identify various RNA biomarkers linked to illnesses by modifying the specified target recognition sequence. It offers considerable practicality and numerous potential applications in disease prevention and treatment, as well as for the quantitative measurement of RNA in medical research, paving the way for innovative studies.

Author Contributions: X.Y. and J.D. performed the experiments and wrote the paper; X.Y., Y.L., F.L. and J.L. analyzed the data, and J.D. conceived and designed the experiments. All authors have read and agreed to the published version of the manuscript.

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Article

NIR-Sensitive Squaraine Dye—Peptide Conjugate for Trypsin Fluorogenic Detection

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Abstract: Trypsin enzyme has gained recognition as a potential biomarker in several tumors, such as colorectal, gastric, and pancreatic cancer, highlighting its importance in disease diagnosis. In response to the demand for rapid, cost-effective, and real-time detection methods, we present an innovative strategy utilizing the design and synthesis of NIR-sensitive dye-peptide conjugate (SQ-3 PC) for the sensitive and selective monitoring of trypsin activity by fluorescence ON/OFF sensing. The current research deals with the design and synthesis of three unsymmetrical squaraine dyes SQ-1, SQ-2, and SQ-3 along with a dye-peptide conjugate SQ-3-PC as a trypsin-specific probe followed by their photophysical characterizations. The absorption spectral investigation conducted on both the dye alone and its corresponding dye-peptide conjugates in water, utilizing SQ-3 and SQ-3 PC respectively, reveals enhanced dye aggregation and pronounced fluorescence quenching compared to observations in DMSO solution. The absorption spectral investigation conducted on dye only and corresponding dye-peptide conjugates in water utilizing SQ-3 and SQ-3 PC, respectively, reveals not only the enhanced dye aggregation but also pronounced fluorescence quenching compared to that observed in the DMSO solution. The trypsin-specific probe SQ-3 PC demonstrated a fluorescence quenching efficiency of 61.8% in water attributed to the combined effect of aggregation-induced quenching (AIQ) and fluorescence resonance energy transfer (FRET). FRET was found to be dominant over AIQ. The trypsin-mediated hydrolysis of SQ-3 PC led to a rapid and efficient recovery of quenched fluorescence (5-fold increase in 30 min). Concentration-dependent changes in the fluorescence at the emission maximum of the dyes reveal that SQ-3 PC works as a trypsin enzyme-specific fluorescence biosensor with linearity up to 30 nM along with the limit of detection and limit of quantification of 1.07 nM and 3.25 nM, respectively.

Keywords: trypsin detection; aggregation-caused quenching; FRET; NIR detection; squaraine dye; dye–peptide conjugate

1. Introduction

Proteases are enzymes that catalyze the hydrolysis of peptide bonds, playing crucial roles in various physiological and pathological processes [1]. Among these, trypsin, a serine protease produced in the pancreas, is particularly significant due to its role in digestion and its involvement in pancreatic disorders such as pancreatitis and pancreatic cancer [2–4]. Trypsin facilitates the detachment of adherent cells and is therefore used therapeutically in tissue regeneration and as an anti-inflammatory agent [5]. The expression of trypsin is significantly elevated in several human cancer cells found in the stomach, colon, lung, and breast [6]. Recently, trypsin has gained recognition as a potential marker for cancer patients [7,8]. This is because the dysregulation of trypsin has been associated with the development of several tumors, such as colorectal cancer (CRC), gastric cancer, and pancreatic cancer [9,10]. In addition, trypsin is often used in the field of food chemistry for protein identification purposes [11]. In light of its biological significance, there is a significant

demand for inventive and accessible assays for trypsin. These assays are essential for the development of effective diagnostic and therapeutic techniques for the treatment of pancreatic diseases and as applications in the proteomics field [12,13]. Various techniques have been developed for trypsin detection, including mass spectroscopy [14,15], gel electrophoresis [16], enzyme-linked immunosorbent assay (ELISA) [17,18], colorimetric [19,20], and fluorometric approaches [21,22]. Among these, fluorogenic sensing technologies have gained prominence due to their simplicity, accessibility, and visual detection capabilities. These methods eliminate the need for sophisticated instrumentation and complex analytical procedures, contributing to their widespread adoption in various fields of study.

Capitalizing on these advantages, researchers have developed a diverse array of sensors employing distinctive fluorescence approaches, specifically designed and validated for the precise and sensitive detection of trypsin activity [23,24]. Recent advances in protease assay methods have focused on developing immunoassays using antibodies to bind target proteases [25,26]. However, while effective for quantifying protease levels, these immunoassays are limited in their ability to map protease activity and its correlation with disease progression. To address this limitation, alternative approaches have been explored, utilizing tailored peptide substrates with optical detection techniques such as absorbance or fluorescence spectroscopy [27–29]. These techniques leverage protease-substrate interactions to measure enzymatic activity directly. One of the challenges in developing fluorescence-based protease assays is the interference from the autofluorescence of biological samples and potential photo-damage to biological species due to high-energy excitation. The use of fluorophores that emit in the near-infrared (NIR) wavelength region offers a solution to these issues, providing increased sensitivity and the potential for deep-tissue bio-imaging [30,31].

In this context, NIR-sensitive squaraine dyes have emerged as promising candidates for biosensing applications due to their exceptional physical and chemical attributes, including strong absorption bands, high molar absorption coefficients, elevated quantum yields, and excellent photostability [32,33]. The absorption and emission properties of squaraine dyes can be tuned from the visible to NIR wavelength region by carefully selecting donor groups with controllable π -conjugation [34,35]. Additionally, SQ dyes exhibit a pronounced aggregation-caused quenching (AIQ) due to the presence of extended π -conjugation leading to a planer molecular framework. This allows for the modulation of their fluorescence properties by controlling their aggregation and dispersion states [36]. This characteristic has been utilized to develop a variety of fluorescent probes, including those for detecting proteins, enzymes, and ions. For instance, a BSA-SQ assembly was designed as a fluorescent "on-off" probe for pepsin detection, where the aggregation of SQ led to fluorescence quenching in the pepsin-catalyzed hydrolysis of BSA [37]. The versatility of SQ dyes in biosensing applications extends to protein detection [38], where their binding to proteins like bovine serum albumin disrupts SQ aggregation, leading to fluorescence enhancement. Furthermore, peptide-conjugated SQ probes based on Förster Resonance Energy Transfer (FRET) have been designed for enzyme detection [39], where enzymatic cleavage terminates the resonance energy transfer process and subsequently increases the fluorescence intensity. Aggregation-induced quenching (AIQ) and FRET can work synergistically to enhance biosensing systems. Hence, there is broad applicability of SQ dyes in developing sensitive and selective fluorescent probes for various biomolecules and ions, capitalizing on their unique aggregation based and energy transfer properties.

In this study, we present the development of an internally quenched homo-labeled fluorescent peptide substrate for the sensitive and selective detection of trypsin enzyme activity. The probe features two NIR-sensitive unsymmetrical squaraine dyes conjugated at each terminal of a trypsin-cleavable peptide sequence. This configuration exhibits fluorescence quenching in its intact state. Upon cleavage by trypsin, the probe demonstrates a significant enhancement of the fluorescence signal in the NIR region, enabling a sensitive detection of trypsin activity. The designed probe offers exceptional sensitivity, capable of detecting trypsin activity at exceedingly low concentrations. Moreover, it demonstrates

remarkable selectivity, with minimal interference from other proteases commonly found in complex biological samples. This approach has the potential to improve diagnostics and monitoring of trypsin-related disorders and also has potential applications in proteomics research and drug discovery.

2. Materials and Methods

2.1. Reagents and Materials

All chemicals, solvents, and reagents were of analytical or spectroscopic grade and used as supplied. Fmoc-protected amino acids, Rink amide 4-methylbenzhydrylamine (MBHA)resin,piperidine,1-((Dimethylamino)(dimethyliminio)methyl)-1H-benzo[d][1,2,3] triazole 3-oxide hexafluorophosphate (HBTU), 1H-Benzo[d][1,2,3]triazol-1-ol (HOBt)·H₂O, diisopropylethylamine (DIEA), trifluoroacetic acid (TFA), and 4 M HCl/Dioxane were purchased from Watanabe Chemical Industries, Ltd. (Hiroshima, Japan). Enzymes (α -chymotrypsin, α -trypsin from bovine pancreas, pancreatic elastase, horseradish peroxidase, bovine serum albumin, and papain) were obtained from Sigma-Aldrich Co., LLC (Tokyo, Japan). All other reagents and solvents were procured from Wako Pure Chemical Industries, Ltd. (Osaka, Japan).

2.2. Theoretical Structure Calculation

The theoretical quantum chemical calculation for the structural optimization of squaraine dye—peptide conjugate (**SQ-3 PC**) was conducted on a Dell workstation (8 Core) using the Gaussian G16 program package [40]. The calculation was performed for the isolated molecule in the gaseous state. In the theoretical calculations, judicious selection of a suitable basis set, theory, and functional are required to make a logical balance between the computation cost and accuracy. Considering the huge number of atoms in the molecule, theoretical calculation for structure optimization was conducted on HF/6-31G level of theory to save computation time.

2.3. Instrumentation

Nuclear magnetic resonance (NMR) spectroscopy (500 MHz for 1H NMR) and TOF/FAB-mass spectroscopy in positive ion monitoring mode were used for structural characterization. UV–visible–NIR absorption spectra were recorded on a JASCO V-530 UV/VIS spectrophotometer (JASCO Corporation, Tokyo, Japan). Fluorescence emission spectra were obtained using a JASCO FP-6600 spectrofluorometer (JASCO Corporation, Tokyo, Japan). Analytical HPLC was performed on a Hitachi L-7100 system (Hitachi High-Technologies Corporation, Tokyo, Japan) with an Xterra MS C18-5 μ M column (4.6 \times 250 mm; Waters Corporation, Milford, MA, USA). The sample was pre-incubated in a EYELA SLI-220 temperature-controlled incubator (Tokyo Rikakikai Co., Ltd., Tokyo, Japan) to ensure thermal equilibration. During spectroscopic measurements, the cuvette holder was equipped with a water-jacketed cell holder connected to a THERMO SUPPLIER EZL-80F circulating water bath system (TAITEC Corporation, Saitama, Japan) maintained at 37 °C.

2.4. Spectroscopic Measurements

Stock solutions of 100 μ M trypsin in phosphate buffer saline (PBS) (pH 7.2, 0.1 M) and 100 μ M **SQ-3 PC** in dimethyl sulfoxide (DMSO) were prepared and diluted as required. The fluorescence response of **SQ-3 PC** (5 μ M) to trypsin was evaluated in H₂O (pH 7.2, 2% DMSO). Fluorescence measurements were taken at 659 nm (excitation and emission slit widths of 5 nm and 6 nm, respectively) at 5 min and then every 10 min for 30 min. Trypsin concentrations ranged from 1 nM to 75 nM. Baseline fluorescence intensity (F₀) was determined using 5 μ M **SQ-3 PC** in H₂O with 2% DMSO.

2.5. Determination of Detection Limit

The limit of detection (LOD) and limit of quantification (LOQ) for trypsin-specific fluorescence biosensor was calculated using the following formula:

$$LOD = 3.3 (SD)/S \tag{1}$$

$$OQ = 10 (SD)/S$$
 (2)

where S is the slope of the linear part of the calibration curve and SD is the standard deviation of the sample for consecutive measurements.

2.6. Fluorescence Quenching Efficiency

Fluorescence quenching efficiency was calculated using the following equation:

Quenching efficiency (%) =
$$[1 - (F_{PC}/F_{dye})] \times 100$$
 (3)

where F_{dye} is the emission intensity of SQ-3 dye and F_{PC} is the emission intensity of SQ-3 PC at the same concentration in H_2O (2% DMSO).

3. Results

3.1. Design and Synthesis of Target Molecules

In this study, three unsymmetrical squaraine dyes (SQ-1, SQ-2, and SQ-3) and a dye-peptide conjugate (SQ-3 PC) were synthesized and investigated as NIR-sensitive fluorescent molecules for determining trypsin enzyme activity. The dyes are based on a squaraine core with two indole-based terminal groups (Figure 1). The molecular design strategy involved several key modifications to the squaraine core structure to enhance its photophysical properties. Iodine was introduced to investigate the heavy atom effect, potentially improving intersystem crossing. A methoxy group was added as an electrondonating substituent to optimize electronic distribution within the π -conjugated system. Additionally, a propionic acid moiety was incorporated at the N-position of one terminal indole ring to establish an intramolecular charge transfer (ICT) system and provide a reactive carboxylic acid for bioconjugation with peptides via amide bond formation. SQ-1 features an asymmetric substitution with an iodo group (-I) and a propionated alkyl chain(-CH₂CH₂COOH) on one indole group and a methoxy group and butyl chain(- $CH_2CH_2CH_3$) on the other. **SQ-2** is similar to **SQ-1** but has butyl chains on both indole groups. SQ-3 resembles SQ-1 but lacks the iodine substitution. The synthesis, structural characterization, and analytical data for these compounds and their intermediates are provided in the Supplementary Materials (Schemes S1 and S2).

Figure 1. Chemical structure of squaraine dyes (SQ-1, SQ-2, and SQ-3).

SQ-3 was identified as the optimal choice for the preparation of the dye–peptide conjugates due to its distinctive structural attributes, aggregation properties, and optical behavior. Although the substitution of hydrogen with a heavy atom such as iodine in

squaraine dyes has been previously reported to enhance fluorescence quantum yield [41], this modification did not prove to be significantly advantageous in the present study. The absence of iodine in SQ-3 avoids potential heavy metal effects and reduces steric hindrance, potentially facilitating a more efficient coupling to peptides and better interaction with the target enzyme trypsin. In a recent study (Soumya et al.) [42], a symmetrical iodinated squaraine dye is designed to be a photosensitizer for photodynamic therapy, where its stability is crucial for targeted action on tumor tissues. This stability may contribute to its resistance to enzyme hydrolysis. SQ-3 exhibits the highest aggregation index (AI) in PB of 3.38, attributed to the combination of electron-donating (-OCH₃) and electron-withdrawing (-COOH) groups promoting intermolecular interactions. The amino acid sequence Gln-Arg-Glu was chosen as the substrate for the trypsin enzyme based on the known specificity of trypsin [43]. Specifically, trypsin cleaves peptide bonds between the carboxyl group of arginine or lysine and the amino group of the adjacent amino acid. To optimize the design of our biosensing probe, we incorporated additional structural elements (Figure 2a). Free amino groups with β-Ala and Lys were integrated to serve dual purposes: first, as spacers to provide flexibility towards the substrate, and second, to facilitate the binding of the SQ-3 squaraine dye to the substrate.

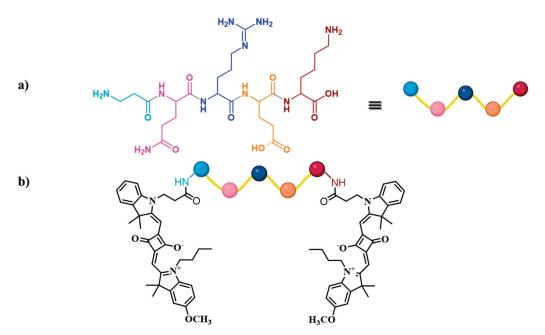
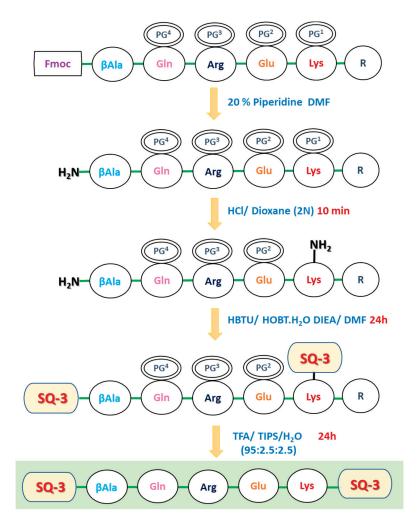


Figure 2. Chemical structure of **(a)** trypsin enzyme-specific peptide and **(b)** dye–peptide conjugate probe (**SQ-3 PC**) for the trypsin enzyme sensing.

The trypsin-specific probe **SQ-3-PC** (Figure 2b) was synthesized using solid-phase peptide synthesis (SPPS). The design incorporated a linear Gln-Arg-Glu tripeptide, chosen for trypsin specificity, with β -Alanine as a spacer and Lysine for the side-chain binding of the terminal squaraine dye. Synthesis began with Fmoc-Lys(Boc)-OH conjugation to Rink Amide MBHA Resin, followed by a stepwise peptide chain construction (Scheme 1). The resin-supported peptide (Fmoc- β Ala-Gln(trt)-Arg(Pbf)-Glu(o'tBu)-Lys(Boc)-Resin) underwent Fmoc deprotection with 20% piperidine and Boc removal with HCl/Dioxane (2 M). **SQ-3** was then conjugated at a 2.2:1 molar ratio using HOBt/HBTU as coupling reagents. The final cleavage and deprotection were achieved using a TFA/triisopropylsilane/H₂O (95:2.5:2.5) cocktail. The crude dye–peptide conjugate was precipitated with ether and purified by silica gel column chromatography (chloroform/methanol, 9:1), yielding the final probe. Successful synthesis was confirmed by Time-of-Flight (TOF) mass spectrometric analysis, which revealed a measured m/z of 1702.9112 [M + H]⁺ (calculated 1701.8901). The detailed synthesis procedure, reaction conditions, and analytical data, including mass spectrometry, are provided in the Supporting Information (Scheme S3).



Scheme 1. Schematic representation for the synthesis of dye–peptide conjugate (**SQ-3 PC**); R, Rink Amide MBHA Resin; PG^1 , Boc; PG^2 , o'tBu; PG^3 , Pbf; PG^4 , trt.

3.2. Optical Characterization of Squaraine Dyes and Dye-Peptide Conjugate

The photophysical properties of squaraine dyes (**SQ-1**, **SQ-2**, **SQ-3**, and **SQ-3 PC**) were characterized using optical absorption and fluorescence emission spectroscopies in DMSO. These measurements aimed to elucidate the impact of structural modifications on the optical properties. Figure 3a,b present the absorption and emission spectra, respectively, at a dye concentration of 5 μ M in DMSO along with the summarization of the deduced optical parameters such as absorption maximum (λ_{abs}), emission maximum (λ_{em}), molar extinction coefficients (ϵ), and Stokes shifts (Δ) in Table 1. All dyes exhibit strong and sharp optical absorption with the absorption maximum (λ_{abs}) between 652 nm and 659 nm with a weak vibronic shoulder between 600 nm and 620 nm, characteristic of squaraine dyes due to the π - π * transition. **SQ-3**, with -OCH₃ and -COOH groups, has a strong, narrow band in the visible to NIR wavelength region. **SQ-1** and **SQ-2**, containing iodine and -OCH₃ groups on the main ring, show relatively higher bathochromic shifts due to the heavy atom effect, broadening the absorption band and shifting it to longer wavelengths, which is a similar effect as observed by Mayerhöffer et al. [41].

These structural differences not only affect the optical absorption characteristics but also influence the emission properties of the dyes under investigation, as shown in Figure 3b. The enhanced emission of iodine-containing **SQ-1** and **SQ-2** is likely due to the iodine heavy atom effect, which may increase the molecular rigidity or reduce non-radiative decay. **SQ-1** with a single iodine atom achieves an optimal balance between the heavy atom effect and structural asymmetry, resulting in the highest emission intensity. The absence of iodine in **SQ-3** may lead to greater molecular flexibility, increased non-radiative decay,

and lower emission intensity. Nevertheless, all of the dyes depict a very small Stokes shift of only 8 nm–15 nm with similar shapes of absorption and emission spectra, which are indicative of the very comparable molecular configurations of the dye in the ground and excited states [44]. Interestingly, the optical behavior observed in the aqueous medium is very different compared to that observed in the DMSO solution. The electronic absorption spectrum of SQ-3 in H₂O containing 2% DMSO showed a broad absorption spectral feature with λ_{abs} at 653 nm and a highly pronounced vibronic shoulder appearing at 610 nm, as shown in Figure 4a, indicating an enhanced dye aggregation in the aqueous medium.

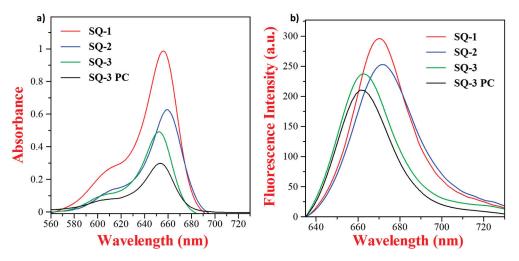


Figure 3. (a) Optical absorption and (b) fluorescence emission spectra of SQ-1, SQ-2, SQ-3, and SQ-3 PC in DMSO (5 μ M).

Table 1. Optical parameters for the squaraine dyes and dye–peptide conjugates in DMSO solution.

Squaraine Dye	λ_{abs}	λ_{em}	Δ	ϵ (M ⁻¹ cm ⁻¹)
SQ-1	656 nm	671 nm	15	1.98×10^{5}
SQ-2	659 nm	672 nm	13	1.26×10^{5}
SQ-3	652 nm	663 nm	11	1.00×10^{5}
SQ-3 PC	654 nm	662 nm	8	0.60×10^{5}

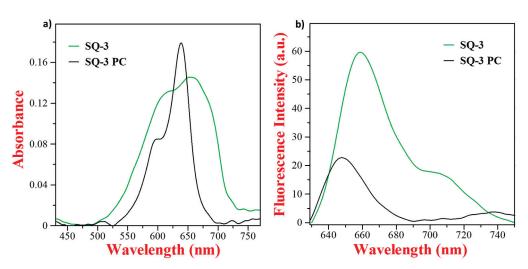


Figure 4. (a) Absorption and (b) emission Spectra of SQ-3 and SQ-3 PC (5 μ M) in H₂O (2% DMSO) solvent.

In contrast, SQ-3 PC in the same aqueous medium exhibits a blue-shifted, sharp absorption peak at 639 nm, with a prominent vibronic shoulder at 598 nm and less spectral broadening compared to SQ-3. This relatively lesser extent of molecular aggregation in SQ-3 PC compared to that of SQ-3 is attributed to the hampering of aggregation due to the presence of the oligopeptide chain and spacers. Squaraine dyes are prone to dye aggregation forming blue-shifted and co-facial H-aggregates and red-shifted head-to-tail packed J-aggregates (Figure 5) compared to monomeric dye absorption, which is influenced by the structure and molecular environment of the dyes' planar molecular structure [45]. The vibronic shoulder in squaraine dyes is very sensitive to dye aggregation and quite often has been used to index the relative extent of dye aggregation in different molecular environments. In the electronic absorption spectrum of squaraine dyes, the ratio of the absorbance around 600 nm associated with vibronic shoulder and around 650 nm associated with the monomeric dye absorption has been widely utilized to estimate the relative extent of the dye aggregation [46]. A perusal of the normalized absorption spectra of SQ-3 and SQ-3 PC in DMSO and water (Supporting Information Figure S1) reveals that the extent of aggregation of both of the dyes in DMSO is almost similar (0.26), while the behavior is very different in water, with nearly a doubled aggregation index for SQ-3 dye (0.91) compared to that of SQ-3 PC (0.46). At the same time, the nature of the aggregates is primarily H-type [45] in SQ-3 PC, while SQ-3 demonstrated the presence of both the Hand J- aggregates [47]. The unexpectedly low molar absorption coefficient of SQ-3 PC, despite containing two squaraine moieties, can be attributed to intramolecular electronic coupling leading to hypochromism, non-optimal conformations due to the flexible linker, and specific solute-solvent interactions in DMSO. These factors collectively contribute to a reduced absorption intensity compared to the expectations for two independent squaraine chromophores.

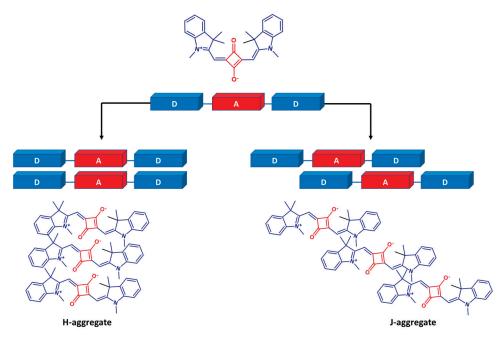


Figure 5. Schematic representation of H- and J-aggregate formation in squaraine dyes.

Notably, in the aqueous solution, **SQ-3-PC** and **SQ-3** exhibited a quenched fluorescence nearly ten times and five times higher (Figure 4b), respectively, than that observed in the DMSO solution. The quantum yield of **SQ-3** and **SQ-3** PC in water was calculated and found to be 0.006 and 0.0026. This is associated with aggregation-induced quenching (AIQ) due to the enhanced intermolecular interaction in water facilitated by hydrogen bonding between the free -COOH of the dye and water molecules. It is interesting to note that despite the nearly halved aggregation in **SQ-3** PC in water compared to that of pure dye **SQ-3**, the

quenching in fluorescence for SQ-3 PC was about three times higher compared to that of **SQ-3**. This clearly corroborates that, apart from AIQ, there might be an alternative pathway for fluorescence quenching in the case of SQ-3 PC. A second possibility is the quenching caused by fluorescence resonance energy transfer (FRET). To have FRET-based quenching, two criteria, such as the presence of the two terminal fluorophores within the Forster radius of typically 2 nm-20 nm and a sufficient overlap between the absorption spectrum of one fluorophore with the emission spectrum of another one, need to be fulfilled [41]. Due to a very small Stokes shift of only 11 nm in the case of SQ-3 dye in DMSO, there is a sufficient overlap between the absorption and emission spectrum of SQ-3, as shown in Figure 6a fulfilling the first criterion of FRET in the case of SQ-3 PC. Interestingly, the distance between the two faces of SQ-3 dye in the dye-peptide conjugate SQ-3 PC is only 1.92 nm (Figure 6b), which is well within the typical Förster radius, ensuring the efficient energy transfer from one dye molecule to another, fulfilling the second criterion for FRET to occur. By analyzing the fluorescence intensities of the dye (SQ-3) and the dye-peptide conjugate (SQ-3 PC), we determined that the fluorescence quenching efficiency was 61.8% in H₂O (2% DMSO). Therefore, comparing the AIQ and FRET mechanisms present in SQ-3 PC, FRET seems to be more dominant compared to AIQ. The property of conjugation allowed us to investigate the potential application of dye-peptide conjugates for enzyme activity monitoring using fluorescence-based ON/OFF biosensing.

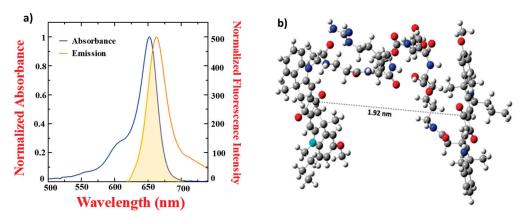


Figure 6. (a) Normalized absorption and emission spectra of **SQ-3** in DMSO and (b) optimized molecular structure of **SQ-3** PC calculated using Gaussian G16 program.

3.3. Enzymatic Hydrolysis of SQ-3-PC with Trypsin

The newly designed NIR-sensitive probe SQ-3-PC exhibited a highly quenched fluorescence signal caused by the AIQ and FRET, which later played a dominant role. However, after incubation with the trypsin enzyme, there was a substantial increase in fluorescence at 654 nm associated with monomeric absorption. Upon the addition of varying concentrations of this enzyme prepared in the phosphate buffer with a (pH of 7.4, 0.1 mM) in a 5 μ M solution of SQ-3-PC in H₂O (2% DMSO) followed by incubation for 30 min, there was a significant restoration of quenched fluorescence of the dye, as shown in Figure 7a. A perusal of Figure 7b shows that there was a significant increase in the fluorescence of the dye within 5 min. After 30 min, the fluorescence reached saturation, resulting in a 5-fold increase compared to the fluorescence in the absence of the trypsin enzyme. The observed enhancement in fluorescence intensity subsequent to enzymatic hydrolysis can be attributed to the spatial separation of the SQ-3 fluorophores beyond their Förster radius. This increased intermolecular distance effectively terminates the resonance energy transfer process, resulting in the cessation of FRET and the concomitant restoration of fluorescence emission. Figure 7c,d show that F₀ indicates the initial fluorescence intensity in the absence of the enzyme, where F represents the fluorescence intensity at a specified point of time after the enzyme is added. Figure 7c demonstrates the fluorescence response of the squaraine-based probe SQ-3 PC (5 μM) to varying concentrations of the target enzyme

(0-75 nM). The fluorescence intensity changes as a function of enzyme concentration, with higher enzyme levels inducing more rapid and pronounced spectral shifts. This indicates a direct correlation between the enzyme concentration and probe activation. Concentrationdependent changes in the fluorescence (F/F_0) at the emission maximum, as shown in Figure 7d, reveal that SQ-3 PC works as a trypsin enzyme-specific NIR fluorogenic probe with the LOD and LOQ of 1.07 nM and 3.25 nM, respectively. The linear regression equation was y = 0.11327x + 1.05815, $R^2 = 0.97043$ (Figure S2). Our innovative approach for trypsin detection demonstrates significant advancements over existing methodologies. The developed probe exhibits a remarkably low LOD, which is comparable to or surpasses the most sensitive techniques reported in the recent literature (Table S1). In healthy humans, the concentration of trypsin varies in serum and in the intestine. In serum under fasting conditions, the concentration of trypsin varies from 4 nm to 30 nM [48,49]. Pancreatic diseases such as cystic fibrosis, acute pancreatitis, or the acute phase of chronic pancreatitis are associated with an increased trypsin level of 2.1–71.42 nM in the serum of patients [50]. Thus, the synthesized probe is able to measure the trypsin level in healthy individuals and individuals with a disease associated with the hyper- and hypo-activities of the trypsin enzyme. Furthermore, the detection of trypsin in the NIR wavelength regions allows the utilization of direct body fluids without prior sample processing, owing to the highly diminished autofluorescence.

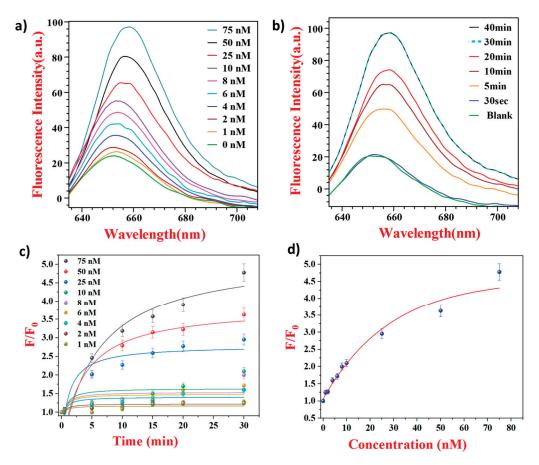


Figure 7. (a) Concentration-dependent fluorescence spectra of SQ-3 PC (5 μ M) in H₂O (2% DMSO) for 30 min with the addition of 0 to 75 nM Trypsin, (b) time-dependent fluorescence spectra of SQ-3PC (5 μ M) in H₂O (2% DMSO) for 30 min with the addition of 75 nM of trypsin enzyme, (c) fluorescence spectra of SQ-3 PC with change in fluorescence intensity as a function of time with different concentrations of enzyme for a fixed concentration of SQ-3 PC (5 μ M), and (d) plot of ratio of fluorescence intensity (F/F₀) as a function of enzyme concentration.

3.4. Enzyme Selectivity of the Enzyme Probe SQ-3-PC

Apart from the sensitivity and detection limit, the selectivity of a probe is highly desired for targeted disease diagnosis and point-of-care testing devices. The selectivity of the present probe towards trypsin was evaluated by analyzing its response in the presence of potentially interfering enzymes. A 5 μ M solution of the newly designed probe SQ-3 PC in H₂O (2% DMSO) was incubated with various enzymes, including trypsin, papain, pancreatic elastase, horseradish peroxidase, chymotrypsin, and bovine serum albumin (BSA), each at a concentration of 75 nM, for 30 min. As shown in Figure 8, the emission intensity at 654 nm remained relatively constant in the presence of the competing enzymes. However, a significant increase in emission intensity was observed upon the addition of trypsin. This marked change in fluorescence signal highlights the probe's high selectivity for trypsin, demonstrating its ability to specifically detect trypsin activity while remaining unresponsive to other enzymes present in the sample. These results underscore the potential of this biosensing approach for accurate trypsin detection in complex biological environments.

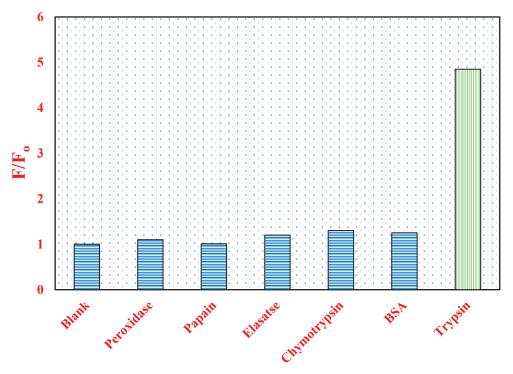


Figure 8. Fluorescence response (F/F₀) of the probe **SQ-3 PC** (5 μ M) in H₂O (2% DMSO) towards different potential and competing enzymes (75 nM).

4. Conclusions

In conclusion, we have successfully designed and synthesized novel unsymmetrical squaraine dyes and a potential dye–peptide conjugate for the sensitive and selective detection of the trypsin enzyme. In the aqueous medium, the dye–peptide conjugate **SQ-3 PC** demonstrates reduced dye aggregation but pronounced fluorescence quenching compared to its constituent dye **SQ-3** alone. A fluorescence quenching of 61.8% in the case of **SQ-3 PC** in an aqueous medium compared to its DMSO counterpart is associated with both the AIQ and FRET, where FRET was demonstrated to be the dominant phenomenon for this pronounced fluorescence quenching. The hydrolysis of the designed probe **SQ-3 PC** with trypsin enzyme restores an appreciable amount of the quenched fluorescence within 5 min, and the enhancement in the fluorescence saturates after 30 min. Notably, the probe exhibited high specificity and good selectivity towards trypsin when tested against a panel of potentially interfering enzymes, including BSA, papain, chymotrypsin, horseradish peroxidase, and pancreatic elastase. This selectivity is crucial for the potential application

of the probe in complex biological samples, though further validation in such matrices remains an important area for future investigation.

Supplementary Materials: The following supporting information can be downloaded at: https://www.mdpi.com/article/10.3390/bios14100458/s1, Scheme S1: Scheme for synthesis for unsymmetrical squaraine dye SQ-1, SQ-2; Scheme S2: Scheme for synthesis for unsymmetrical squaraine dye SQ-3; Scheme S3. Scheme for the synthesis of peptide sequence; Figure S1: Normalized Absorption Spectra of SQ-3 and SQ-3 PC in DMSO (solid line) and H_2O (2% DMSO) (dash line) at a concentration of 5 μ M.; Figure S2: Linear correction curve of ratio of fluorescence intensity against trypsin concentration (0 nM to 10 nM). Table S1: Comparison of different methods for the determination of trypsin.

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Communication

All-Dielectric Metasurface-Based Terahertz Molecular Fingerprint Sensor for Trace Cinnamoylglycine Detection

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Abstract: Terahertz (THZ) spectroscopy has emerged as a superior label-free sensing technology in the detection, identification, and quantification of biomolecules in various biological samples. However, the limitations in identification and discrimination sensitivity of current methods impede the wider adoption of this technology. In this article, a meticulously designed metasurface is proposed for molecular fingerprint enhancement, consisting of a periodic array of lithium tantalate triangular prism tetramers arranged in a square quartz lattice. The physical mechanism is explained by the finite-difference time-domain (FDTD) method. The metasurface achieves a high quality factor (Q-factor) of 231 and demonstrates excellent THz sensing capabilities with a figure of merit (FoM) of 609. By varying the incident angle of the THz wave, the molecular fingerprint signal is strengthened, enabling the highly sensitive detection of trace amounts of analyte. Consequently, cinnamoylglycine can be detected with a sensitivity limit as low as 1.23 μg·cm⁻². This study offers critical insights into the advanced application of THz waves in biomedicine, particularly for the detection of urinary biomarkers in various diseases, including gestational diabetes mellitus (GDM).

Keywords: terahertz; molecular fingerprint sensor; all-dielectric metasurface; cinnamoylglycine; gestational diabetes mellitus

1. Introduction

Terahertz (THz) spectroscopy, a portion of the electromagnetic spectrum, deals with radiation frequencies from 0.1 THz to 10 THz. This radiation can penetrate a wide variety of non-conducting materials and exhibits non-ionizing properties [1]. Consequently, it is acclaimed as a potent spectroscopic technique due to its label-free and non-destructive attributes. Previous research in THz time-domain spectroscopy has demonstrated it to be a state-of-the-art technology capable of providing distinct molecular fingerprints by detecting vibrational and rotational modes sensitive to molecular structure and environmental factors within the THz frequency range [2,3]. This underscores its substantial potential for applications in biological detection [3,4]. However, traditional THz spectroscopy encounters limitations, notably its reduced sensitivity for detecting trace analytes, which impedes its extensive application. This issue arises because the dimensions of target molecules are significantly smaller than the micrometer scale, whereas the wavelength of THz waves spans approximately from 30 μm to 3000 μm [5], leading to a remarkable mismatch between the absorption cross-section of the analyte and the wavelength. Hence, the interactions between them are too feeble to generate adequate vibrational absorption signals, necessitating a substantial quantity of analyte [6,7]. Recent studies have proposed various approaches to address this issue, encompassing surface plasmon resonances (SPRs) [8], electromagnetic subwavelength structures [9], waveguides [10], metal hole arrays [11], and graphene metamaterials [12]. These sensors can achieve sub-wavelength spatial resolution to enhance detection sensitivity [5]. But they are constrained by substantial ohmic losses and high dispersion of the metal materials [13,14], leading to limited improvements in sensitivity [15,16].

All-dielectric metasurface has been demonstrated as an alternative to traditional metallic metasurfaces. It is based on high-refractive-index dielectrics and leverages Mie-type resonances rather than surface plasmon resonances to obtain strong magnetic resonances as well as electric resonances, thereby enhancing the interaction between electromagnetic waves and materials [17-19]. This innovative design significantly reduces energy losses while elevating the quality factor (Q-factor) of the resonators, thus greatly improving detection sensitivity [17,20]. Furthermore, intrinsic heating can alter the local refractive index and potentially damage biomolecules. A dielectric metasurface can mitigate the plasmonic ohmic loss and heating issues, thereby enabling better repeatability and stability for biosensing [21]. Currently, a precise angle-scanning strategy has been proposed, which is based on guided-mode resonance theory [22]. By monitoring the transmission characteristics of THz waves at different incident angles, richer spectral information can be obtained, thereby improving the detection capability of trace analytes [22,23]. Employing this strategy, a previous study successfully identified lactose and glucose with detection limits of 1.53 μg·cm⁻² and 1.54 μg·cm⁻², respectively [24]. On this basis, maintaining the polarization direction can further refine the enhancement effect and broaden the bandwidth [24].

Cinnamoylglycine, a metabolic byproduct of cinnamic acid metabolism, is derived from dietary sources such as cinnamon [25]. In the human body, it is produced by gut microbes [26] and excreted in urine [27]. Gestational diabetes mellitus (GDM), which involves significant metabolic alterations [28,29], can pose severe health risks to both mother and fetus, such as preeclampsia [30], preterm birth [31], and increased risk of type-2 diabetes later in life [32]. These harms highlight the need for efficient GDM diagnostic methods for regular monitoring and timely intervention. Current methods, including fasting plasma glucose [33,34], oral glucose tolerance testing [35], and the detection of glycated hemoglobin [36], have limitations such as the need for fasting, multiple blood draws, and variability due to external factors. In contrast, studies have suggested a link between urinary cinnamoylglycine levels and diabetes, with diabetic patients exhibiting a clearance rate of 77 mL/min compared to 317 mL/min in healthy controls [37-39]. This proves urinary cinnamoylglycine levels are a practicable biomarker for monitoring metabolic changes associated with GDM and offer a non-invasive alternative. But existing detection methods for cinnamoylglycine, such as high-performance liquid chromatography (HPLC), are time-consuming and costly due to the pretreatment and processing of such highly complex matrices [40]. This underscores the critical need for a rapid, convenient, and accurate detection method for urinary cinnamoylglycine, thereby enhancing the overall management and prognosis of GDM. THz spectroscopy has proven effective for analyzing urine samples and biomarkers [41]. Furthermore, the use of metamaterials has created additional opportunities over the past decade, such as the detection of proteins [42,43] and cancer biomarkers [44]. Therefore, this study introduces a THz sensor to urinary cinnamoylglycine detection, utilizing a sophisticatedly designed all-dielectric metasurface to enhance molecular fingerprints and improve detection sensitivity and specificity.

In this study, we propose an all-dielectric metasurface-based THz sensor, utilizing a periodic symmetric tetramer of lithium tantalate triangular prisms on a quartz substrate. This periodic unit structure, composed of the four high-index tetramer clusters, can induce a strong magnetic dipole (MD) resonance, with a high Q-factor of 231 and figure of merit (FoM) of 609, thus significantly enhancing the molecular fingerprint for detecting trace amounts of analytes. By manipulating the incident angle of a THz wave, a broad envelope curve can be measured due to the shift of resonance frequency, which corresponds to the absorption resonances of the analyte. Consequently, this innovative sensor allows for the detection of cinnamoylglycine, with a detection limit as low as 1.23 $\mu g \cdot cm^{-2}$. Moreover, the proposed metasurface can be efficiently fabricated using micro/nanotechnology, and the testing can be performed with a terahertz time-domain spectrometer [45].

2. Structure and Design

Figure 1 depicts the schematic of the proposed all-dielectric metasurface. It consists of a periodic array of tetramer clusters arranged on a square quartz substrate. Each tetramer cluster is formed by four high-index triangular prisms. The dielectric property is a crucial characteristic of a material, as it fundamentally defines how the material responds to electromagnetic waves. It uniquely determines how radiation propagates through the material by describing its ability to polarize in response to an electric field, thus influencing the propagation of waves within the medium [46]. Lithium tantalate is chosen as the building material for the tetramer clusters due to its minimal imaginary part, resulting in negligible dissipation loss within the relevant frequency range [47]. The complex dielectric permittivity of lithium tantalate can be described as a Lorentz-type dispersion, as is described by the following equation [48]:

$$\varepsilon = \varepsilon_{\infty} \frac{\omega^2 - \omega_L^2 + i\omega\gamma}{\omega^2 - \omega_T^2 + i\omega\gamma} \,, \tag{1}$$

In the equation, the transverse and longitudinal optical phonon frequencies are $\omega_T/2\pi=26.7$ THz and $\omega_L/2\pi=46.9$ THz, respectively. The damping factor is $\gamma/2=0.94$ THz and the quartz substrate has a refractive index of n=2. For frequencies below the longitudinal optical phonon frequency, the permittivity of lithium tantalate is calculated to be 41.4 [49].

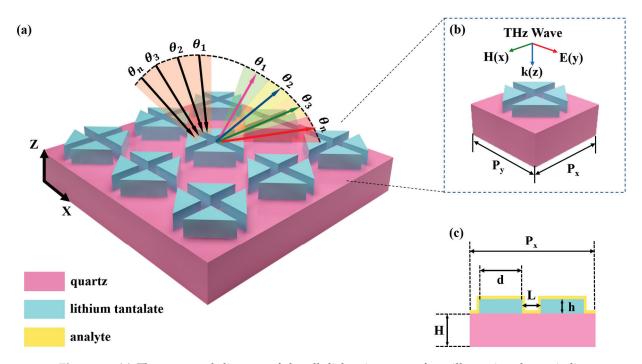


Figure 1. (a) The structural diagram of the all-dielectric metasurface, illustrating the periodic arrangement of the high-index triangular prism tetramer based on the quartz substrate; (b) a unit cell of the periodic structure with a y-polarized source incident downwards in the z direction; (c) the main view of the unit cell (y-z plane) and corresponding parameters.

The geometric parameters of the square unit cell are illustrated in Figure 1b,c, with the side length $P_x = P_y = 135~\mu m$ and the substrate thickness $H = 130~\mu m$. The prisms are shaped as isosceles right-angled triangles, of which the side length is $d = 90.4~\mu m$. As shown in the main view of the unit cell in Figure 1c, the distance between adjacent triangular prisms is $L = 33~\mu m$, and the height of each prism is $h = 49.6~\mu m$. By directing the incident THz wave to propagate downward along the z direction, the resonance dip within the transmission spectrum can be manipulated, resulting in the formation of an

envelope curve that captures the absorption resonances of analytes. This dynamic shifting of the resonance dip is essential for identifying the molecular fingerprints of various substances.

In order to evaluate the performance of the proposed sensor, extensive numerical simulations and spectral response analysis have been carried out with the commercially available three-dimensional finite-difference time-domain (FDTD) software. In the simulation, periodic boundary conditions were applied along the x and y directions to emulate an infinite array, while a perfectly matched layer boundary condition was set along the z-axis to prevent wave reflections. The simulations were conducted with a minimum meshing step of 1 nm to achieve high-resolution accuracy in modeling the electromagnetic fields.

3. Results and Discussion

Figure 2 illustrates the transmission spectra of the proposed metasurface and the electric and magnetic field at the resonance frequency. In Figure 2a, a single resonance at 0.624 THz is observed for both x-polarized and y-polarized waves under normal incidence due to the symmetric arrangement of the metasurface structure. However, as shown in Figure 2b, when the incident angle is adjusted to 37° , the sensor exhibits a strong sensitivity to polarization due to the breaking in symmetry, leading to a frequency shift. The resonance is measured at 0.487 THz for y-polarized waves, showing strong correspondence with the absorption peak of cinnamoylglycine and a high Q-factor of approximately 231. The Q-factor is defined as $Q = f_0/\Delta f$, where the full-width at half-maximum (FWHM) Δf is 2.11 GHz and the resonance frequency f_0 is 0.487 THz. In contrast, the frequency shift for x-polarized waves does not align sufficiently with the targeted absorption peak. Therefore, y-polarized waves were selected for angle-scanning to achieve frequency shifts and obtain the transmission envelope curve.

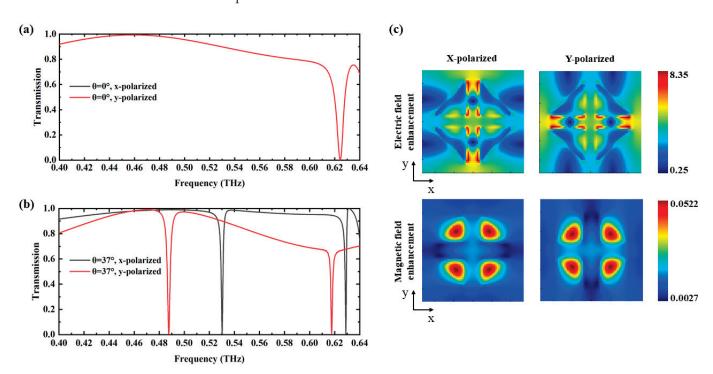


Figure 2. (a) Transmission spectra for x-polarized and y-polarized incident waves at 0° ; (b) transmission spectra for x-polarized and y-polarized incident waves at 37° ; (c) the electric and magnetic field distribution measured at the surface of the quartz substrate at vertical incidence. The left and right figures correspond to the x-polarized and y-polarized incident wave, respectively.

Figure 2c shows the electric and magnetic field distributions at the substrate surface in the x–y plane. The left and right panels correspond to x-polarized and y-polarized

incident waves, respectively. The resonance is identified as an MD mode, characterized by the collective response of four longitudinal MDs [50]. The electric field is predominantly localized in the center region of the cluster, indicating the excitation of the MD resonance in this area, which enhances the interaction between the incident THz waves and analytes. The magnetic field distribution confirms the presence of strong MD resonances, further supporting the sensor's sensitivity to different polarization states and incident angles. By leveraging these distinct field distributions, the sensor can detect analytes more effectively, which improves both the sensitivity and specificity of the detection process.

To investigate the combined spectral response and the relevancy between resonance dip and the incident angle, various angles were selected for analysis. Figure 3 illustrates that each incident angle corresponds to a distinct narrowband unity transmission, with the resonance frequency progressively decreasing as the incident angle increases. This frequency shift results from the alteration of the structure's symmetry due to the change in the incident angle. As this asymmetry intensifies, the frequency shift becomes more observable. Furthermore, modifying the parameters of the tetramer leads to a shift in the resonance frequency of the MD [49]. These influencing factors together create a comprehensive spectral cluster within the transmission spectrum, covering a broad frequency range that aligns closely with the absorption frequency of the analyte. This relevancy is crucial for enhancing the precision of sample identification, as it ensures that the spectral response could precisely mirror the unique absorption characteristics of the analyte.

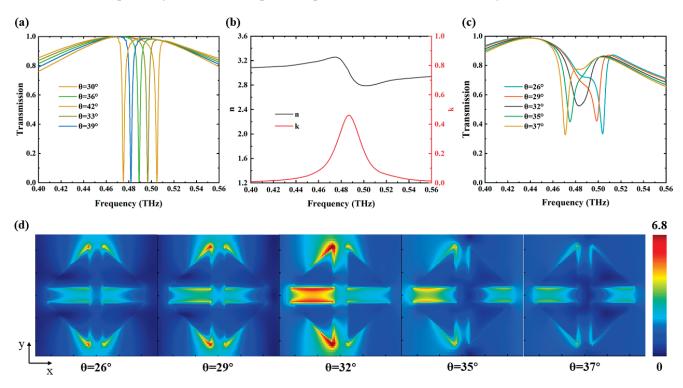


Figure 3. (a) Transmission spectra at different incident angles without any analyte; (b) the experimentally measured refractive index (n) and extinction coefficient (k) of cinnamoylglycine across the relevant frequency range; (c) transmission spectra at different incident angles with a 1 μm thick layer of analyte; (d) the electric field distribution measured at the substrate surface in the x–y plane at 0.487 THz for specific incident angles, corresponding to the transmission spectra shown in (c), respectively.

When the THz wave interacts directly with the metasurface in the absence of analyte, as shown in Figure 3a, both the minimum transmission value and the linewidth of the transmission spectra remain constant. In Figure 3b, the experimentally measured complex refractive index of cinnamoylglycine is presented, with data extracted using Fresnel formulas [51]. The real part, n, signifies the refractive index, whereas the imaginary

part, k, denotes the extinction coefficient. This measurement indicates that the fingerprint absorption peak for cinnamoylglycine is located at 0.487 THz.

The fingerprint detection capability of the metasurface is evaluated by observing the transmission envelope curve with a 1 μm thick cinnamoylglycine layer. Figure 3c displays a unique envelope that peaks at 0.487 THz at an incident angle of 32°, matching the extinction coefficient curve for cinnamoylglycine. This phenomenon is explained by the electric field distributions at specific angles, shown in Figure 3d. At 0.487 THz and an incident angle of 34°, a notable electric field enhancement is observed, with pronounced concentration in the central region of the structure. This enhancement and localization increase the likelihood of wave—matter interactions near the resonant frequency for the analyte on the metamaterial. In contrast, when the incident angle is reduced to 25° or increased to 42°, the electric field strength diminishes.

The sensing sensitivity (S) of the designed metasurface is evaluated based on $S = \Delta f/\sigma$. As is shown in Figure 3c, the resonant frequency at 37° is 0.471 THz, resulting in a frequency shift Δf of 0.0158 THz. With the 1 μ m thick cinnamoylglycine layer, the surface concentration is calculated as $\sigma = \rho \times h$, where ρ represents the volume density of cinnamoylglycine, noted as 1.23 g/cm³. Accordingly, $\sigma = 123 \, \mu g/cm^2$. Figure of merit (FoM) is used to assess the comprehensive performance of the sensor, which is defined as $FoM = S/FWHM = S \times Q/f_0$ [52]. Consequently, FoM is obtained as 609.

Figure 4 presents the comprehensive angle-scanning transmission spectra of the metasurface, demonstrating that this strategy effectively distinguishes cinnamoylglycine from other substances due to its distinct absorption resonances. In Figure 4a, the transmission spectra show stable frequency shifts across various incident angles ranging from 13° to 70° in the absence of any analyte, indicating a consistent performance of the metasurface. Conversely, Figure 4b presents the transmission spectra with a 1 μm thick cinnamoylglycine layer on the metasurface. The maximum transmission rate for each incident angle was extracted, and the results were fitted into an envelope curve, plotted as the red line in Figure 4b. The envelope curve peaks at 0.487 THz, which reaches approximately 56.85%, aligning with the characteristic fingerprint spectrum of cinnamoylglycine. The observed phenomenon is attributed to the optical loss of cinnamoylglycine, indicating that the angle-scanning strategy enhances detection sensitivity and accuracy by capturing comprehensive transmission envelopes, thereby facilitating precise analyte identification.

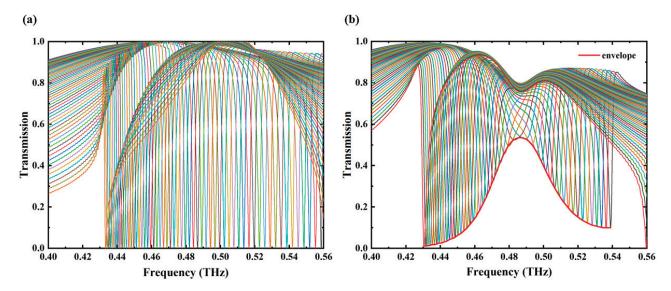


Figure 4. (a) Comprehensive transmission spectra without any analyte, with the incident angle ranging from 13° to 70° . Specifically, the rightmost line represents the transmission curve for an angle of 13° , while the leftmost line corresponds to 70° ; (b) comprehensive transmission spectra with 1 μ m thick cinnamoylglycine, with the incident angle ranging from 13° to 62° . The corresponding envelope curve has been plotted by red line in the figure.

Figure 5a depicts the transmission envelope curves for cinnamoylglycine across a series of layer thicknesses, spanning a frequency range from 0.40 THz to 0.56 THz. Notably, the transmission peak increases from 49.55% to 56.85% at 0.487 THz as the analyte layer thickness increases from 0.01 μ m to 1 μ m, whereas the transmission envelope remains approximately at zero in the absence of analyte. The detection limit is determined by $\sigma = \rho \times h$, among which the volume density of analyte is $\rho = 1.23 \text{g/cm}^3$, and the minimum thickness of analyte layer is $h = 0.01 \,\mu$ m. Consequently, the detection limit is 1.23 μ g/cm². It can be observed in Figure 5b that the transmission at 0.487 THz exhibits a linear relationship with the analyte thickness when the thickness exceeds 0.01 μ m. The fitted equation is y = 0.044x + 0.5242, where y represents the transmission and x denotes the analyte thickness. The correlation coefficient of approximately 0.99 indicates a strong linear dependence of the fitting line, suggesting that the thickness of cinnamoylglycine can be predicted from its transmission at 0.487 THz.

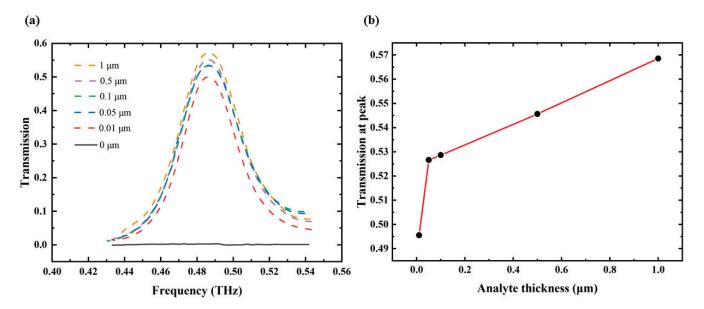


Figure 5. (a) Transmission envelope curves for analytes of varying thicknesses; (b) the relationship between the thickness of the analyte and the transmission at 0.487 THz.

4. Conclusions

In conclusion, this study presents a novel all-dielectric metasurface-based THz sensor for the sensitive detection of cinnamoylglycine, an emerging urinary biomarker for GDM. The proposed sensor, utilizing lithium tantalate triangular prism tetramers on a quartz substrate, achieves a high Q-factor of 231, significantly enhancing molecular fingerprint detection. The angle-scanning strategy enables broad spectral coverage and captures comprehensive transmission envelope curves that precisely align with specific analyte absorption frequencies. Simulation results indicate a detection limit as low as 1.23 μg·cm⁻², demonstrating the sensor's exceptional sensitivity and specificity. Considering the severe health risks associated with GDM, including preeclampsia and preterm birth, this approach offers a rapid, convenient, and accurate alternative to traditional diagnostic methods, with potential applications in detecting other biomarkers and monitoring various diseases. It meets the critical need for non-invasive and efficient GDM monitoring, facilitating regular monitoring and timely intervention. The successful identification of cinnamoylglycine underscores the robust performance of the metasurface in practical sensing applications, establishing a solid foundation for future research and advancements in THz spectroscopy and biomedical diagnostics.

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Q.X. and W.W.; resources, Y.S.; data curation, Q.X., M.S., W.W., and Y.S.; writing—original draft preparation, Q.X.; writing—review and editing, Q.X. and Y.S.; visualization, Q.X.; supervision, Y.S.; project administration, Y.S.; funding acquisition, Y.S. All authors have read and agreed to the published version of the manuscript.

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Article

Asterias forbesi-Inspired SERS Substrates for Wide-Range Detection of Uric Acid

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Abstract: Uric acid (UA), the final metabolite of purine, is primarily excreted through urine to maintain an appropriate concentration in the bloodstream. However, any malfunction in this process can lead to complications due to either deficiency or excess amount of UA. Hence, the development of a sensor platform with a wide-range detection is crucial. To realize this, we fabricated a surface-enhanced Raman spectroscopy (SERS) substrate inspired by a type of starfish with numerous protrusions, *Asterias forbesi*. The *Asterias forbesi*-inspired SERS (AF-SERS) substrate utilized an Au@Ag nanostructure and gold nanoparticles to mimic the leg and protrusion morphology of the starfish. This substrate exhibited excellent Raman performance due to numerous hotspots, demonstrating outstanding stability, reproducibility, and repeatability. In laboratory settings, we successfully detected UA down to a concentration of 1.16 nM (limit of detection) and demonstrated selectivity against various metabolites. In the experiments designed for real-world application, the AF-SERS substrate detected a broad range of UA concentrations, covering deficiencies and excesses, in both serum and urine samples. These results underscore the potential of the developed AF-SERS substrate as a practical detection platform for UA in real-world applications.

Keywords: uric acid; wide-range detection; bioinspired; gold nanoparticles; surface-enhanced Raman scattering

1. Introduction

Uric acid (UA), primarily synthesized in the liver, is an organic compound produced endogenously from purine metabolism in the human body. It circulates through the bloodstream and is excreted in urine via kidneys [1]. UA functions as a biomarker for oxidative stress, and the changes in its concentration are associated with various physiological and pathological conditions [2]. Generally, the normal concentration range of UA in the blood is approximately 200 to 500 μ M [3,4], and deviations from this range can lead to complications. If UA levels in the blood drop below the normal range (a condition known as hypouricemia), then the risk of acute kidney injury can significantly increase after intense physical activity. Additionally, low UA levels in the blood are associated with renal hypouricemia, which is linked to kidney stones due to an increased level of urinary UA, correlating with the impaired tubular reabsorption of UA [5]. Conversely, heightened UA levels, referred to as hyperuricemia, can result in the formation of urate crystals, which accumulate in joints, cartilage, and adjacent tissues, giving rise to conditions such as gout [6,7]. Additionally, hyperuricemia is associated with inflammatory diseases, kidney issues, and hypertension [8–11]. Therefore, detecting UA concentrations outside the normal range in both blood and urine is crucial for assessing health status and diagnosing diseases.

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Due to the critical importance of detecting UA levels, tests for UA in blood and urine are performed via standard procedures in hospitals. However, these tests have limitations, including the need for periodic sample collection and intricate preparation procedures, and these constraints have driven the development of various UA detection methods over the past decades [12-14]. Techniques involving electrochemistry [15,16], fluorescence [17], and spectroscopy [18,19] have been widely employed in UA detection. In particular, electrochemical detection methods are prevalent because of their effectiveness in UA detection. However, UA produces a redox signal similar to dopamine and ascorbic acid, which may cause selectivity problems [20-22]. Recently, spectroscopy based on surfaceenhanced Raman scattering (SERS) has garnered significant attention in the biosensing field due to its remarkable sensitivity and exceptional selectivity [23,24]. However, the current SERS-based platforms still face limitations in detecting UA in both deficiency and excess ranges [25]. To overcome this challenge, it is crucial to design nanostructures with high sensitivity and extensive surface areas. Fabricating nanostructures with numerous nanogaps and forming SERS substrates with broad surface areas capable of accommodating high concentrations of analytes are essential for addressing this limitation.

In this study, we drew inspiration from the morphology of *Asterias forbesi*, commonly known as the Forbes Sea Star, to enhance the performance of the SERS substrate. *Asterias forbesi* exhibits numerous protrusions on its torso and legs, resembling structures in which the SERS effect can be maximized by creating nanoscale hotspots. To replicate this structure, we strengthened the SERS effect by fabricating small-scale gold nanoparticles (GNPs) in a gold–silver bimetallic structure, imitating starfish legs and the protuberances of *Asterias forbesi*. This design maximized the SERS effect by increasing the adhesion area with UA and promoting hotspot creation. The *Asterias forbesi*-inspired SERS (AF-SERS) substrate demonstrated an outstanding SERS effect; this result was further confirmed through electromagnetic field simulations. Using the AF-SERS substrate, UA was detected across a wide range of samples, including serum and human urine. Moreover, the substrate enabled the specific identification and detection of UA among various endogenous metabolites. These findings showcase the promising potential of the developed AF-SERS substrate as a sensor platform not only for UA but also for various targets in real-world field settings.

2. Experimental Section

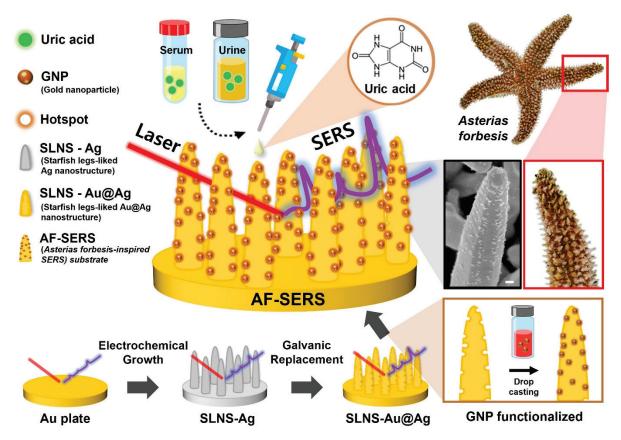
2.1. Chemical Agents

The reagent for SERS substrate fabrication and SERS performance analysis (sodium citrate, chloroauric acid (HAuCl₄), potassium dicyanoargentate (KAg(CN)₂), Rhodamine 6G(R6G)), UA, selective experimental group (L-ascorbic acid, hydrochloride dopamine, L-cysteine, D-(+)-glucose, creatine), and human serum and sigmatrix urine diluent (artificial urine) were purchased from Sigma Aldrich (St. Louis, MO, USA). A glass slide was purchased from Paul Marienfeld GmbH & Co. KG (Lauda-Königshofen, Germany). UA solution was generated by dissolving it in deionized (DI) water, and the selectivity analysis group was prepared in the same manner. All glassware was washed with piranha solution in concentrated sulfuric acid (98% w/v) and hydrogen peroxide (30% w/v) prior to use. All solutions were prepared using Millipore DI water with a resistivity of at least 18.2 M Ω cm at 25 °C. Additionally, all target samples for final substrate production were reacted with GNP at a 1:1 ratio based on volume.

2.2. Uric Acid Detection Strategy Using AF-SERS Substrate

To develop a sensor capable of simultaneously detecting UA deficiencies and excesses outside the normal range in both blood and urine, a SERS structure with a broad surface area is crucial. This type of structure ensures an effective response even at high UA concentrations [26]. Therefore, we developed a detection strategy based on a SERS substrate inspired by *Asterias forbesi*, commonly known as Forbes Sea Star (Scheme 1). Initially, an electrochemical method was employed to produce a silver nanostructure resembling the legs of a starfish. This was achieved by reducing silver using the KAg(CN)₂ solution on

the surface of the gold plate through overpotential [27]. However, various substances present in serum and urine can promote silver oxidation, leading to a decline in the intensity of the SERS signal [28]. To address this, we employed galvanic substitution for reducing gold ions using an HAuCl₄ solution, thereby coating the surface with gold while simultaneously inducing surface roughening. Subsequently, through the integration of gold-coated silver nanostructures with GNPs, a nanostructure resembling the arms of *Asterias forbesi* was successfully fabricated. This was confirmed through SEM images, demonstrating that the developed structure was similar to the spikes in *Asterias forbesi*. These nanostructures exhibited numerous hotspots, a wide surface area facilitated by high Z-axis and particle characteristics, and a bimetallic enhancement effect, which combined the effective plasmonic properties of silver with the chemical stability of gold, ensuring a strong SERS effect [29]. Through these enhancements, the SERS substrate effectively detected UA in serum and urine samples.



Scheme 1. Schematic of Raman spectroscopic detection of UA derived from serum and urine using an *Asterias forbesi*-inspired SERS substrate.

2.3. Synthesis and Optimization of AF-SERS Substrate

Starfish-leg-like Ag nanostructure (SLNS-Ag) and starfish-leg-like Au@Ag nanostructure (SLNS-Au@Ag) were fabricated using the methods described in previous studies [30,31]. Briefly, bare Au plate was fabricated using an E-beam evaporator, and SLNS-Ag was produced through the reduction in KAg(CN)₂ by employing chronoamperometry. Subsequently, 50 μ L of 0.2 mM HAuCl₄ was placed in a polydimethylsiloxane (PDMS) mold for 30 min to induce galvanic replacement, resulting in the fabrication of SLNS-Au@Ag.

GNPs were synthesized using the Turkevich method [32]. For the synthesis of GNP, 1 mL of 1 wt% HAuCl₄ solution and 5 mL of 1 wt% sodium citrate solution were added to 100 mL of DI water and heated at 100 °C for 1 h with vigorous stirring. Afterward, the target substance to be detected was mixed with the synthesized GNP solution in a 1:1 ratio and dropped onto the SLNS-Au@Ag substrate to complete the production of the AF-SERS

substrate. For sensitive and precise measurements, each AF-SERS substrate was exclusively used for a single sample, ensuring that none was reused to maintain the highest standards of accuracy.

2.4. SEM- and FEM-Based Simulation Data

The morphologies of SERS substrates were observed via field-emission scanning electron microscopy (FE-SEM; JSM-IT800, JEOL, Tokyo, Japan). SEM images were obtained at 15 kV HV at different magnifications.

Finite elements method (FEM)-based plasmonic simulations of the optimal SERS substrate analysis were performed using the electromagnetic wave module of COMSOL Multiphysics 6.0. The structures of the nanopillar SERS substrate (Au plate, SLNS-Ag, SLNS-Au@Ag) and nanopillar-combined GNP were based on the SEM images. The time harmonic maxwell equation with boundary conditions was calculated for plasmonic SERS substrates. The simulation was vertically projected onto the surface of the SERS substrate with an incident light of 785 nm. The refractive index and simulation function were referenced from a previous study [26], and RGB values were used to analyze the FEM simulation results using Image J software(Image J v.1.53).

2.5. SERS Analysis

SERS spectra were analyzed using Raman spectroscopy (In Via Reflex, Renishaw, Wotton-under-Edge, UK). The laser was focused using a $100\times$ objective lens (Leica DM2700 M, DEU, and Renishaw Centrus Detector, GBR), and the diameter of the spot of the laser beam was approximately 0.4 μ m. All SERS spectra were measured using a 785 nm laser with a power of 5.05 mW after 1 s exposure and 10 accumulations.

The Raman spectra of R6G were measured in the range of $567-1675 \text{ cm}^{-1}$, whereas those of UA and the selective experimental group were measured in the range of $497-1622 \text{ cm}^{-1}$.

2.6. Optimization of SERS Substrate

To analyze the optimal SERS substrate, 5 μ L droplets of 100 μ M R6G were reacted on the following prepared substrates: Au plates, SLNS-Ag, and SLNS-Au@Ag. Additionally, to evaluate the SERS performance of GNP-functionalized substrates, 5 μ L droplets of R6G and GNP mixed in a 1:1 volume ratio were applied to each substrate (Au plate, SLNS-Ag, and SLNS-Au@Ag).

To optimize the GNP concentration, the synthesized GNPs were precipitated via centrifugation (Labogene 1524, LABOGENE, Lillerød, Denmark) at $6500 \times g$ rpm (2669 g) for 30 min, and the supernatant was removed or more DI water was added to prepare GNPs at various concentrations ($\times 0.2-\times 5$). To optimize the quantity of GNPs, various concentrations of GNP solutions were prepared using the synthesized GNP stock solution ($\times 1$). Diluted GNP solutions were prepared by diluting the original solution fivefold ($\times 0.2$) and twofold ($\times 0.5$) using DI water. Concentrated solutions, obtained through centrifugation with concentrations twofold ($\times 2$) and fivefold ($\times 5$), were prepared for use. The resulting different concentrations of GNPs were mixed with R6G at a 1:1 ratio, and 5 μ L droplets were reacted with SLNS-Au@Ag for subsequent SERS measurements.

Initially, a series of SLNS-Ag and SLNS-Au@Ag were prepared and exposed to a challenging $10\times$ PBS (phosphate-buffered saline, 100 mM) environment. Subsequently, these particles were individually retrieved on specified days and thoroughly washed with ethanol (EtOH) and DI water. Afterward, 5 μL of GNP solution combined with 100 μM R6G in a 1:1 ratio were deposited onto the washed substrates containing SLNS-Ag and SLNS-Au@Ag. Subsequently, the variations in the Raman spectra of R6G on each substrate were measured daily.

2.7. Uric Acid Detection with AF-SERS SERS Sensors

All sample preparations involved mixing the sample with 1X GNP at a 1:1 volume ratio and subsequently applying a 5 μ L droplet onto the GNP substrate, and samples were allowed to dry prior to SERS measurements.

The reproducibility and uniformity analyses of the fabricated AF-SERS substrate were conducted using a 100 μ M UA solution. Reproducibility analysis entailed reacting the UA of the same concentration with six independently prepared substrates to measure the Raman intensity. Uniformity analysis involved measuring the Raman intensity of UA at 50 randomly selected spots on the prepared substrates.

In laboratory conditions, UA detection was performed in the concentration range from 1 nM to 1 mM. The selectivity of the UA sensor was evaluated using various biological fluid materials, including dopamine, L-ascorbic acid, D-(+)-glucose, L-cysteine, and creatine, each at a concentration of 1 mM.

The actual application of the UA sensor involved the use of 10% diluted human serum and artificial urine samples. Various concentrations of UA ranging from 1 μ M to 1 mM were prepared using 10% diluted human serum. Additionally, using artificial urine samples, we prepared samples with UA concentrations within the normal range, representing 300 μ M, as well as samples indicating UA deficiency with a concentration of 30 μ M and excess conditions with a concentration of 3 mM. The Raman measurement conditions were maintained consistent with those used for Raman analysis.

3. Results and Discussion

3.1. Characterization of AF-SERS Substrate

To confirm the characteristics of the AF-SERS substrate, various analyses, including morphology, Raman intensity, and electromagnetic field simulations, were conducted (Figure 1). Initially, SEM was employed for morphological analysis. Figure 1a confirms the successful fabrication of starfish-leg-shaped silver nanostructures (SLNS-Ag) using electrochemical methods. Subsequently, after galvanic replacement, the surface of SLNS-Ag was replaced with gold, resulting in SLNS-Au@Ag. This replacement led to a darker appearance of the substrate surface (Figure S1), and SEM images revealed a slightly roughened surface compared with the smooth SLNS-Ag surface (Figure 1b). This roughening was attributed to the deposition of gold in the areas where silver was displaced. The displacement reaction is described by the following equation [31]:

Anode Reaction: $3Ag \rightarrow 3Ag^+ + 3e^-$

Cathode Reaction: $AuCl4^- + 3e^- \rightarrow Au + 4Cl^-$

Finally, the SLNS-Au@Ag substrate functionalized with GNP assumed a structure similar to Asterias forbesi (Figure 1c,d). Figure S2 further validates the successful attachment of galvanic and gold nanoparticles through EDS analysis, which identifies the components of each Au and Ag. Additionally, the low-magnification SEM images for each condition are shown in Figure S3. Following that, we conducted Raman intensity analysis using R6G, a representative Raman indicator, based on the morphology observed at each step. R6G is a Raman indicator with fluorescent characteristics, and in order to minimize the impact of fluorescence in Raman analysis, we selected a laser with a wavelength of 785 nm [33]. When using the SLNS-Ag substrate as opposed to a gold plate, enhanced Raman signals at the characteristic peak of R6G (1508 cm^{-1}) were observed [34], attributed to the hotspots created by the nanostructures (Figure 1e). Typically, in a Raman substrate utilizing silver, the intrinsic Raman peak of silver appears around 1000 cm^{-1} . However, in Figure 1e, the Raman intensity attributed to the SERS effect of the Raman indicator, R6G, is relatively high, making it challenging to distinguish the Raman peak of silver, which appears at a much lower value. As seen in Figure S4, when the Raman intensity scale is reduced for closer examination, it is confirmed that, unlike the case with the Au plate, a Raman signal around 1000 cm^{-1} is observed in the SLNS-Ag substrate composed of silver [35]. SLNS-Au@Ag exhibited a slight enhancement in Raman intensity compared with SLNS-Ag, attributed to the gold-silver bimetallic effect resulting from the coexistence of gold and silver (Figure 1f) [36]. When evaluating the SERS intensity using the AF-SERS substrate, a significant amplification in the R6G SERS signal was observed (Figure 1g). This outcome might be anticipated solely due to GNP. However, when only GNPs were employed, the performance was inferior compared to when the AF-SERS was employed (Figure S5). These results suggest a synergistic effect between the structures with excellent Z-axis characteristics and nanoparticles. For the analysis of the Raman intensity of each substrate, emphasis was placed on 1508 cm⁻¹, representing the peak value, which manifests the most conspicuous concentration change among the Raman peaks of R6G. Figure 1h presents a bar graph illustrating the average and variance in the Raman intensity at 1508 cm⁻¹, facilitating a quantitative assessment of the Raman intensity for each substrate. The Raman intensity of the AF-SERS substrate surpassed that of the SLNS-Au@Ag substrate by more than 6.2 times, underscoring its outstanding SERS performance. These findings emphasize the highly effective characteristics of Asterias forbesis-inspired nanostructures as superior SERS substrates. Notably, the presence of smaller scale nanoparticles, in contrast to a singular metal nanostructure, leads to an increased reaction surface area with the target and the generation of additional hotspots [37,38].

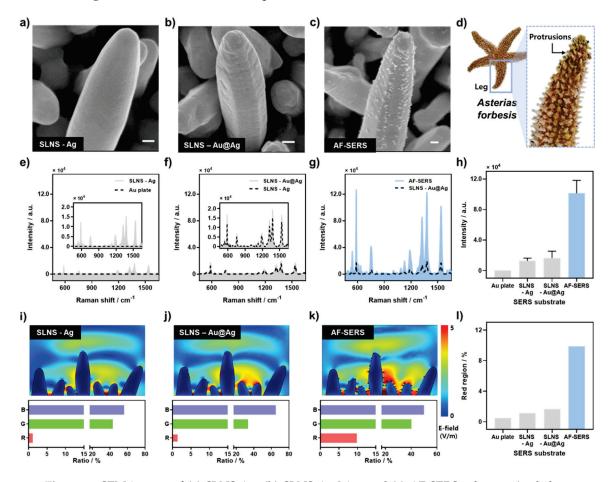


Figure 1. SEM images of (a) SLNS-Ag, (b) SLNS-Au@Ag, and (c) AF-SERS substrate (scale bar: 100 nm). (d) Image of the legs and protrusion structures of the *Asterias forbesi*. Raman spectrum comparison data for 100 μM R6G, a Raman emitter: (e) Bare Au plate and SLNS-Ag substrate (Gray scale: SLNS-Ag, Black dotted line: Au plate), (f) SLNS-Ag and SLNS-Au@Ag substrate (Gray scale: SLNS-Au@Ag, Black dotted line: SLNS-Ag), and (g) SLNS-Au@Ag and AF-SERS substrate (Blue scale: AF-SERS, Black dotted line: SLNS-Au@Ag), respectively. (h) Raman intensities at 1508 cm⁻¹, a specific peak of R6G on various SERS substrates (Au plate, SLNS-Ag, SLNS-Au@Ag, and AF-SERS,

respectively). FEM-based electromagnetic simulation results and RGB value graph for each area of (i) SLNS-Ag, (j) SLNS-Au@Ag, and (k) AF-SERS substrate. (l) Ratio of the red region (in (i-k)) for each substrate.

To validate these results, we conducted electromagnetic field simulations, designing the morphologies to resemble the SEM images of each state (Figure 1i-k). The electron density generated through the simulations is depicted using colors, with red indicating high intensity, green representing intermediate intensity, and blue indicating low intensity. Initially, on the Au plate, hotspot formation was minimal, and the red area was scarce (Figure S6). However, as the nanostructures formed at each stage, the electromagnetic field signal was enhanced, leading to an increased presence of the red region (Figure 1i-k). Although SLNS-Ag and SLNS-Au@Ag appeared similar, slightly higher electron density was observed in SLNS-Au@Ag. Remarkably, the AF-SERS structure exhibited a broader red region, particularly on the GNP surface, indicating enhanced hotspots due to the numerous nanogaps formed by particles. This contributed to the reinforcement of SERS intensity. Quantitative analysis revealed that the red region for AF-SERS was 497.32% larger than that for the SLNS-Au@Ag substrates (Figure 11). Therefore, it can be inferred that AF-SERS demonstrated increased SERS intensity due to the abundant formation of hotspots in both the nanostructures and starfish-leg-shaped structures. Ultimately, through this validation, we confirmed consistency between experimental results and simulation outcomes, providing robust support for our claims.

3.2. Optimization and Sensor Performance of AF-SERS Substrate

The analysis of Raman intensity and simulations underscored the significant impact of the presence or absence of GNPs on the starfish-like nanostructure in determining SERS intensity. Consequently, optimizing the quantity of GNPs corresponding to the protuberances of Asterias forbesi was crucial for achieving optimal sensor performance. To explore this, we conducted an analysis of Raman intensity for R6G by varying the amount of GNP to identify the point of optimal effect. The quantity of GNP was adjusted by considering the amount of GNPs initially included in the synthesized GNP solution as 1, allowing for subsequent dilution or concentration to modify the relative GNP quantity. Detailed methods for this adjustment are described in the Experimental Section. Figure 2a presents the Raman spectrum data corresponding to the amount of GNP. As the amount of GNP increases from SLNS-Au@Ag without GNP, the Raman intensity values show an upward trend. To provide a detailed analysis of the Raman intensity of R6G under each condition, a bar graph was generated to represent the Raman intensity at 1508 cm $^{-1}$ (Figure 2b). However, an excessive amount of GNPs can lead to aggregation and overlap, reducing hotspots and consequently diminishing the SERS effect [39,40]. Therefore, the highest Raman intensity was observed at the x1 concentration, confirming the importance of optimizing the amount of GNP for an effective response. Therefore, for subsequent experiments, the AF-SERS substrate fabricated under x1 conditions was used.

Hence, replacing silver with gold through galvanic substitution is a clever strategy to address the oxidation issue, as gold is more resistant to oxidation. This approach can significantly enhance the stability of the sensor, ensuring long-term consistent and reliable performance [41]. It provides a practical solution to overcome the limitations associated with silver structures, particularly when dealing with biomaterials. To assess the stability of the sensor, it was exposed to stringent environmental conditions (immersion in $10\times$ PBS solution for 7 days), and Raman signals (R6G, $100~\mu\text{M}$) were compared. The experiment involved using the AF-SERS substrate and GNP + SLNS-Ag substrate, where the galvanic replacement step was omitted. The results indicated that the AF-SERS substrate maintained a consistent Raman intensity over the entire 7-day period, with a relative standard deviation (RSD) value of 7.21% (Figures S7 and 2c). The minimal changes observed over the 7-day period under rigorous environmental conditions imply that AF-SERS is also advantageous for long-term storage. In contrast, the control group, GNP + SLNS-Ag substrate, exhibited a noticeable reduction in Raman intensity at the same peak over time, resulting in a significant

54.76% decrease after 7 days (Figures S8 and 2d). The AF-SERS substrate, coated with gold, demonstrated exceptional resistance to oxidation, effectively preventing a decrease in Raman intensity due to oxidation conditions.

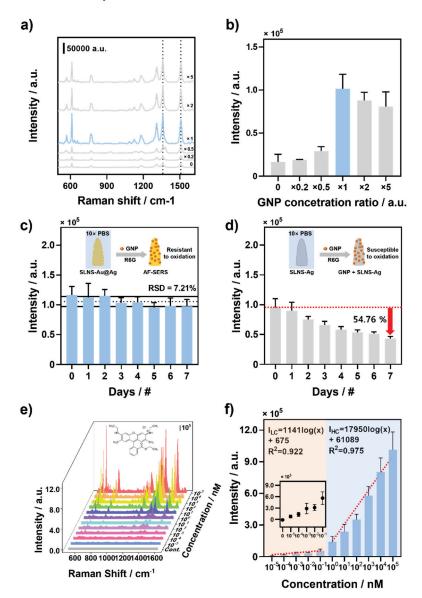


Figure 2. Optimization and SERS performance evaluation of AF-SERS with R6G. (a) SERS spectra of R6G depending on GNP concentration. (b) Measured values of the SERS peak intensities at 1508 cm⁻¹ depending on the GNP concentration. Intensity at specific Raman peaks of R6G in (c) AF-SERS and (d) GNP + SLNS-Ag substrate measured under extreme conditions ($10 \times PBS$ solution) for 7 days. (e) SERS sensitivity of R6G detection at various concentrations (10^{-14} – 10^{-4} M). (f) Measured values of SERS peak intensities at 1508 cm $^{-1}$ depending on the R6G concentrations. Inset depicts the magnified SERS peak intensities of R6G at a low concentration (0– 10^{-10} M).

To evaluate the sensitivity of the AF-SERS substrate, we utilized various concentrations of R6G, ranging from 100 μM to 10 fM. The specific Raman peak of R6G increased with higher concentrations (Figure 2e). For the quantitative analysis of Raman intensity, the Raman intensity at a specific R6G peak (1508 cm $^{-1}$) was expressed using mean and variance (Figure 2f). Distinct linear behaviors were observed in both the low concentration (10 $^{-5}$ –10 $^{-1}$ nM) and high concentration (10 2 –10 5 nM). The linear equations for low and high concentration ranges are I_{LC} = 1141 \times log(X) + 675 (R 2 = 0.922) and I_{HC} = 17,950X + 61,089 (R 2 = 0.975), respectively, where X represents the concentration of R6G. The limit

of detection (LOD) was 1.395 fM, calculated using the formula LOD = $3.3 \times \text{standard}$ deviation/slope. Based on these findings, we computed the analytical enhancement factor (AEF) for the AF-SERS substrate using the following formula [42]:

$$AEF = \frac{I_{SERS}/C_{SERS}}{I_{OR}/C_{OR}}$$

Here, I_{SERS} represents the SERS intensity; I_{OR} denotes the SERS intensity observed on the bare substrate; C_{SERS} stands for the concentration of the Raman indicator on the SERS substrate (i.e., LOD); and C_{OR} is the concentration of the Raman indicator on the bare substrate. The computed analytical enhancement factor (AEF) for the AF-SERS was 3.658×10^{11} , signifying enhancement in performance compared to other SERS substrates (Table S1). Consequently, the AF-SERS substrate demonstrated highly sensitive detection of R6G and successfully detected a wide concentration range. These findings suggest a promising potential of the developed substrate for the detection of UA and other analytes in future applications.

3.3. Performance Evaluation of AF-SERS Substrate with UA

Subsequently, the suitability of the AF-SERS substrate for UA detection was examined, considering the reproducibility between substrates, random spot reproducibility, and UA detection performance (Figure 3). The normal Raman spectrum of uric acid is depicted in the following figure, with band assignment displayed in Table S2. Initially, six different AF-SERS substrates were fabricated, and each substrate was reacted with a consistent 5 μ L of 100 μ M UA (Figure 3a). The specific Raman peak of UA occurred at 640 cm⁻¹, and each substrate exhibited a consistent spectral pattern (Figure 3a). Subsequently, the Raman intensity at 640 cm⁻¹ was compared for the six substrates (Figure 3d), revealing excellent repeatability with an RSD value of 7.227%.

After reacting UA on a single AF-SERS substrate, we presented the measured Raman spectra for 50 random spots as a heat map image (Figure 3b). The bright line at 640 cm⁻¹ in the image indicated the strong SERS intensity of UA on AF-SERS SERS substrates. Similarly, a comparison of Raman intensity at a specific UA peak of 640 cm⁻¹ for each spot revealed excellent uniformity with an RSD value of 13.46% (Figure 3e).

Following the examination of various UA concentrations on the AF-SERS substrate, detection experiments were conducted for 10 concentrations, in the range from 1 mM to 1 nM, covering the normal physiological range of 200–500 μ M (Figure 3c). As shown in Figure 3c, there is a noticeable tendency for the Raman intensity to increase with higher UA concentrations (10^0 – 10^{-5} mM) at a specific Raman peak of 640 cm⁻¹, aligning with the molecular structure of UA. The evaluation of the Raman intensity of UA at the specific Raman peak value of 640 cm⁻¹ revealed a linear increase in the high concentration range, spanning from the normal physiological range (200–500 μ M) to 1 mM (Figure 3f). The linear equation for UA concentration was calculated as I_{HC} = 211.9 X + 31,674 (R^2 = 0.99) in the high concentration range. Additionally, a logarithmic linear increase was observed in the low concentration range (10 nM–100 μ M), and the logarithmic linear equation for low concentration UA was calculated as I_{LC} = 9835 × log(X) – 22,636 (R^2 = 0.987; X represents UA concentration (mM)). The LOD for UA was determined to be 1.16 nM using the following formula: LOD = 3.3 × standard deviation/slope of the regression curved line.

Based on these findings, we experimentally validated the ability of our AF-SERS substrate to detect UA concentrations within the normal range, as well as at deficient and excess concentrations. As indicated in Table S3, our results not only exhibit the highest sensitivity of the substrate in terms of LOD but also boast the broadest detection range compared to various UA detection methods, including electrochemical, fluorescence, and spectroscopic methods.

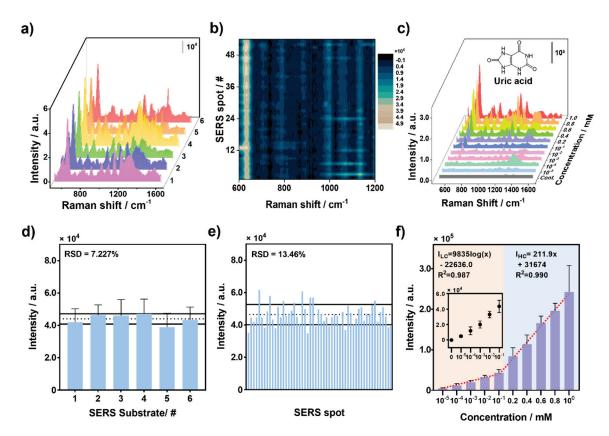


Figure 3. Efficiency and sensitivity analysis of the AF-SERS-based UA sensor. (a) R6G Raman spectra of six different AF-SERS substrates and (b) Heat map image of R6G Raman spectra at 50 random spots. (c) Raman spectra for UA detection at various concentrations on AF-SERS SERS substrate. Inset image is the molecular structure of UA. (d) Raman intensity at 1508 cm⁻¹, which is the specific peak of R6G for each substrate. (e) Raman intensity measurements at 1508 cm⁻¹ for 50 random data points extracted from heat map data. (f) Measured value of the SERS peak intensities at 640 cm⁻¹ depending on the UA concentration. Inset image shows the results in the low concentration range, presented at an optimal scale.

3.4. Detection Selectivity of UA Sensor with AF-SERS

To evaluate the selective UA detection capability of the fabricated AF-SERS substrate, signals from various metabolites were analyzed (Figure 4). Figure 4a presents the chemical formulations of target molecules used in the selectivity experiments, including UA (I), ascorbic acid (II), creatine (III), dopamine (IV), glucose (V), and L-cysteine (VI), which are representative metabolites [43–46]. UA was tested at a concentration of 800 μ M, whereas the remaining substances were experimented with at a high concentration of 1 mM. The Raman spectrum results revealed a distinct and robust intensity at 640 cm⁻¹, the UA-specific Raman peak, exclusively for UA. Other target molecules, even at high concentrations, did not exhibit peaks, which overlapped with UA (Figure 4b). Further analysis focused solely on the Raman intensity at 640 cm $^{-1}$ demonstrated a value of 196,617 \pm 18,102.54 for UA, whereas the remaining molecules displayed significantly lower values, consistently below 15,000 (ascorbic acid: $7559.80 \pm 30,472.24$, creatine: 1307.59 ± 917.42 , dopamine: 917.57 ± 1525.35 , glucose: 8298.72 ± 2247.93 , and L-cysteine: $14,426.45 \pm 4684.40$) (Figure 4c). This rigorous comparison highlights the specificity of the proposed AF-SERS-based UA sensor, demonstrating its ability to selectively detect UA among other metabolites in the body, despite the relatively high concentrations of other target substances.

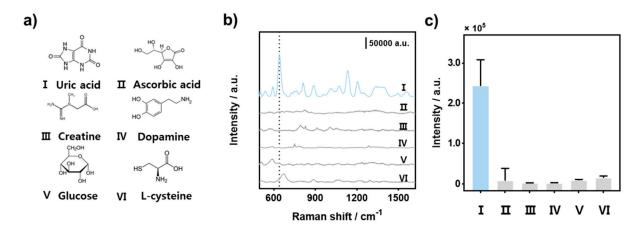


Figure 4. Selectivity analysis of the AF-SERS-based UA sensor. (a) Structural formula of selective groups: uric acid (I), ascorbic acid (II), creatine (III), dopamine (IV), glucose (V), and L-cysteine (VI). (b) SERS spectra and (c) SERS peak intensities at 640 cm⁻¹ of UA and the selective groups.

3.5. UA Detection in Real Sample: Human Serum and Urine

Actual samples for detecting UA include human blood and urine. Therefore, to validate the detection capability of UA within the body using the AF-SERS-substrate-based UA sensor, we designed experiments for UA detection in human serum and real urine samples (Figure 5). First, for UA detection in serum, human serum was diluted at a 1:10 ratio (Figure 5a). Subsequently, UA was mixed into the diluted serum at various concentrations to prepare the samples. Samples were prepared from 1 mM to 1 µM, including a concentration of 200–500 µM, which corresponds to the normal range of UA in human serum. As shown in Figure 5b, the Raman signal at 640 cm⁻¹ changes significantly as the concentration of UA changes. Specifically, the Raman intensity results by concentration at 640 cm^{-1} and specific Raman peak of UA are shown in Figure 5c. A log-linear increase is observed at low concentrations (1–100 μM), corresponding to conditions in which UA deficiency is anticipated. In contrast, a linear increase is observed at high concentrations (200 μM–1 mM), which are high-risk UA concentrations. The log-linear equation for low concentrations was calculated as $I_{LC} = 2825 \times \log(X) + 3132$ ($R^2 = 0.967$), and the linear equation for high concentrations was calculated as $I_{HC} = 35.59X + 3253$ ($R^2 = 0.9644$). Specifically, X denotes the concentration of UA (mM) in serum. The LOD of UA in serum was determined to be 1.05 μM using the LOD formula. The results indicate a decrease in sensitivity compared to the laboratory environment because various proteins in the serum act as interfering factors. However, despite this, the AF-SERS substrate exhibited detection performance, which far exceeded the normal range.

Following this, urine samples were prepared using synthetic human urine without UA (Figure 5d). Uric acid concentrations in the urine of healthy individuals typically range from 160 to 320 μ M [47]. Therefore, we prepared a sample reflecting normal conditions by dissolving 300 μ M of uric acid, within the normal range, in artificial urine. Additionally, urine under UA deficiency conditions was diluted to a concentration of 30 μ M (1/10 of the normal range), and UA excess conditions were produced by adding UA to 3 mM (10 times beyond the normal range). Figure 5e presents the Raman spectra results for each urine sample. Artificial urine contains not only UA but also various proteins, especially urea. Furthermore, the presence of a Raman-specific peak (1000 cm $^{-1}$) of urea was confirmed [48]. Among these, a specific peak related to UA appears at 640 cm $^{-1}$, signifying that it is possible to detect normal UA levels, as well as deficiency and excess states, as shown in Figure 5f. Our UA detection sensor encompasses a broad spectrum within the normal range of uric acid, as illustrated in Figure S9, and in fact boasts a wider coverage compared to conventional detection range [49]. This level of detection capability, which identifies deficiencies and excesses based on UA concentration in urine samples, is also applicable to

real urine samples in the future and is deemed suitable for potential application in future real-world field scenarios.

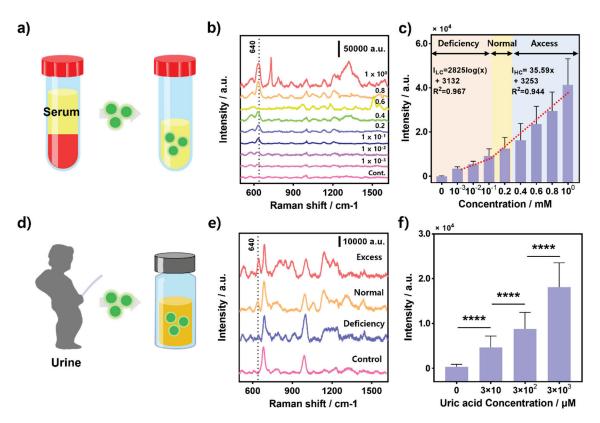


Figure 5. UA detection using the AF-SERS substrate in human serum and artificial urine. (a) Schematic of UA sampling in 10% diluted human serum. (b) SERS spectra and (c) SERS peak intensities at 640 cm⁻¹ of 10% diluted human serum in UA in the concentration range of 1 μ M-1 mM. (d) Schematic of UA sampling in urine. (e) SERS spectra and (f) SERS peak intensities at 640 cm⁻¹ of various UA concentration ranges (deficiency, normal, excess) in urine. Each detection stage can be distinctly differentiated statistically (**** p-value < 0.0001).

4. Conclusions

We developed an AF-SERS substrate for detecting UA by mimicking the ecological structure of a type of starfish. The leg and protrusion morphology of Asterias forbesi exhibited a highly enhanced Raman signal, and simulations confirmed that its structure possessed numerous hotspots. The AF-SERS substrate not only exhibited excellent SERS performance but also minimized oxidation-related damage, making it suitable for practical applications. Moreover, we proved the high repeatability and uniformity of the AF-SERS substrate for UA detection. Using the proposed AF-SERS substrate, we successfully detected UA in a laboratory environment down to a concentration of 1.16 nM and demonstrated selectivity against various metabolites. For real-world application of the sensor, we detected UA in human serum and actual urine samples. In serum, we could differentiate UA concentrations beyond the normal range, covering deficiencies and excesses, with a LOD of 1.05 µM. Similarly, we successfully detected UA in artificial urine samples corresponding to deficiency, normal, and excess states. In conclusion, our research confirms that the proposed AF-SERS substrate enables highly sensitive UA detection in real samples. Our findings suggest its potential as a practical detection platform for real-world UA applications.

Supplementary Materials: The following supporting information can be downloaded at: https: //www.mdpi.com/article/10.3390/bios14010008/s1. Figure S1: Optical images of a (a) SLNS-Ag and (b) SLNS-Au@Ag substrates; Figure S2: SEM-EDS data of AF-SESR substrate; Figure S3: (a) Raman spectrum of 100 μM R6G on only GNP and AF-SERS substrate respectively. (b) Raman intensity graph at 1508 cm^{-1} , the specific Raman peak of R6G. (Scale bar: 500 nm); Figure S4: Raman spectral data of Au plate and SLNS-Ag substrate for R6G 100 μM. Upon reducing the y-axis scale in a specific region, the Raman signal of silver was identified at around 1000 cm⁻¹ (indicated by the red square); Figure S5: (a) Raman spectrum of 100 μ M R6G on only GNP and AF-SERS substrate respectively. (b) Raman intensity graph at 1508cm-1, the specific Raman peak of R6G; Figure S6: FEM-based electromagnetic simulation results for Au plate and RGB ratio spectrum of each area; Figure S7: Raman spectra of R6G for each day on AF-SERS substrate when exposed to harsh conditions (PBS buffer solution) for 7 days; Figure S8: Raman spectra of R6G for each day on GNP + SLNS-Ag substrate when exposed to harsh conditions (PBS buffer solution) for 7 days; Figure S9: A graph comparing the excretion concentration range of uric acid in the body using conventional technology, the normal range of uric acid, and the detection range of uric acid using AF-SERS substrate; Table S1: Comparison of enhancement factor performance of SERS substrates using various nanomaterial [37,44,50–53]; Table S2: Vibrational SERS band assignments for Uric acid [54,55]; Table S3: Comparison of the SERS sensing performances of our AF-SERS and other various sensing platform about UA [25,56–63].

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Article

Complex Spatial Illumination Scheme Optimization of Backscattering Mueller Matrix Polarimetry for Tissue Imaging and Biosensing

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Abstract: Polarization imaging and sensing techniques have shown great potential for biomedical and clinical applications. As a novel optical biosensing technology, Mueller matrix polarimetry can provide abundant microstructural information of tissue samples. However, polarimetric aberrations, which lead to inaccurate characterization of polarization properties, can be induced by uneven biomedical sample surfaces while measuring Mueller matrices with complex spatial illuminations. In this study, we analyze the detailed features of complex spatial illumination-induced aberrations by measuring the backscattering Mueller matrices of experimental phantom and tissue samples. We obtain the aberrations under different spatial illumination schemes in Mueller matrix imaging. Furthermore, we give the corresponding suggestions for selecting appropriate illumination schemes to extract specific polarization properties, and then provide strategies to alleviate polarimetric aberrations by adjusting the incident and detection angles in Mueller matrix imaging. The optimized scheme gives critical criteria for the spatial illumination scheme selection of non-collinear backscattering Mueller matrix measurements, which can be helpful for the further development of quantitative tissue polarimetric imaging and biosensing.

Keywords: polarimetry; Mueller matrix; backscattering imaging; spatial illumination; bio-imaging and sensing

1. Introduction

Polarization techniques have been widely applied to biomedical imaging and sensing with their label-free, non-invasive, and microstructure-sensitive advantages [1–3]. Mueller matrix (MM) polarimetry is more and more prevalently used nowadays, as it encodes numerous polarization properties comprehensively [4–6]. To quantitatively characterize the polarization properties such as diattenuation, retardance, and depolarization, polarization basic parameters (PBPs) are conventionally obtained using Mueller Matrix Polar Decomposition [7], Mueller Matrix Transformation [2], and other MM decomposition methods [8–10]. PBPs can increase the image contrast of specific tissue structures with characteristic properties, such as size, shape, fiber orientation, and alignment, thus enhancing polarimetric biosensing [2]. The primary techniques for MM measurement can be categorized as transmission MM microscopy [11] and backscattering MM polarimetry [12,13]. Thin tissue sample (with a thickness less than 3 × 10¹ μ m) measurements can be taken using transmission microscopy, while bulk tissue sample (with a thickness larger than 3 × 10³ μ m) measurements require backscattering polarimetry, such as bio-structural

and optical properties sensing [14,15], polarimetric endoscopy [16,17], minimally invasive surgery [5], and skin tissue analysis [18–20] for in vivo scenarios. Despite the fact that backscattering MM polarimetry has great potential for various applications, there are still largely unexplored questions. The precision of MM polarimetry can be affected by many factors, including the low signal-to-noise ratio induced by the light field [21], the azimuthal dependence of the data obtained [22], and other measurement factors [2]. One of these is the uncertain incidence and detection angles resulting from uneven biomedical tissue surfaces, namely the surfaces of internal organs with complicated topography [13].

In general, the backscattering MM polarimetric setup consists of a polarization state generator (PSG) and a polarization state analyzer (PSA), which are located on the same side of the sample. The angle formed between the PSG (or PSA) and the normal of the tissue surface is named incidence angle θ (or detection angle ζ), and the angle formed between the PSG and PSA is named absolute spatial angle ψ , which is the sum of θ and ζ . There are four backscattering polarimetry schemes, including the following: (i) normal incidence with normal emergent light; (ii) oblique incidence with normal emergent light; (iii) normal incidence with oblique emergent light; (iv) oblique incidence with oblique emergent light. The complex spatial illumination scheme refers to the spatial relationship between the PSG, PSA, and sample, including the four different schemes above. Ideal collinear reflection MM measurement, as in scheme (i), can reduce the amount of measuring errors by avoiding the orientation effects [23]. However, it requires a non-polarizing beam splitter to ensure the emergent light is colinear with incidence, which increases the cost of the device, induces cumulative errors after frequent use, and asks for an intricate calibration [2]. On the other hand, uneven biomedical tissue surfaces make it extremely difficult to take this way. Conversely, the non-linear reflection approach, although inducing orientation effects, can effectively overcome the above problems [13]. Therefore, non-colinear reflection MM measurement equipment has been increasingly used in biomedical studies and applications.

As an optical method, backscattering polarimetry can be used to detect chemical compounds in tissues or cells by optical signals as part of a biosensor, together with other techniques such as surface plasmon resonance [24,25]. For instance, biosensors based on Mueller matrix measurements have been applied to the detection of miRNA [24], bovine serum albumin [25], dengue virus, and glucose [26-28]. Specifically, a decomposition Mueller matrix polarimetry was proposed for detecting miRNA [24], showing that the polarization parameter could be used as a quantitative sensing index of chemical compounds. A cutting-edge biosensing method that tracks the binding reaction between bovine serum albumin and its antibodies by measuring the phase difference between p- and s-polarization was demonstrated [25]. Moreover, Mueller matrix polarimetric techniques were also used to sense the glucose concentration in aqueous solutions or detect dengue virus [26-28]. However, in the above-mentioned works, we can hardly find a clear criterion of the spatial illumination angle selection, which may have a potential influence on the polarimetric measurement. Both the incidence and detection angles of non-collinear backscattering polarimetric devices are inconsistent in different Mueller matrix polarimetry-based biosensors. For instance, the detection angle was 80° in [24], and the absolute space angle was 60° in [26]. Therefore, we believe that analyzing the impact of complex spatial illumination on backscattering Mueller matrix polarimetry is crucial, which is the focus of this study.

Some recent works have demonstrated that non-colinear reflection can have a complex impact on MM measurement [2]. For instance, the polarimetric aberrations in the MM elements and PBPs at different θ , as in scheme (ii), have been analyzed previously [13]. In this study, we analyze the detailed features of complex spatial illumination-induced aberrations for backscattering MM. We measure the MMs of an anisotropic silk phantom designed previously [13] and porcine liver tissue to explore the influence of the different schemes (ii)—(iv), on MM imaging. The experimental results demonstrate that different polarimetric aberrations can be induced in two-periodic MM elements, as in schemes (ii) and (iii), which are verified using the mean-square error (MSE) and energy spectral density. Furthermore, we find that when measuring MMs with an adjusting distribution of θ and ζ , as in scheme

(iv), the aberrations can be significantly alleviated. Additionally, as ζ increases, obvious image distortions can be observed, performing as transverse compressions parallel to the incident plane. The relationship between ζ and the extent of distortions is analyzed using linear regression (LR), and the MM elements can be reconstructed accordingly. Based on the analysis, this study gives corresponding suggestions for selecting appropriate polarimetric schemes to extract specific polarization properties, which can be helpful for the further development of quantitative tissue polarimetric imaging and biosensing.

2. Materials and Methods

2.1. Experimental Setup and Tissue Samples

In this study, we adopt an experimental setup for 3×3 backscattering MM imaging based on a division of focal plane (DoFP) camera [29–31]. As shown in Figure 1a, the monochromatic light emitted from the LED (3 W, 633 nm, $\Delta\lambda$ = 20 nm, Daheng Optic, Beijing, China) is modulated by the PSG consisting of a collimating objective lens L1 (Hengyang Optic, Guangzhou, China) and three fixed polarizers P1 (extinction ratio > 1000:1, LBTEK Optic, Changsha, China), and then scattered by a tissue sample. The P1 is driven by a screw linear motor M (FSK30, FUYU Technology, Chengdu, China) to generate 0°, 45°, and 90° linear states of polarization illuminations in series during the measurement. The scattered photons from the tissue samples are captured by the DoFP camera (PHX050S-P, Lucid Vision Labs, Richmond, BC, Canada) after passing through the imaging objective L2 (Hengyang Optic, Guangzhou, China). The DoFP camera contains numerous units capable of autonomously extracting intensity information from four different linear polarizers within different emergent lights. As illustrated in Figure 1b, by utilizing the grayscale images, we can construct the sample's MMs and further analyze them by extracting the corresponding PBPs, which characterize the polarization properties of the tissue sample.

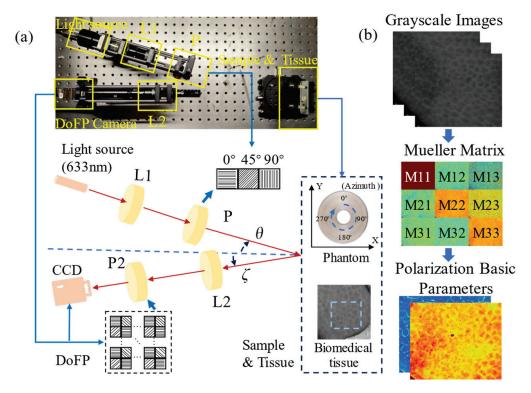


Figure 1. Schematics of the experimental setup and samples: (a) Backscattering MM setup using a DoFP camera. L1 and L2, lenses; P1 and P2, polarizers; θ , the oblique incidence angle; ζ , the oblique detection angle; sample and tissue, concentrically aligned silk phantom, and porcine liver tissue; (b) flowchart of polarization parameter acquisition based on the 3 × 3 MMs.

Here, we adopt an anisotropic scattering phantom consisting of concentrically aligned silk fibers, as shown in Figure 1a, which can generate the MMs of cylindrical scatterers along all azimuthal directions in the imaging X—Y plane in a single measurement [4]. Additionally, we utilize the porcine liver tissue as an ex vivo experimental sample for validation. This animal experimentation work was approved by the Ethics Committee of Tsinghua Shenzhen International Graduate School, Tsinghua University.

Prior to measurement, calibration of each optical component is conducted using a polarimeter (PAX1000, Thorlabs, Newton, NJ, USA) to ensure that systematic errors are maintained within 1%. During measurement, a constant distance between the PSA (or DoFP camera) and the tissue surface is rigorously maintained to minimize the errors arising from focal length variations. Meanwhile, the angles between the center line and both the PSG and sample arms are systematically adjusted to modulate θ and ζ as indicated in Figure 1a. The 3 × 3 MMs of the sample are reconstructed according to Equations (1) and (2), where DoFP, L_2 , P, and L_1 correspond to the 3 × 3 MMs of the optical components, S_{in} represents the input Stokes vector, and S_{out} represents the output Stokes vector [13].

$$S_{out} = DoFP \times L_2 \times MM_S \times P \times L_1 \times S_{in}, \tag{1}$$

$$MM_S = DoFP^{-1} \times S_{out} \times S_{in}^{-1} \times P^{-1}.$$
 (2)

2.2. Azimuthal-Dependent Curves of Mueller Matrix Elements

To comprehensively illustrate the aberrations in MM elements induced by complex spatial illumination, the azimuthal-dependent curves of nine MM elements are constructed according to the method described in our previous work [13]. As shown in Figure 2a,b, we maintain ζ (or θ) at 0° and systematically modulate θ (or ζ) from 0 to 40° in steps of 10° , respectively. Subsequently, the acquired curves are shown in Figure 2d,e. Upon comparing the two groups of curves, we can evidently notice that the curves of M12, M21, M13, and M31 exhibit opposite variations with increasing θ or ζ , while the curves of the other MM elements have identical variations with increasing θ or ζ . Additionally, as depicted in Figure 2f, we measure the azimuthal-dependent curves by simultaneously modifying both θ and ζ , shown in Figure 2c. These curves maintain the identical ψ and illustrate the influence of different distributions of θ and ζ on the aberrations. A detailed analysis will be provided in Section 3.

2.3. Analysis Methods

For quantitative analysis, we define the MMs measured when ψ is 10° and θ or ζ are equal to 0° as the referenced MMs. The MSEs between the curves for the other θ or ζ and the corresponding reference curves are calculated as Equation (3) for the two-periodic elements M12, M21, M13, and M31, respectively. In Equation (3), $C(M_i)$ represents the calculated azimuthal-dependent curve, $C(M_{ref})$ represents the corresponding reference curve, and N represents the point number in the curves. In Figure 2d,e, the M12 and M21 curves exhibit a noticeable shift along the vertical axis. Thus, we calculate the direct current components of energy spectral density for different curves.

$$MSE(C(M_j), C(M_{ref})) = \frac{1}{N} \sum_{i=1}^{N} \left\| C(M_j)_i - C(M_{ref})_i \right\|^2.$$
 (3)

Since the anisotropic silk fiber phantom is circularly shaped, we introduce the ellipticity as a parameter to characterize the image distortions induced by the increasing ζ , as shown in Equation (4), where y and x are the longitudinal and transverse distances of the phantom and r is the radius of the phantom. The analytical relationship between ζ and the ellipticity is also shown in Equation (4).

Ellipticity =
$$\frac{y-x}{y} = \frac{r-r\cos\zeta}{r} = 1-\cos\zeta$$
. (4)

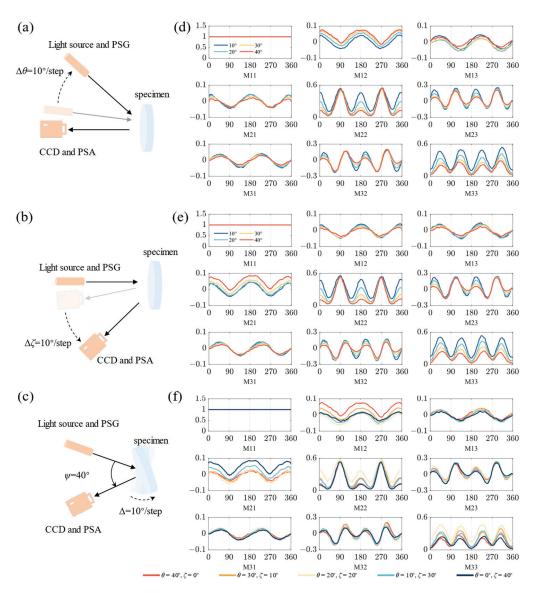


Figure 2. Azimuthal-dependent curves of 3×3 MM elements constructed as schemes (ii)—(iv): (a) incident angle θ modulated; (b) detection angle ζ modulated; (c) absolute spatial angle ψ modulated. The corresponding curves: (d) incident angle θ modulated; (e) detection angle ζ modulated; (f) absolute spatial angle ψ modulated. The horizontal axis shows the azimuth angle of the fibers, and the vertical axis represents the MM element values. All the elements are normalized by M11.

The ellipticity results are analyzed through the LR and least squares method in order to validate the correctness of the analytical relationship. In Equations (5) and (6), x is the deg of θ or ζ , y is the ellipticity of corresponding x, \overline{x} and \overline{y} are the means of x and y, and y is the data number.

$$Variable = \frac{\sum xy - n\overline{xy}}{\sum x^2 - n\overline{x}^2},$$
 (5)

$$Intercept = \overline{y} - Variable \cdot \overline{x}. \tag{6}$$

To validate the significant alleviation of aberrations in PBPs achieved by measuring MMs with an adjusting distribution of θ and ζ , we characterize the depolarization of the tissue sample by measuring the LDoP (linear degree of polarization) and b_2 (linear depolarization and anisotropies) parameters [32], which are sensitive to the oblique incidence [13,33]. They can be calculated as Equations (7)–(9).

$$b = \frac{M22 + M33}{2},\tag{7}$$

$$b_2 = 1 - b, \tag{8}$$

$$LDoP = \frac{M21 + M22}{M11 + M12}. (9)$$

We construct the frequency distribution histograms (FDHs) of the LDoP and b_2 parameters, which can systematically characterize the distribution of the overall polarization properties with PBP images. Several image distance measurements, such as the histogram distance (HD), Euclidean norm, and Structure Similarity Index Measure (SSIM) [34], are utilized to evaluate the alleviation effects quantitatively. Meanwhile, these measurements can also be used to identify the structural information as different schemes. The formulas for these measurements can be described as Equations (10) and (11),

$$HD(x,y) = MSE(histogram(x), histogram(y)),$$
 (10)

$$SSIM(x,y) = \frac{(2\mu_x\mu_y + C_1)(2\sigma_{xy} + C_2)}{(\mu_x^2 + \mu_y^2 + C_1)(\sigma_x^2 + \sigma_y^2 + C_2)}, s.t.C_1 = (0.01 \times L)^2, C_2 = (0.03 \times L)^2, (11)$$

where μ_x , μ_y , σ_x , σ_y , and σ_{xy} are the local means, standard deviations, and cross-covariance for MM images x, y, and L = 1 for MM images.

3. Results

3.1. Aberrations Induced by Complex Spatial Illumination of Two-Periodic MM Elements

It is widely recognized that the D_L derived from M12 and M13 represents the linear diattenuation, while the P_L derived from M21 and M31 represents the linear polarizance of the tissue sample [32]. D_L and P_L can be calculated as Equations (12) and (13).

$$D_L = \sqrt{M12^2 + M13^2},\tag{12}$$

$$P_L = \sqrt{M21^2 + M31^2}. (13)$$

Previously, we demonstrated that the symmetries of the M12 and M21 pairs, together with the M13 and M31 pairs, can be broken when θ is more than 20° [13]. In this work, we measure MMs and construct the azimuthal-dependent curves of the anisotropic silk phantom, as in schemes (ii) and (iii), as shown in Figures 2d and 2e, respectively. The azimuthal-dependent curve values of M12, M21, M13, and M31 are within the range of -0.1 and 0.1 when θ and ζ are less than 20° . As θ increases, the M12 curve progressively shifts towards positive, while the M21 and M31 curves exhibit amplitude variations but maintain both positive and negative values. Similarly, as ζ increases, the M21 curve shifts towards positive, while the M12 and M13 curves have variations in amplitude but still exist as positive and negative values. Furthermore, these can be more clearly seen in the MM images of the phantom in Figure 3a,c. It should be noted that the curves' variations of other MM elements are identical, as in both schemes (ii) and (iii), like period degeneracy in four-periodic elements. These observations imply that the different polarimetric aberrations in two-periodic elements with identical ψ may be directly correlated to different schemes.

In Figure 3b,d, we quantitatively analyze difference aberrations in MM elements by calculating the MSE of curves using schemes (ii) and (iii), respectively. Different ring areas with various radii for each MM element are collected to calculate the means and standard deviations (SDs) of the MSE. The error bars on the plots represent the SDs of the data. In consideration of the image distortions due to the oblique emergent light, the SDs, which do not exceed 5% of the means, validate the reliability of our experimental results. As shown in Figure 3b,d, there is a significant difference between the two conditions, as in schemes (ii) and (iii). Figure 3b shows significant increases in the MSE between the M12 curves with θ exceeding 20° and the referenced M12 curve. Notably, when θ increases to 40°, the MSE even rises to 1.5 × 10⁻³, indicating a significant aberration in M12. Similarly, Figure 3d shows significant increases in the MSE between the M21 curves with ζ exceeding 20° and the referenced M21 curve. Notably, when θ increases to 40°, the MSE even rises

to 2.3×10^{-3} , indicating a significant aberration in M21. Additionally, the curves of M12, M13, M21, and M31 are similar to the trigonometric functions, and the curves of M12 and M21 shift along the vertical axis with an increase in θ or ζ . Therefore, we calculate the direct current components of the energy spectral density at the above eight different angles to quantitatively analyze the aberrations, as shown in Table 1. For most curves, their direct current components tend to be zero, and the variations do not exceed 3×10^{-4} , excluding the curves of M12 or the curves of M21. When θ or ζ attain significant magnitudes, such as achieving 40° , the direct current components of the M12 and M21 curves can escalate to 2×10^{-3} . These results also show that false-positive linear diattenuation is induced as θ increases and false-positive linear polarizance is induced as ζ increases. Through these quantitative analyses, it can be proven that the aberrations are different from schemes (ii) and (iii), which cannot be ignored in the polarization properties analysis.

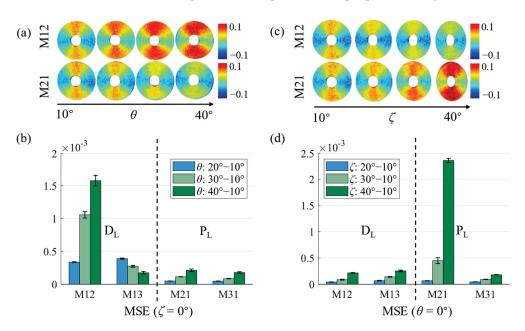


Figure 3. M12 and M21 images and analysis of anisotropic silk phantom under different θ and ζ : (a) M12 and M21 images with θ from 10° to 40°; (b) MSE with different θ ; (c) M12 and M21 images with ζ from 10° to 40°; (d) MSE with different ζ .

Table 1. Direct current components of the energy spectral density with different MM elements.

Degrees	M12 ($\times 10^{-5}$)	M13 ($\times 10^{-5}$)	M21 ($ imes 10^{-5}$)	M31 ($ imes 10^{-5}$)
$\theta = 10^{\circ}$, $\zeta = 0^{\circ}$	0.87	0.38	0.30	0.13
$\theta = 20^{\circ}, \zeta = 0^{\circ}$	40.83	25.66	0.29	0.0080
$\theta = 30^{\circ}$, $\zeta = 0^{\circ}$	119.83	12.61	0.07	0.03
$\theta=40^{\circ}$, $\zeta=0^{\circ}$	173.45	2.49	2.42	0.01
$\theta=0^{\circ}$, $\zeta=10^{\circ}$	0.89	1.04	0.68	0.10
$\theta=0^{\circ}$, $\zeta=20^{\circ}$	3.11	2.98	4.27	0.19
$\theta = 0^{\circ}$, $\zeta = 30^{\circ}$	4.13	0.29	48.64	1.13
$\theta=0^{\circ}$, $\zeta=40^{\circ}$	5.87	3.81	252.78	4.24

In summary, when θ or ζ is beyond 30°, we can obtain more accurate M21 and M31, as in scheme (ii), or obtain more accurate M12 and M13, as in scheme (iii); however, it should be noted that image distortions will be included by the oblique emergent light.

3.2. MM Image Distortions Induced by Oblique Emergent Light

We find that the MM images of the phantom undergo transverse compressions, defining the transverse direction as parallel to the incident plane. The transverse compressions

result in the MM image distortions and the destruction of their structural information. To quantitatively analyze the distortions, ellipticity is utilized as a characterization parameter, which represents the distortions of the MM images and can be obtained by Equation (4). As shown in Figure 4a, it is evident that the variations in θ exert minimal influence on the distortions of the MM images, while the shape of the MM images transforms into an ellipse with an increasing ζ . In Figure 4b, the lengths of the columns with the identical color, representing the identical θ , progressively escalate as ζ increases from 10° to 40° . Notably, this phenomenon is most obvious when θ is 0° or 10° . This intuitively illustrates that the increase in ζ aggravates the extent of distortions, while the varying θ has no obvious influence on it.

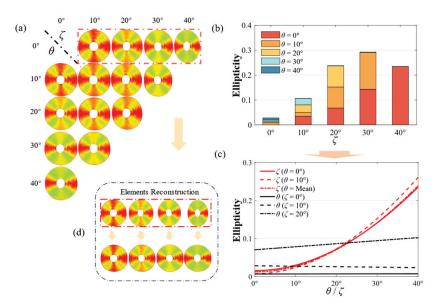


Figure 4. Schematic of ellipticity variation: (a) the M22 elements for different θ and ζ ; (b) histogram of ellipticity for different θ and ζ ; (c) LR fitting curves obtained by the ellipticity data; (d) the reconstruction of the MM elements.

Therefore, according to the experimental setup model, the analytical relationship between ellipticity and ζ is constructed as shown in Equation (4). The analytical relationship illustrates that the cosine value of ζ has a linear relationship with the ellipticity of the MM images. As shown in Figure 4c, we utilize the LR to verify the correctness of the analytical relationship using the above ellipticity data. The fitting curves for the variable θ and the constant ζ are predominantly close to horizontal, indicating a limited association between ellipticity and θ . Additionally, these curves possess different intercepts, indicating that ζ has a predominant influence on ellipticity. Conversely, the curves for the constant θ and the variable ζ are extremely aligned with the curves constructed in Equation (4), affirming the accuracy of the analytical relationship. Moreover, we can introduce a compensation coefficient σ to mitigate the distortions in the transverse direction through Equation (14), where x_{mit} is the distance of the reconstructed MMs and x_{dis} is the distance of the distorted MMs. The reconstructed MM images can be obtained using the bilinear interpolation algorithm. As shown in Figure 4d, this approach effectively mitigates distortions in the MM images.

$$x_{mit} = \sigma x_{dis}, \ \sigma = \frac{1}{1 - Ellipticity}.$$
 (14)

To systematically verify the reliability of our conclusions, the calculated results of LR based on the least squares method are demonstrated in Table 2. When θ is treated as the variable, the slope of the regression equation converges to approximately zero, while the intercept enlarges with an increasing ζ . Conversely, when ζ is considered as the variable, the slope and the intercept are approximately equal to -1 and 1, aligning closely with the

values outlined in Equation (4), respectively. The conspicuous linear correlation between ζ and ellipticity leads to an exceedingly small F below 0.5%, indicating a statistically significant correlation. The exceedingly small p-value, which is also below 0.5%, further reinforces the reliability of the above conclusions. All of these conclusions are applicable to the majority of distortion models caused by oblique emergent light and can be used for MMs' reconstruction, which can recover their structural information. Consequently, we advocate for obtaining MM measurements using scheme (ii) rather than scheme (iii) to acquire undistorted MM images. In instances where measuring MMs using scheme (iii) is deemed indispensable, we suggest restoring the structural integrity of the images using Equation (14).

Table 2. LR results.

Different Degrees	Intercept	X Variable	Significance F	<i>p</i> -Value
$\theta (\zeta = 0^{\circ})$	0.0066	0.000013	0.92	0.21
$\theta \ (\zeta = 10^{\circ})$	0.028	-0.00012	0.80	0.070
$\theta \ (\zeta = 20^{\circ})$	0.070	0.00080	0.30	0.052
$\zeta (\theta = 0^{\circ})$	0.95	-0.93	0.0014	0.0011
$\zeta (\theta = 10^{\circ})$	1.091	-1.08	0.0058	0.0051
ζ (θ = Mean)	0.98	-0.96	0.000054	0.000040

3.3. Calibration of Polarization Properties by Adjusting the Distribution of θ and ζ

It should be noted that MMs measured using schemes (ii) or (iii) with ψ exceeding 20° display significant aberrations, which could potentially induce inaccuracies in characterizing the polarization properties. However, we find that measuring MMs with an adjusting distribution of θ and ζ , as in scheme iv), can remarkably alleviate the aberrations. In Figure 2f, by observing the azimuth dependence curve with θ equal to ζ , some qualitative variations are revealed, as follows: (1) The vertical displacements between the M12, M21, M13, and M31 curves and referenced curves are maintained within 0.03. (2) The M22 curve has a sharp peak at the special azimuth location of 90°. (3) Both the M23 and M32 curves have two sharp peaks at the special azimuth locations of 150° and 210°, respectively. (4) The M33 curve amplitude and minimum approximately increase to 0.4 and 0.2, respectively. These variations make the MMs' curves extremely congruent with the referenced curves. Meanwhile, period degeneracy occurs in the central-block MM elements when θ or ζ exceed 20°, as in schemes (ii) and (iii), which achieves excellent periodic recovery when θ approaches approximate equality to ζ , as in scheme (iv), as shown in Figure 4a. These results reveal that the MMs measured with an adjusting distribution of θ and ζ exhibit a great resemblance to the referenced MMs, which may provide a strategy for obtaining accurate MMs with ψ exceeding 30°.

To further explore the improvement of the polarization properties through this strategy, we measure MMs on porcine liver tissue to validate the alleviation effect on the aberrations in LDoP and b_2 . To avoid the destruction of the structural information, we adopt scheme (ii) to measure MMs when ζ is 0° . As depicted in Figure 5a, subtle structural distortions are observed in the measured MM images when θ and ζ are 20° . Notably, the polarization properties of the microstructure, including LDoP and b_2 , are perceived to be more accurate when using scheme (iv). When θ is 40° and ζ is 0° , the aberrations of the PBPs validate the false-positive and false-negative calculated results of depolarization. It should be noted that b_2 is commonly utilized in cancer detection due to its sensitivity to small particles, so the incidence of the above phenomena needs to be reduced [32]. Furthermore, in Figure 5b, when θ is 40° and ζ is 0° , the curves exhibit substantial peak shifts along the horizontal axis compared to the referenced curves, which can be recovered by measurement with an adjusting distribution of θ and ζ . We calculate HD, where a smaller distance indicates a higher precision in polarization properties, to verify the reliability of the conclusion. The tabulated results reveal that parameters LDoP and b_2 with b_2 and b_3 equal to b_3 are closer

to the corresponding results of the reference parameters, which merely demonstrate 1.4% and 12% HDs in comparison to the parameters measured in scheme (ii). These results illustrate the capability of measuring MMs using scheme (iv) to preserve nearly all accurate depolarization information. Furthermore, measuring polarization properties using scheme (iv) can be more accurate and avoid uncertain aberrations caused by oblique incidence or oblique emergent light in biological tissues.

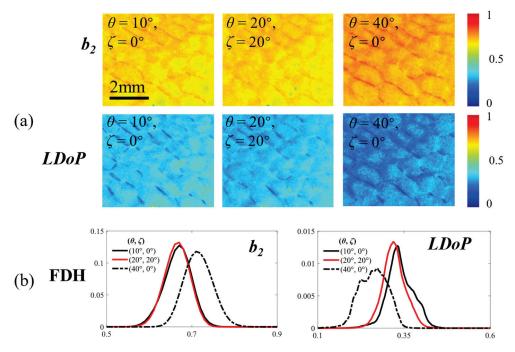


Figure 5. Polarization property recovery: (a) parameter images of LDoP and b_2 at different θ and ζ ; (b) corresponding FDHs of the curves.

To assess the precision of the recovery, we calculate the Euclidean norm and SSIM between the PBP images using scheme (ii) (or scheme (iv)) and the referenced PBP images. In Table 3, employing the Euclidean norm as a metric reveals that the distances between the b_2 and LDoP images using scheme (iv) and the referenced images are only 11.53 and 19.96, respectively, indicating an effective overall recovery. Meanwhile, despite the MM images being distorted, as in scheme (iv), the higher SSIM values of b_2 and LDoP, which imply greater similarity to the original structure, suggest that the precision of the polarization properties in the microstructure can partly compensate for the loss of macroscopic structural information. Therefore, when ψ exceeds 30°, measuring MMs with an adjusting distribution of θ and ζ is more precise if the overall MMs and PBPs are of the utmost concern.

Table 3. Quantitative results of the recovery by different similarity metrics.

Similarity Metrics	$b_2(\theta,\zeta)$		$LDoP(\theta,\zeta)$	
	(20°, 20°)	(40°, 0°)	(20°, 20°)	(40°, 0°)
$HD (\times 10^{-6})$	0.16	10.83	1.29	10.67
Euclidean norm	11.53	49.67	19.96	79.15
SSIM	0.90	0.88	0.81	0.75

4. Discussion

MM polarimetry plays a crucial role in optical biosensors due to its significant performance and advantages in mapping the microscopic morphological information of biological tissues. In our previous work, we focused on the influence of oblique incidence with normal emergent light, as in scheme (ii), and proposed specific optimizing strategies [13]. Here,

we further explore the detailed characteristics of aberrations induced in all non-collinear conditions, as in schemes (ii)–(iv), which are adaptable to more scenarios in practical polarimetric applications. The conclusions of this study for scheme (ii) are consistent with our previous work, while the analysis and optimization strategies of schemes (iii) and (iv) are the essential aims of this work, providing a thorough complement to the previous conclusions. Furthermore, this study serves as a methodological contribution to offer valuable insights into the design and construction of optical systems for backscattering MM polarimetry in bioimaging and sensing applications, such as bio-structural and optical properties sensing [14,15], polarimetric endoscopy [16,17], and skin tissue evaluation [18–20].

5. Conclusions

In this study, we systematically analyzed the polarimetric aberrations induced by three different illumination schemes in backscattering MM imaging. These schemes were distinguished by incidence and emergent light in non-collinear conditions. Additionally, comprehensive comparisons were performed among the MMs obtained under these schemes, considering ψ as 10° , 20° , 30° , and 40° , respectively. We found that measuring MMs with oblique emergent light can induce specific distortion in MM images, hindering the characterization of tissue structural information. Notably, a linear relationship emerged between the cosine value of ζ and the extent of transverse compression, and this could be used to reconstruct the MM images. Additionally, when measuring MMs as in scheme (ii) with θ exceeding 20°, the aberrations of M12 and M13 exhibited increasing intensities. When it came to scheme (iii) with ζ exceeding 20°, M21 and M31 exhibited significantly increasing intensities. However, the variations in other MM elements remained identical. It should be noted that measuring MMs as different schemes may have different effects on the polarization parameters. Furthermore, we found that scheme (iv) could alleviate the polarimetric aberrations. Quantitative analysis of PBPs using porcine liver tissues revealed a significant improvement in the alleviation of PBP aberrations when MMs were measured with an adjusting distribution of θ and ζ . These findings delivered some crucial guidance for choosing the appropriate spatial illumination for non-collinear MM imaging, as follows: (1) If possible, measure MMs with a normal emergent light to obtain accurate structural information. (2) For linear polarizance, measure MMs as in scheme (ii) to accurately obtain M21 and M31. (3) For linear diattenuation, measure MMs as in scheme (iii) to obtain M12 and M13 accurately. (4) To obtain overall MMs and PBPs, measure MMs as in scheme (iv) with an adjusting distribution of θ and ζ . In summary, the optimized schemes provided critical criteria for the spatial illumination scheme selection of non-collinear backscattering MM measurements, which can be helpful for the further development of quantitative tissue polarimetric imaging and biosensing.

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Article

Thermo-Visco-Elastometry of RF-Wave-Heated and Ablated Flesh Tissues Containing Au Nanoparticles

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Abstract: We report non-contact laser-based Brillouin light-scattering (BLS) spectroscopy measurements of the viscoelastic properties of hyperthermally radiofrequency (RF)-heated and ablated bovine liver and chicken flesh tissues with embedded gold nanoparticles (AuNPs). The spatial lateral profile of the local surface temperature in the flesh samples during their hyperthermia was measured through optical backscattering reflectometry (OBR) using Mg–silica-NP-doped sensing fibers distributed with an RF applicator and correlated with viscoelastic variations in heat-affected and ablated tissues. Substantial changes in the tissue stiffness after heating and ablation were directly related to their heat-induced structural modifications. The main proteins responsible for muscle elasticity were denatured and irreversibly aggregated during the RF ablation. At T > 100 °C, the proteins constituting the flesh further shrank and became disorganized, leading to substantial plastic deformation of biotissues. Their uniform destruction with larger thermal lesions and a more viscoelastic network was attained via AuNP-mediated RF hyperthermal ablation. The results demonstrated here pave the way for simultaneous real-time *hybrid optical sensing of viscoelasticity and local temperature* in biotissues during their denaturation and gelation during hyperthermia for future applications that involve mechanical-and thermal-property-controlled theranostics.

Keywords: radiofrequency; heating; hyperthermia; ablation; fiber-optic; sensors; Brillouin spectroscopy; viscoelastic; gold nanoparticles; thermometry; temperature; protein denaturation; flesh; metal nanoparticle; theranostics

1. Introduction

For localized treatment of cancer in the liver, brain, and breast [1,2] and of Barrett's esophagus syndrome, varicose veins, and uterine fibroids, *radiofrequency ablation (RFA)* at mild powers is often employed as a minimally invasive, safe, and low-risk therapeutic procedure [3,4]. During RFA, the electromagnetic (EM) waves passing through biological tissues via an applied needle cause their local heating due to the propagation of an alternating electrical current (AC) modulated at frequencies ranging from 100 kHz to 300 GHz [5], which causes ionic agitation and cellular necrosis, which can ultimately cause tissue coagulation [6]. However, a lack of control of thermal energy deposition during RFA can lead to such undesirable therapeutic results as malignant transformation [7]. Therefore, the control of the RFA procedure is essential for defining how much tissue is destroyed while avoiding charring and excessive damage to surrounding healthy tissues.

Compared to other EM wave sources, such as visible and near-infrared light, RFA causes much deeper heat penetration into tissues and, therefore, is applicable for the treatment of deeply localized solid tumors [8]. However, it is impossible to generate desirable ablation zones that cover spatially extended tumors, and the employment of a large RF

power and a longer time in order to extend the ablation zone will lead to undesirable tissue burns. Adding a mediating absorptive agent, such as *metal nanoparticles* (*NPs*), expands the ablation zone during RFA. Hyperthermal heating occurs in the nanometer-scale vicinity of embedded NPs, thus increasing the local temperatures up to hundreds of degrees on a micrometer scale, leading to a more spatially extended temperature distribution within the heated sample, which causes the disruption of cells in the vicinity. More importantly, loading selected NPs with high thermal and electrical conductivities driven by free electrons into tissues promotes wider and deeper ablation zones due to random spatial distributions of localized heating sources attributed to these metal NPs [9].

Gold nanoparticles (AuNPs) are promising for the targeted thermal destruction of cancer due to their abilities of surface functionalization and hyperthermal RF heating. Moreover, AuNPs are less toxic than other metallic NPs and do not chemically interact with most biological substances [10]. The multifunctionality of AuNPs makes them ideal candidates for theranostic applications [11]. Furthermore, with the increase in NP concentration, the ablation efficiency increases. However, NPs with sizes smaller than 5 nm and larger than 100 nm are not biologically active. Very small NPs are easily cleared by the kidney, while NPs larger than 100 nm are sequestered by the white blood cells that defend the immune system [12].

The viscoelastic properties of tissues vary due to tissue dehydration, protein denaturation, and coagulation during hyperthermal treatments, and these properties make them good candidates for dosimetry [13]. Monitoring *viscoelastic variations* in tissues is critical during thermal ablation, since cell death under hyperthermal treatment is driven by AC-induced thermal stresses in the presence of magnetic nanofluids [14]. It was revealed that thermal conduction and tissue volume relaxation parameters play a significant role in tissue bio-thermo-mechanics [15]. A fully nonlinear thermo-visco-hyperelastic finite-element algorithm was developed to describe the nonlinear characteristics of bioheat transfer in deformed soft tissues with thermal expansion/shrinkage during hyperthermal treatment [16].

Turning to the effects of AuNPs on the elastic properties of tissue, the conjugation of AuNPs to fibrilized collagen, as a soft tissue filler, leads to the formulation of AuNP-collagen constructions and increases the collagen's crosslinking longevity. The mechanical properties of tissues are enhanced as a function of crosslinking [17]. However, to the best of our knowledge, no investigations have been performed on the *thermo-viscoelasticity of tissues during and after AuNP-mediated hyperthermal* treatment. Additionally, the viscoelasticity monitoring of RF-heated tissues is very critical if one wants to understand and control the evolution of the mechanical properties of tissues during and after RF heating and ablation. In addition, it is well known that the mechanical properties of the extracellular matrices of biological tissues and their biopolymer components are essential for normal tissue function, and disturbances in these properties frequently occur in disease [18].

The awareness that diseases ranging from osteoarthritis to atherosclerosis, cancer, and diabetes are linked to alterations in the mechanical properties of the affected tissue has stimulated research on the mechanical properties of entire tissues [18]. Since proteins are major components of flesh (e.g., beef or chicken), it is important to emphasize that the viscoelastic properties of RF-heated and ablated flesh will be dominated by the viscoelasticity of the major RF-heated and ablated proteins constituting these flesh tissues. In terms of the percentage of mass content, chicken flesh actually has more protein than beef flesh [19]. Heat-ablated tissue is stiffer than untreated tissue and is used to measure spatial stiffness variations and to quantify Young's modulus in tissues after thermal treatment [20]. The elastic and viscous components of a gelated myofibrillar protein network started to increase from 30 to 45 °C [21]. The formation of a viscoelastic gel network in ribbonfish meat was observed during heating, and the maximum rate of increase in the storage and loss of elastic moduli was found to be in the temperature range of 56.8 to 63.3 °C [22]. However, the viscoelastic properties of beef muscle showed a slight decrease for T < 55 $^{\circ}$ C, then a sharp increase for $55 \,^{\circ}\text{C} < T < 80 \,^{\circ}\text{C}$, and a saturation at higher temperatures [23]. During the heating of cancerous biological tissues, the cytotoxicity of tumors occurred at 42 °C, while

protein coagulation started at 60 $^{\circ}$ C. The temperature range of 60 to 100 $^{\circ}$ C is known as an effective parameter for obtaining the largest ablation area. However, temperatures above 100 $^{\circ}$ C cause tissue evaporation and carbonization, which have undesirable effects [24].

For the assessment of the elastic and viscous properties of tissues, various techniques have been utilized, such as ultrasound elastography [25], dynamic nanoindentation [20], dynamic rheology [21], and Brillouin light-scattering (BLS) spectroscopy [26]. The latter technique is based on the inelastically scattered laser light from thermally excited GHz acoustic waves in a probed condensed medium. BLS spectroscopy is a unique nondestructive and non-contact technique for probing the viscoelastic properties of various materials, and it was traditionally applied in soft matter and materials science studies [27,28], geosciences [29], biomedical applications [30,31], and the detection of biomarkers [32,33]. More importantly, advances in confocal micro-Brillouin light scattering have enabled the rapid noninvasive monitoring of tissue biomechanics down to the micrometer scale, thus opening the way for live cell imaging [34,35]. Turning to cancer detection, Brillouin shifts in non-regressing and regressing melanomas and in healthy tissue samples were found significantly different, with the healthy tissue being the softest and non-regressing melanoma being the stiffest [36]. The elastic moduli of muscle and rat tail tendon collagen were obtained, and hydrogen bond force constants were estimated with BLS [37,38]. In addition, BLS allowed the observation of the propagation of sound waves along a one-dimensional periodic array of sarcomeres in rabbit psoas muscle myofibrils [39]. A biomechanical contrast of muscular, connective, epithelial, and nervous tissues was presented through high-resolution Brillouin microscopy imaging [26]. A temperature-dependent stiffening and gelation process of albumin from chicken eggs was reported through BLS [40].

As for local optical thermometry, the OBR technique stands as the most modern approach to temperature sensing with fiber-optic sensors (FOS) during thermal ablation. In particular, Mg–silicate-NP-doped fibers have greater backscattering signals than those of standard single-mode fibers, thus allowing the extension of OBR to the real-time measurement of the spatial distribution of temperature over the inner plane of thermal ablation, providing a valuable alternative to thermal imaging with a simpler implementation [41]. Thermal modifications in photodamaged biological samples can be experimentally monitored by chaotic attractors, which depend on optical transmittance [42].

In this work, we report local viscoelastic variations through BLS spectroscopy within chicken muscle and bovine liver tissues with embedded AuNPs after RF heating and ablation. The distributions of Mg–silicate-NP-doped sensing fibers along an RF applicator during ablation allowed for the planar measurement of the local temperature profile and assessment of the ex vivo RF-ablated zone. BLS combined with Mg–silicate-NP-doped sensing fibers allowed us to obtain viscoelastic sensing correlated with the detection of the local temperature of AuNP-mediated hyperthermally RF-heated and ablated flesh tissues. The nature of the high Brillouin peak shift and broadening contrast between normal tissue and the created thermal lesions was studied. In addition, the relationship between the change in viscoelasticity and the obtained thermal dose of the tissues was deeply investigated to better understand the AuNP-assisted hyperthermal RF ablation phenomenon.

2. Materials and Methods

Brillouin spectra were recorded in the 180° -backscattering configuration by using a 6-pass tandem Fabry–Perot Interferometer TFP-2 (Table Stable Ltd., Mettmenstetten, Switzerland) [43]. The free spectral range was set to 25 GHz. Laser light with a 532 nm wavelength from Verdi-G2 (Coherent, Santa Clara, CA, USA) was used with a mild incident beam power kept below 10 mW to prevent sample damage. The laser spot diameter with a $5\times$ microscope objective was 5 μ m. The Brillouin spectral changes were measured at different spatial points on RF-heated and ablated mammal tissues corresponding to the temperatures measured with temperature-sensing fibers.

The ex vivo ablation of fresh chicken and bovine liver tissues was conducted by using an RF/MW Hybrid generator (Leanfa s.r.1.) according to the method in [44], and this was

aligned with the European Union's "Three Rs" principle. Before use, the meat and livers were stabilized to room temperature at 20–21 °C. The RF power for the ablation procedure was set to 60 W, while the frequency value was 450 kHz. To control the temperature change during radiofrequency ablation, four Mg–silicate-NP-doped fibers spliced to single-mode optical fibers (SMFs) were placed in parallel on the y-axis on meat at a 5 mm distance from each other, as depicted in Figure 1. The generator was set to the safe mode, allowing the termination of the ablation procedure when the impedance of 800 Ω was reached. The active electrode (AE), which was in the form of a research-grade single-tip applicator with a 3 mm diameter and a 160 mm length, was inserted between the 2nd and 3rd fibers to deliver the RF waves, while the meat was positioned on a metallic plate that was used as a passive electrode (PE). The positions of both electrodes and sensing fibers were fixed for all of the experiments. A schematic representation of the experimental setup containing the RF generator, Fabry–Perot interferometer, optical backscatter reflectometer (OBR), and the distributed temperature-sensing optical fiber system is shown in Figure 1.

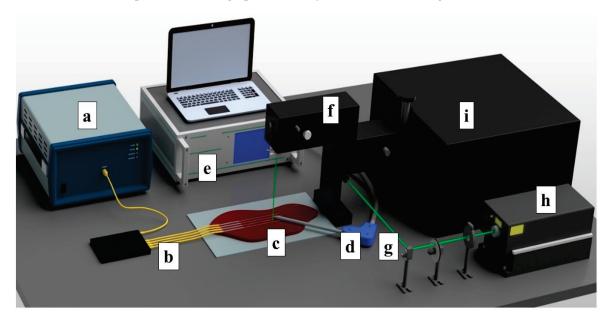


Figure 1. Schematic view of the radiofrequency ablation setup. The setup consisted of: (a) OBR Luna 4600 with a computer used for data acquisition, (b) optical fibers, (c) a bovine liver, (d) the RFA applicator, (e) a hybrid RF/MWA generator used in the RF mode, (f) a confocal microscope, (g) a laser path, (h) a Verdi-G2 532 nm laser source, and (i) a tandem Fabry–Perot interferometer TFP-2.

The synthesis of AuNPs was conducted by using the citrate reduction method, as presented by Turkevic et al. [45]. The average size (d = 20 nm) of the NPs was estimated with a transmission electron microscope (TEM, JEOL JEM—1400 Plus, Indianapolis, IN, USA) image, as shown in Figure 2e. Based on previous studies, Au NPs with a size of >20 nm could take a longer time to be excreted from the body during in vivo experiments, as described in [46]. The optimal density of NPs was selected as 1 mg/mL based on outcomes from previous works [47]. The RF ablation was performed in two conditions: on pristine tissues and on tissues treated with AuNPs with a size of 20 nm dispersed in 0.2% agarose solution. A total of 200 μ L of AuNPs dispersed in an agarose solution were injected into the tissue by using a syringe in proximity to the active electrode. The ablation was performed three times under the same ablation conditions to measure the post-heating Brillouin spectral variations with a coupled confocal microscope.

The temperature-sensing setup was based on a commercial OBR (Luna Inc., OBR4600, Roanoke, VA, USA) that worked in the continuously distributed sensing mode. The OBR instrument was connected to the sensing fibers with different lengths by 1×8 wideband splitter operating in the third optical window. The OBR delivered a signal via a fiber link and then measured the Rayleigh-scattering-induced reflection. Return losses experienced

along the fiber length could be estimated by recording the propagation time of backscattered light. The RFA wavelength of each fiber section's reflection spectrum shifted depending on the change in the temperature. The measurements were taken at 45 mm intervals at the tip of each fiber's spectrum, and these points were grouped into a 2D matrix based on the physical arrangement of the fibers. The obtained 2D thermal maps in the lateral planes of the tissues were arranged in different colors that indicated the temperature changes. The choice of sensing parameters represented a trade-off between the spatial resolution and temperature accuracy; in the experiments, the spatial resolution was set to 2.0 mm, while the overall 'useful' sensing length was set to 45 mm while taking all four fibers into account in this window. All optical connections, including the led-in fibers, splitters, and extenders, were based on a standard single-mode fiber (SMF, corning SMF-28), while the sensing fibers were based on Mg–silicate-NP-doped fibers and had a high spatial resolution (2 mm).

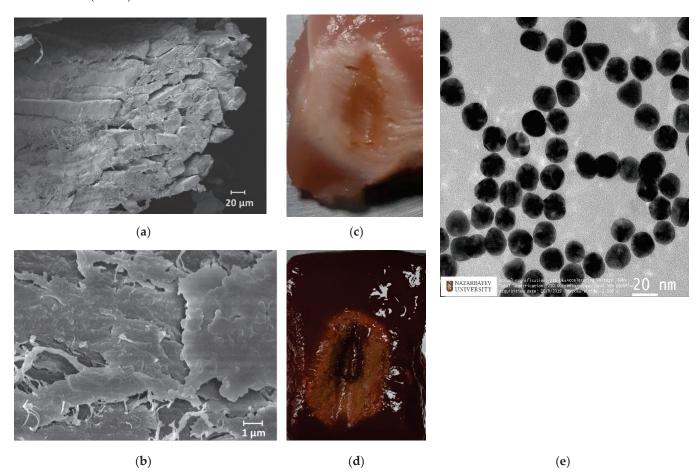


Figure 2. Scanning electron microscopy images of skeletal muscle tissue (a) and a closer view of the connective tissues (b). Ablated chicken muscle (c) and bovine liver (d). Transmission electron microscopy image of AuNPs (e).

3. Results and Discussion

Generally, both chicken and bovine liver flesh consists of approximately 71–75% moisture, 16–17% protein, 5–8% fat, and 1% ash [19]. Chicken muscle tissue contains muscle fibers and connective tissue, where connective tissue covers the muscle fibers, as can be seen in Figure 2a,b. The *myofibrillar proteins* constitute 50 to 55% of the chicken muscle's total protein content, while *sarcoplasmic proteins* account for approximately 30–34%. The remaining 10–15% of the proteins are *connective tissue proteins*. Myosin and actin are the most important proteins that constitute the myofibrillar structure [48]. The bovine liver is composed of smaller histological structures called lobules, which are roughly hexagonal

in shape. In addition, the bovine liver is covered with a connective tissue capsule, which is a thin but tough fibrous supporting connective framework of densely interwoven collagen fibers [49]. The fibers of connective tissue consist of collagen, elastic, and reticular fibers; they are, thus, responsible for the liver's elasticity and tensibility [50].

The heating process had a considerable effect on the biological tissues' properties, as the changes in toughness with denaturation temperatures or indirect interpretations of structural changes in tissue components occurred [51]. Figure 3 demonstrates the chicken muscle's 2D thermal map in the x-y plane, which was processed from the data obtained along the lengths of all four Mg-silicate-NP-doped sensing fibers. The thermal maps of the tissues present the changes in the temperature during the thermal ablation procedure due to the dissipation of RF power in the bovine liver and chicken muscle tissues. The active electrode was positioned at 3.5 cm from the left side, while the sensing system was placed on the right side. The heat pattern clearly showed the maximum temperature increase at the tip of the active electrode and the heat distribution within the tissue, which are depicted in different colors. Images of the RF-ablated bovine liver and chicken meat are shown in Figure 2c,d. Adding AuNPs to the tissues led to a more extensive spatial temperature distribution with a lower maximum temperature; thus, much more uniform temperature distributions along the tissue could be obtained with AuNPs. Hyperthermal RF heating occurred, whereas the AuNPs absorbed the RF waves and produced heat via the limited movement of free electrons inside the NPs. The asymmetric heat pattern was likely due to the heterogeneous properties of the tissue, which deviated from the RF ablation pattern, particularly for fast ablation phenomena.

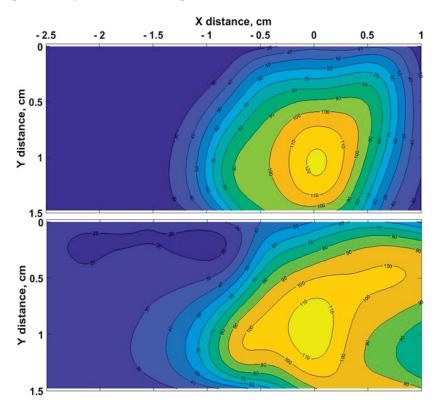


Figure 3. Typical two - dimensional thermal map of pristine chicken muscle (top) and tissues with AuNPs added (down) in the xy lateral plane (x = direction parallel to the RFA applicator and to the sensing fibers).

Figure 4 demonstrates the quantification of the (a) cytotoxic (temperatures > 42 °C) and (b) thermally damaged (temperatures > 60 °C) regions of the chicken muscle tissue for each experimental condition. As can be seen from the bar charts, the AuNP-mediated RFA achieved a larger area of thermal heating and ablation results than it provided in the

Pristine

(a)

400 Mean I Max-m 400 350 350 300 E 300 <u>5</u> 250 200 ç 200 150 100 100 50 50

pristine chicken meat. The same trend was observed for the bovine liver, and both tissues had similar thermal properties.

Figure 4. Evaluation of the surface area exposed to temperatures exceeding (**a**) 42 $^{\circ}$ C and (**b**) 60 $^{\circ}$ C at the maximum ablation temperature; the bar charts show the maximum, minimum, and mean values of the areas over three experiments for each RFA condition.

(b)

Generally, chicken meat was found to enhance its stiffness upon heating in two phases. The first phase took place at T=40–60 °C, which was likely due to the heat-induced denaturation of myofibrillar proteins, especially myosin. The second stage led to further stiffening at T=65–80 °C, which could be ascribed to the denaturation of intramuscular collagen, which was associated with the initial breakage of hydrogen bonds, thus loosening the fibrillar structure with the subsequent contraction and dissolution of the collagen molecules, as well as their eventual gelation [52]. Most of the water in living muscle is held within the myofibrils, and longitudinal shrinkage of the myofibrils at high temperatures causes great water loss. Other studies suggested that the *thermal injury threshold* for RFA in chicken tissue is 65 °C [53]. After thermal treatment at temperatures in the range from 53 to 63 °C, the bovine liver's collagen and reticular fibers were denatured, and this led to the breaking of hydrogen bonds and an irreversible transformation of the crosslinked triple-helical structure into a more random coiled structure [54]. At T>63 °C, the cells had shrunk as much as they could, and most of the water in the cells was forced out [55].

Figure 5a demonstrates the typical Brillouin spectra of pristine and AuNPs-treated chicken muscle tissue corresponding to the measured local temperatures. The back-scattering configuration employed in our BLS measurements enabled us to measure only the bulk longitudinal acoustic waves. The Brillouin peaks decayed in intensity with the increase in the RF-induced temperature due to the heat-induced removal of water content from the tissues, thus effectively reducing the elasto-optic coupling. The ablation temperatures correlated well with the Brillouin shift and linewidth changes, with a frequency contrast from 8 and 1.5 GHz, respectively, at the outer transition zone's boundary and of 14.3 and 9 GHz at the condensation boundary. The rate of the Brillouin shift and linewidth initially increased sharply and changed slightly at the condensation boundary. At high temperatures from 65 to 100 °C, low mechanical contrast was observed, as most of the cells were destroyed due to coagulative necrosis taking place in this zone. Similar results were obtained by using shear wave elastography methods [48] in which the velocities and attenuation of propagated shear acoustic waves were measured [56].

Figure 5b illustrates the temperature-dependent Brillouin shift (a) and linewidth (b) of pristine and AuNP-treated bovine liver and chicken muscle tissues. Initially, at T = 20–30 °C, the Brillouin shift and linewidth stayed the same, i.e., the tissues did not undergo denaturation, and the mechanical changes were typically reversible for this temperature elevation range. However, at $T \ge 40$ °C, a slight increase in the Brillouin shift was noticed, which corresponded to the cytotoxic region. A sharp increase in stiffness was observed with further temperature elevation until the onset of ablation. Once ablation occurred, almost

constant values of the Brillouin frequency were registered at $T > 60\,^{\circ}$ C. The viscous property showed a generally increasing trend as a function of the temperature for both tissues before and after adding AuNPs, as evidenced by the variations in the Brillouin width.

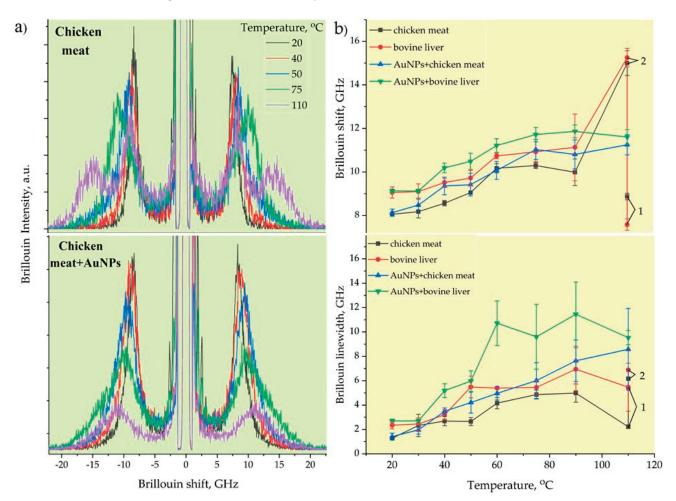


Figure 5. Typical Brillouin spectra of RF - ablated chicken muscle tissues without and with AuNPs (a). The Brillouin frequencies and linewidths of chicken meat and bovine liver with and without AuNPs (b).

At the initial temperatures, the differences in the Brillouin frequency and linewidth between the muscle and liver tissues were almost 1 GHz. The bovine liver was stiffer and more viscous than the chicken muscle tissue. However, the heating of the pristine tissues up to 110 °C for 40-50 s led to incomplete cell destruction, which was observed in the split Brillouin peaks. The position of the first peak corresponded to the tissue's initial Brillouin frequency value and was related to the non-damaged tissues, while the second peak shifted to a much higher Brillouin frequency of 15 GHz, corresponding to the likely formation of an organic polymeric phase related to completely burned (dead) tissues. Furthermore, the AuNP-mediated hyperthermal RF ablation even enabled the uniform destruction of tissues at the micrometer scale, where no Brillouin peak splitting took place. The AuNPs with a size of 20 nm absorbed RF energy and quickly released heat into the surrounding region due to the increased electron-surface scattering, since the size of the NPs was significantly smaller than the mean free path of the electrons in gold. As can be seen from Figure 5b, the tissues with AuNPs were stiffer and more viscous compared to the pristine ones. The implementation of metallic AuNPs in soft tissue led to the formation of an NP-biotissue composite medium, and further heating of this metal-organic medium caused the formation of a more viscoelastic gel network with NPs inside through protein denaturation.

4. Conclusions

We investigated viscoelastic properties and correlated them with the temperature profiles of AuNP-assisted hyperthermally RF-heated and ablated tissues through confocal Brillouin micro-spectroscopy combined with Mg-silicate-NP-doped temperature-sensing fibers. It was revealed that the significant spectral changes in the Brillouin peak position and linewidth were directly related to the obtained thermal dose and the consequent protein denaturation processes. Increased thermal lesion areas with uniform and complete cell destruction were obtained via the loading of Au NPs into tissues. In addition, compared to pristine tissues without any NPs, further heating of the NP-tissue composite medium led to an increase in the composite tissues' viscoelastic properties at the same temperatures. These results suggest that our hybrid Brillouin–OBR technique can possibly define a tissue's structural deformations during thermal therapies, thus encouraging the use of this technique for the dosimetric control of hyperthermia in future applications. This study opens an avenue for real-time simultaneous localized monitoring of the viscoelastic and thermal properties of metal-nanoparticle-embedded bio-tissues across their heat-driven structural phase transitions in hyperthermal and theranostic applications.

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Article

Ternary Heterojunction Graphitic Carbon Nitride/Cupric Sulfide/Titanium Dioxide Photoelectrochemical Sensor for Sesamol Quantification and Antioxidant Synergism

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Abstract: Sesamol (SM) is a potent natural antioxidant that can quench free radicals and modulate the cholinergic system in the brain, thereby ameliorating memory and cognitive impairment in Alzheimer's disease patients. Moreover, the total antioxidant capacity can be amplified by synergistic interactions between different antioxidants. Here, we constructed a ternary heterojunction graphitic carbon nitride/cupric sulfide/titanium dioxide (g-C₃N₄/CuS/TiO₂) photoelectrochemical (PEC) sensor for the quantification of SM and its synergistic interactions with other antioxidants. Crucially, the Schottky barrier in ternary semiconductors considerably enhances electron transfer. The PEC sensor showed a wide linear range for SM detection, ranging from 2 to 1277 μ mol L⁻¹, and had a limit of detection of 1.8 μ mol L⁻¹. Remarkably, this sensing platform could evaluate the synergism between SM and five typical lipid-soluble antioxidants: tert-butyl hydroquinone, vitamin E, butyl hydroxyanisole, propyl gallate, and butylated hydroxytoluene. Owing to its low redox potential, SM could reduce antioxidant radicals and promote their regeneration, which increased the overall antioxidant performance. The g-C₃N₄/CuS/TiO₂ PEC sensor exhibited high sensitivity, satisfactory selectivity, and stability, and was successfully applied for SM determination in both soybean and peanut oils. The findings of this study provide guidance for the development of nutritional foods, nutrition analysis, and the treatment of diseases caused by free radicals.

Keywords: photoelectrochemical sensor; sesamol; antioxidant synergism

1. Introduction

Many people live unhealthy lifestyles, which often include smoking, drinking, getting insufficient sleep, and having an unbalanced diet, all of which result in the accumulation of a large number of oxidizing substances (reactive oxygen and nitrogen free radicals) in the body. These substances disrupt the redox balance of the body and damage biological macromolecules, organelles, and tissue [1]. This process is called oxidative stress [2]. Antioxidants are crucial for slowing down or preventing harmful oxidation and function by scavenging excess free radicals to combat oxidative stress and maintain redox homeostasis. Natural foods, such as vegetables, fruits, and tea, are rich in antioxidants, and their consumption can lower the risk of some chronic diseases [3]. Thus, there is interest in identifying natural sources of antioxidants and adding them to food formulations to improve human health and treat diseases [4]. Some natural antioxidants, such as vitamins E

and C and natural polyphenols (catechins, resveratrol, and tea polyphenols), have received considerable attention from researchers because of their remarkable antioxidant ability and bioavailability [5]. Similarly, sesamol (SM) is a natural oil-soluble antioxidant extracted from sesame oil that has potent anti-aging, anti-mutagenic, and anti-cancer properties. SM is widely used in medicine and food technology [6,7]. For instance, sesame oil has shown potential in treating conditions such as cancer, neurodegenerative diseases, and acute liver injury. This is attributed to its ability to enhance the blood lipid profile and reduce oxidative stress. Moreover, SM has been found to influence the choline system, specifically by amplifying the activity of choline acetyltransferase and reducing that of cholinesterase. Such modulation benefits memory and cognitive functions, which are important in the prevention of Alzheimer's disease (AD) [8]. However, few studies have focused on the determination of the antioxidant capacity of SM.

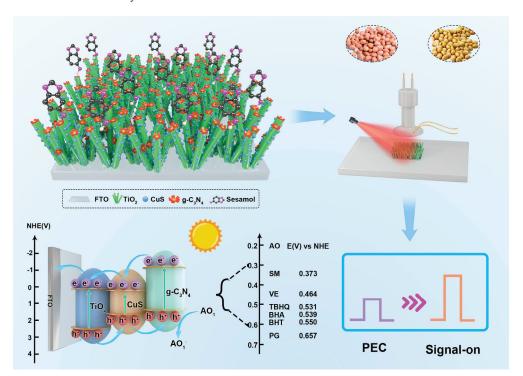
In view of the importance of SM in the fields of medicine and food technology [9], several detection and quantification methods, including electrochemical methods [10], thermogravimetric analysis [11], and spectrophotometry [12], have been developed. However, these technologies have disadvantages, such as low stability and reproducibility, time-consuming operating procedures, low detection limits, and susceptibility to interference, which hinder their large-scale application. Sensors have recently attracted considerable research attention owing to their high sensitivity, low cost, and rapid response [13]. Photoelectrochemistry (PEC) is an innovative technique that combines the advantages of photochemistry and electrochemistry [14]. However, to date, a PEC sensor for the detection of systemic antioxidants based on SM has not been reported [15]. Crucially, using simple and portable equipment, PEC sensors can quantify the activity of antioxidants and the synergistic effect between different antioxidants in an intuitive and rapid way. As a result, PEC sensors are widely used in the analysis of food nutrients and the screening of active ingredients [16].

The excellent performance of PEC sensormostly depends on the properties of the photoactive material [17]. However, single semiconductors have high electron–hole recombination rates and low photocurrent signals as a result of limited visible-light absorption. To enhance photoelectric performance and overcome the limitations of semiconducting photosensitive materials, the development of multicomponent semiconductors and the design of multiple heterojunction structures have drawn attention [18,19]. In particular, the formation of a heterojunction can effectively improve the optical properties of semiconductors and promote electron–hole separation and transfer, thus expanding the light-absorbing window to the visible region and extending the operating life of the material [20].

Recently, nanorods having a large aspect ratio have attracted considerable attention because of their unique optical structure and electronic properties. Among the nanorod-forming materials, titanium dioxide (TiO_2) exhibits a direct channel for electronic transmission and has a large specific surface area, making it especially suitable for various photocatalytic and sensing applications. However, TiO_2 can only be excited by ultraviolet light, which comprises only 4–6% of the solar spectrum. Therefore, the modification of TiO_2 to narrow the band gap is necessary to expand the absorption range to the visible region [21]. Copper sulfide (CuS) is a nontoxic semiconductor material that has a narrow band gap (1.2–2.0 eV) and high visible-light absorption coefficient, exhibiting particularly high absorption in the red and near-infrared regions [22]. Further, as a non-metallic semiconductor polymer, graphitic-phase carbon nitride (g-C₃N₄) has a large number of reactive sites and a unique band gap (2.7 eV). In addition, it is easily functionalized and doped with other semiconductors [23].

To the best of our knowledge, this is the first study to construct a ternary heterojunction $g-C_3N_4/CuS/TiO_2$ for the determination of SM. Briefly, a feasible and sensitive PEC sensing platform was constructed by depositing nanoparticles of CuS on TiO_2 nanorods, followed by loading with $g-C_3N_4$ nanosheets that serve as the photoactive material. The sensing material was then applied for SM detection (Scheme 1). Further, the synergistic effects of SM and other antioxidants, including *tert*-butyl hydroquinone (TBHQ), vitamin E (VE), butyl

hydroxyanisole (BHA), propyl gallate (PG), and butylated hydroxytoluene (BHT), were investigated and confirmed using the ternary heterojunction $g-C_3N_4/CuS/TiO_2$ sensor. The findings of this study will aid the development of sensors for food safety, quality control, and food component compatibility testing and will aid the discovery of treatments for diseases caused by free radicals.



Scheme 1. Schematic of antioxidant capacity analysis based on the g-C₃N₄/CuS/TiO₂ PEC sensing platform. AO, antioxidant; FTO, fluorine-doped tin oxide; SM, sesamol; VE, vitamin E; TBHQ, *tert*-butyl hydroquinone; BHA, butyl hydroxyanisole; BHT, butylated hydroxytoluene; PG, propyl gallate.

2. Experimental Procedure

2.1. Synthesis of g-C₃N₄/CuS/TiO₂ Composites

TiO₂ nanorods were synthesized using the classic hydrothermal method. Briefly, hydrochloric acid (15 mL), ultrapure water (15 mL), and titanium butoxide (500 μL) were added sequentially to a beaker and stirred for a few minutes. The mixture was transferred to a Teflon-lined autoclave containing a conductive fluorine-doped tin oxide (FTO) electrode. After react on at 150 °C for 5 h, the FTO electrode with TiO₂ nanorods was rinsed with deionized water and annealed at 500 °C for 2 h at a heating rate of 2 °C min⁻¹. Then, copper sulfate solution (0.25 μL, 0.1 mol L⁻¹) and sodium thiosulfate solution (4.75 mL, 0.1 mol L⁻¹) were added to a Teflon-lined autoclave containing an electrode modified with TiO₂ nanorods and reacted at 120 °C for 5 h to obtain a sample, denoted CuS/TiO₂. Using the same conditions, different doping amounts of CuS can be obtained by adding different proportions of copper sulfate and sodium thiosulfate solutions.

G- C_3N_4 was prepared by heating melamine at 550 °C for 2 h at a rate of 2 °C min⁻¹. The products were washed three times with ultrapure water and ethanol sequentially, and dried at 60 °C. Then, g- C_3N_4 (10 mg) was added to ultrapure water (20 mL) and ultrasonicated for 3 h. The suspension was then centrifuged at a low speed for 10 min, and the supernatant was collected for further use.

To obtain the g-C₃N₄/CuS/TiO₂ composite, the supernatant of g-C₃N₄ was added to a watch glass containing the CuS/TiO₂ FTO electrode and soaked for 2 h at 70 °C. Then, the g-C₃N₄/CuS/TiO₂ electrode was dried with nitrogen and calcined at 550 °C for 2 h under nitrogen atmosphere at a heating rate of 2 °C min⁻¹. Under the same reaction conditions

as those used to synthesize the composite materials, different loading amounts of g- C_3N_4 were obtained by adding different masses of g- C_3N_4 powder.

2.2. Antioxidant Capacity and Synergistic Effects

First, an antioxidant standard solution having a concentration of 10 mmol L $^{-1}$ was prepared by dissolving SM, GA, CHA, EGC, EGCG, and PC in ethanol. For SM detection, the g-C₃N₄/CuS/TiO₂ electrode was fastened to the PEC cell and suspended in an electrolyte comprising a phosphate-buffered saline (PBS) solution (3.0 mL; 0.1 mol L $^{-1}$) of SM, and the photocurrent was measured at least three times at room temperature. The photocurrent was obtained using the formula: $\Delta I = I_{sample} - I_{blank}$, where I_{sample} is the photocurrent in the presence of the target and I_{blank} is the photocurrent in the absence of the target. To investigate the synergistic effect of the SM-based antioxidant systems, 10 mmol L $^{-1}$ mixtures of SM+VE, SM+BHA, SM+BHT, SM+TBHQ, SM+PG, SM+BHA+BHT+TBHQ+PG, and SM+BHA+BHT+TBHQ+PG+VE were prepared. The molar quantities of all antioxidants in these mixtures were equal. Similarly, VE, BHA, BHT, TBHQ, SM+VE, SM+BHA, SM+BHT, SM+TBHQ, SM+PG, SM+BHA+BHT+TBHQ+PG+VE were added to the detection cell, and the photocurrent was measured. The synergistic effects can be calculated and quantified by comparing the photocurrents of a mixture of antioxidants.

3. Results and Discussion

3.1. Characterization of g-C₃N₄/CuS/TiO₂ Composites

Figure 1A shows a schematic of the g-C₃N₄/CuS/TiO₂ composite synthesized in three steps. Scanning electron microscopy (SEM) images revealed that TiO₂ grew uniformly on FTO and had a smooth surface (Figure S1A in Supplementary Materials). However, the surface became rough after the deposition of the CuS nanoparticles (Figure S1B). Figure 1B shows that the TiO₂ nanorod structure did not change after loading with CuS and g-C₃N₄. As shown in Figure 1C, the tops of the nanorods have a flower-like appearance, which should increase the contact area available for reaction and is conducive to the collection and utilization of visible light. The transmission electron microscopy (TEM) image shown in Figure 1D was used to calculate the cross-sectional width of the g-C₃N₄/CuS/TiO₂ composites: approximately 127 nm, whereas the length of the nanorods was approximately 500–900 nm. These dimensions should promote the longitudinal migration of photogenerated carriers. As shown in Figure 1E, the CuS nanoparticles were uniformly doped on the surfaces of the TiO₂ nanorods, and the figure shows that the particle size of CuS was less than 5 nm. Lattice fringes having interplanar spaces of 0.268, 0.338, and 0.264 nm can be clearly seen in Figure 1F, and these can be ascribed to the (200) and (204) crystal planes of anatase TiO₂ and (110) crystal plane of CuS, respectively. Energy-dispersive X-ray spectroscopy (Figure 1G) measurements revealed the homogeneous distributions of Ti, O, Cu, S, C, and N in the g- C_3N_4 /CuS/TiO₂ composite, suggesting that the composite had been successfully synthesized.

The crystallinity of g-C₃N₄/CuS/TiO₂ was determined through X-ray diffraction (XRD) analysis (Figure 2A). The XRD pattern of g-C₃N₄/CuS/TiO₂ shows several characteristic peaks, corresponding to the (101), (103), (204), and (220) crystal planes of TiO₂; (002) plane of g-C₃N₄; and (102) plane of CuS [24,25]. Furthermore, the XRD pattern of the as-prepared g-C₃N₄/CuS/TiO₂ matches the patterns of standard hexagonal CuS (JCPDS No. 01-076-1725), anatase TiO₂ (JCPDS No. 21-1272), and g-C₃N₄ (JCPDS No. 87-1526) well, suggesting that the stepwise synthesized composites were free of impurities. Next, X-ray photoelectron spectroscopy (XPS) measurements of g-C₃N₄/CuS/TiO₂ were conducted to analyze the electronic states of the surface elements. As shown in Figure S2, the high-resolution (HR)-XPS spectra of the elements ing-C₃N₄/CuS/TiO₂ confirmed that the g-C₃N₄ thin-layer and CuS-nanoparticle-loaded TiO₂ nanorod composite had been successfully synthesized via our designed route [26–28].

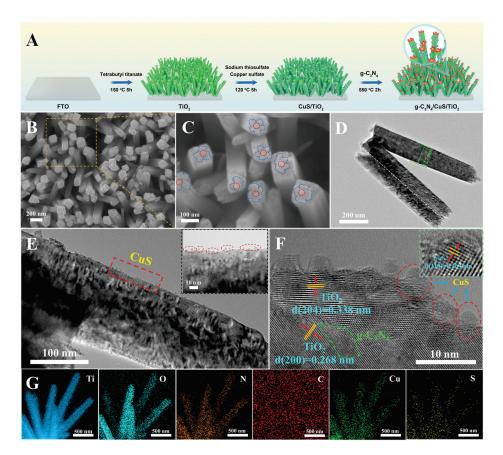


Figure 1. (**A**) Schematic showing the preparation of $g-C_3N_4/CuS/TiO_2$. (**B**,**C**) Scanning electron microscopy (SEM) images, (**D**,**E**) transmission electron microscopy (TEM) images (inset: corresponding magnified picture of CuS nanoparticles), (**F**) high-resolution transmission electron microscopy (HR-TEM) images (inset: corresponding magnified HR-TEM image of CuS nanoparticles) and (**G**) energy-dispersive X-ray spectroscopy (EDX) of $g-C_3N_4/CuS/TiO_2$.

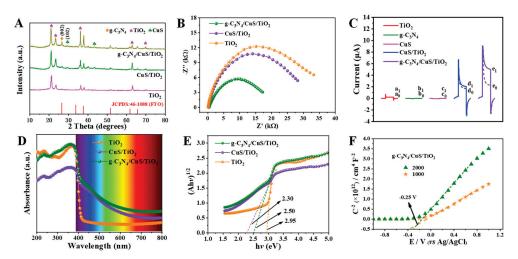


Figure 2. (A) X-ray powder diffractometry (XRD) and (B) Electrochemical impedance spectroscopy (EIS) plots of TiO_2 , CuS/TiO_2 , and $g-C_3N_4/CuS/TiO_2$. (C) Photocurrent responses of TiO_2 , CuS, $g-C_3N_4$, CuS/TiO_2 , and $g-C_3N_4/CuS/TiO_2$ -modified FTO electrodes in the (a_0-e_0) absence and (a_1-e_1) presence of 123.46 μ mol L^{-1} SM. (D) UV-Vis absorbance spectra and (E) Kubelka–Munk plots of TiO_2 , CuS/TiO_2 , and $g-C_3N_4/CuS/TiO_2$. (F) Mott–Schottky plots of $g-C_3N_4/CuS/TiO_2$ at frequencies of 1000 and 2000 Hz.

3.2. Photoelectrochemical Properties of the g-C₃N₄/CuS/TiO₂ PEC Sensing Platform

Next, electrochemical impedance spectroscopy (EIS) was performed to determine the conductivities of the prepared samples. As shown in Figure 2B, g- $C_3N_4/CuS/TiO_2$ yielded the smallest Nyquist semicircle radius, indicating that the composite had the lowest electron transfer resistance and higher conductivity than the other samples. This result can be attributed to the formation of a large number of heterointerfaces between the ternary semiconductor materials, which is conducive to charge transport.

The photocurrent response provides information about the optoelectronic properties and detection sensitivity of a material. Therefore, we investigated the photocurrents of the TiO₂, g-C₃N₄, CuS, CuS/TiO₂, and g-C₃N₄/CuS/TiO₂-modified FTO electrodes (Figure 2C). Owing to the slow separation of photoinduced electron-hole pairs, the single semiconductor materials (i.e., TiO_2 , g- C_3N_4 , and CuS) showed relatively low photocurrents both in the presence and absence of antioxidants. In contrast, the CuS/TiO2 electrode exhibited higher photoelectric activity: its photocurrent increased by 1.49 μA with respect to that of TiO₂ alone, without the addition of the antioxidant. This is because CuS nanoparticles have high absorbance in the visible-light region. However, the photocurrent only increased by $0.32~\mu A$ after the addition of $123.46~\mu mol~L^{-1}$ SM. Therefore, the sensitivity of the CuS/TiO₂ electrode is low and is not sufficient to detect SM. Thus, further modification of CuS/TiO2 by loading with g-C₃N₄ nanosheets is required. As expected, the photoelectric response and detection sensitivity of the ternary semiconductor composite (g-C₃N₄/CuS/TiO₂) were significantly improved. Notably, the photocurrent response of the ternary semiconductor composite increased by 0.62 µA in relation to that of the binary semiconductor composite, and the photocurrent was increased by 2.68 μ A in the presence of 123.46 μ mol L⁻¹ SM, suggesting that the ternary heterojunction PEC sensing platform is suitable for evaluating the antioxidant properties of SM. The significant increase in photocurrent can be attributed to two factors: (1) loading with g-C₃N₄ nanosheets increased the number of active sites for the reaction and enhanced light absorption, and (2) the heterojunctions in the ternary semiconductor composites promoted the reaction between antioxidants and holes and greatly improved the electron-hole separation.

To determine the optical properties, UV-Vis absorbance spectroscopy was carried out, as shown in Figure 2D. TiO₂ nanorods can only absorb light between 200 and 400 nm because of their wide band gap. However, the band gap was narrowed and more active sites for the reaction were produced after doping with CuS and g- C_3N_4 , which increased the collection and utilization of light. Therefore, the ternary g-C₃N₄/CuS/TiO₂ semiconductor exhibited greater light absorption at 200-800 nm, which is consistent with the EIS results. The ability of a sensor to detect antioxidants is strongly dependent on the band structure of the composites. Therefore, to achieve selective detection, the careful design of the valence band (VB) position of the semiconductor is required. As shown in Figure 2E, the relationship between the Kubelka-Munk function and the photon energy of the samples can be obtained from the UV-Vis absorbance spectra. The VB values of TiO2, CuS/TiO2, and $g-C_3N_4/CuS/TiO_2$ are 2.95, 2.30, and 2.50 eV, respectively. The Mott–Schottky (M-S) plots ofg-C₃N₄/CuS/TiO₂ semiconductor exhibited positive slopes (Figure 2F), suggesting that the composites are n-type heterostructures, and the flat-band potential (V_{fb}) was -0.25m V (vs. Ag/AgCl). In addition, the conduction band (CB) of the n-type semiconductor is 0.1 V smaller than its V_{fb} ; thus, the CB of g-C₃N₄/CuS/TiO₂ is -0.35 V (vs. Ag/AgCl).

3.3. Optimization of Experimental Conditions

The intrinsic properties of photosensitive materials are the main factors affecting the photocurrent signal, but other factors, such as the degree of doping of the semiconductor material, laser wavelength, and applied voltage, cannot be ignored. Optimization experiments were performed in PBS solution containing 123.46 μ mol L⁻¹ SM. As shown in Figure S3A, we investigated the influence of the amount of deposited CuS on the photocurrent. Crucially, an appropriate amount of CuS nanoparticles will increase the photocurrent, but an excessive amount could result in a drastic decrease in photocurrent. Based on the experiments, a

doping level of 5% was selected to prepare the CuS/TiO $_2$ composite. The mass ofg-C $_3$ N $_4$ is another key factor that affects the properties of semiconductors. As shown in Figure S3B, g-C $_3$ N $_4$ nanosheets provide more active sites for the reaction, and those prepared using a mass of 10 mg generated the highest photocurrent. As shown by the UV-Vis spectra, the g-C $_3$ N $_4$ /CuS/TiO $_2$ exhibited different absorbances in different bands. In particular, as an excitation light source, 630 nm red light has strong penetration. In addition, red light accounts for a large proportion of sunlight, suggesting its suitability for enhancing the sensitivity of the sensor for SM detection (Figure S3C). As shown in Figure S3D, there is a significant gradient between the semiconductor and FTO electrode in the initial state. When a potential is applied, electrons flow rapidly from the conductor band of the semiconductor to the FTO electrode to generate a photocurrent signal. The current at 0 V was 77.3% of that at +0.08 V, indicating that the PEC sensor has adequate sensitivity for SM detection. Therefore, 0 V was chosen as the working voltage for further experiments.

3.4. Antioxidant Assay and Detection Mechanism of the g-C₃N₄/CuS/TiO₂ PEC Sensing Platform

To analyze the detection performance of the as-prepared g-C₃N₄/CuS/TiO₂ PEC sensor upon 630 nm laser irradiation, the photocurrent signal was measured in PBS solutions (0.1 mol L^{-1}) having different concentrations of SM, PG, and the SM+PG mixture under the optimized assay conditions. As shown in Figure 3A₂, the photoresponse in SM detection exhibited a linear relationship with a detection limit (S/N = 3) of 1.8 μ mol L⁻¹, with three linear regression equations obtained for three concentration ranges, as follows: y = 0.005x + 2.121, $R^2 = 0.983$ $(2.157-100.196 \mu mol L^{-1})$, y = 0.003x + 2.353, $R^2 = 0.997 (100.196-725.196 \mu mol L^{-1})$, and y = 0.002x + 3.512, $R^2 = 0.999$ (725.196–1276.667 µmol L⁻¹). Compared with the SM detection performance of previously reported methods (Table S2), the SM detection performance of the g-C₃N₄/CuS/TiO₂ PEC sensing platform is excellent. In particular, the proposed sensing platform exhibits a distinctly higher detection sensitivity and a wider linear detection range, which completely satisfy actual sample detection requirements. The photocurrent response of the g-C₃N₄/CuS/TiO₂ PEC sensor to PG was lower than that of the SM sensor (Figure 3B₁), and two linear relationships were obtained (Figure 3B₂): y = 0.003x + 2.832, $R^2 = 0.991$ $(71.787-484.812 \text{ }\mu\text{mol }\text{L}^{-1})$ and y = 0.001x + 3.853, $R^2 = 0.987$ (484.812–1276.667 $\mu\text{mol }\text{L}^{-1}$). Concerning the antioxidant mixtures of SM+PG, the photoresponse displayed favorable linear relationships in the three concentration ranges (Figure $3C_1$), and the corresponding linear equations are y = 0.079x + 3.606, $R^2 = 0.983$ (15.990–94.627 μ mol L⁻¹); y = 0.017x + 9.492, $R^2 = 0.991$ (94.627–735.132 µmol L⁻¹); and y = 0.005x + 18.292, $R^2 = 0.989$ (735.132–1276.667 μ mol L⁻¹), respectively. These results indicate that the photoelectric response from the two-component mixture of antioxidants was significantly greater than the combined photocurrents of individual components at the same concentration. This indicates a synergistic effect between SM and PG, especially at high concentrations. Based on these findings, we investigated the synergistic effects of SM and VE, BHA, BHT, and TBHQ. As shown in Figure S4, SM was mixed in pairs with these four typical antioxidants, and the obtained trends are similar to that of the two-component mixtures of SM+PG. Presumably, an additive effect occurs at low concentrations, and a synergistic effect occurs at high concentrations. It has been reported that long distances suppress electron transport between antioxidants at low concentrations. Therefore, at relatively high concentrations, frequent electron transfer and the coupling with the oxidant may promote the regeneration of antioxidants, thus resulting in a more pronounced synergistic effect [29–31].

To understand the sensing principle of our sensor, as well as the observed synergistic effects, the electrochemical properties and photoresponses of the antioxidants were investigated in detail. As shown in Table S1, cyclic voltammetry of the antioxidants at a concentration of 484.812 μ mol L⁻¹ was carried out, yielding the following redox potentials: SM (0.373 V) < VE (0.464 V) < TBHQ (0.531 V) < BHA (0.539 V) < BHT (0.550 V) < PG (0.657 V). SM exhibited the smallest redox potential, suggesting its ability to reduce antioxidants having higher redox potentials [29,32]. As shown in Figure 4 and Table S1, the synergistic effect of binary antioxidant mixtures SM+VE, SM+TBHQ, SM+BHA, SM+BHT,

and SM+PG resulted in increases of 1.57-, 1.38-, 2.31-, 2.42-, and 2.67-times, respectively, compared to the sum of the single antioxidant. This result is similar to the trend in the redox potential of the antioxidants. However, the synergistic effect of SM+VE was more obvious than that of SM+TBHQ, which is inconsistent with the increasing trend in redox potential. This is because both SM and VE are natural antioxidants, and their molecular structure and electronic properties (nucleophilic, electrophilic, ionization potential, and phenolic bond dissociation energy) are similar, thereby showing better performance [33]. Furthermore, we investigated the antioxidant effects of mixtures of natural antioxidants (SM+VE), four artificial antioxidants (BHA+BHT+TBHQ+PG), and six natural and artificial antioxidants (SM+VE+BHA+BHT+TBHQ+PG). All the samples showed improvements in the photocurrent responses, SM+VE, BHA+BHT+TBHO+PG, and SM+VE+BHA+BHT+TBHO+ PG showed increases of 1.57-, 1.64-, and 1.56-times, respectively, in relation to the sum of photoresponses of all single antioxidants. These results indicate that there are synergistic and antagonistic effects between the antioxidant molecules. In particular, there is a negative correlation between complexity of the antioxidant components and the antioxidant effect; thus, it is crucial to select a suitable formula to maximize the antioxidant performance [34]. These data also indicate that natural antioxidants exhibit good antioxidant properties and could be used to replace synthetic antioxidants as food stabilizers.

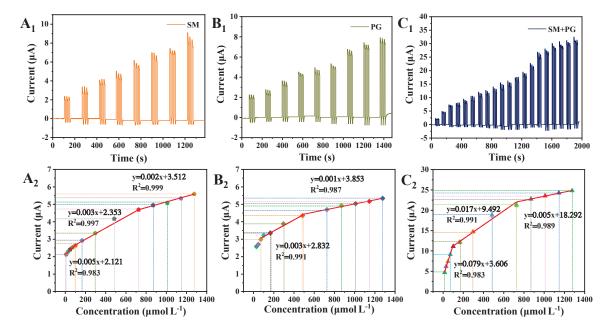


Figure 3. Photocurrent response curves of SM (A_1) , PG (B_1) , and SM+PG in equal proportions (C_1) generated for the g-C₃N₄/CuS/TiO₂-based PEC sensing platform at different analyte concentrations. Calibration plot of photocurrent versus different concentrations of SM (A_2) , PG (B_2) , and SM+PG (C_2) .

Based on the above results, we conclude that coupling oxidation plays a key role in the synergism of the studied antioxidants: specifically, it reduces the potential difference between pairs of antioxidants and promotes their regeneration, thus improving the performance of binary mixtures of antioxidants. Additionally, the heterojunctions in the $g-C_3N_4/CuS/TiO_2$ composite promote the reaction between antioxidants and holes, thus improving the separation and transition of electrons and holes (Scheme 1) and amplifying the detectable signal. Therefore, the $g-C_3N_4/CuS/TiO_2$ -modified platform could be applied for the detection and evaluation of the synergistic effect of SM based antioxidants.

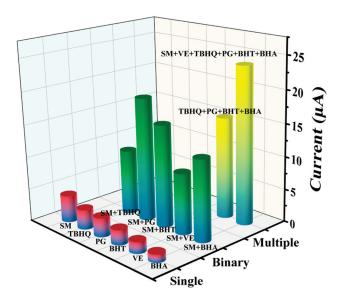


Figure 4. Photocurrents of the g- $C_3N_4/CuS/TiO_2$ PEC sensor in the presence of 484.812 μ mol L⁻¹ SM, VE, TBHQ, BHA, BHT, PG, and a mixture of SM and other antioxidants in equal molar ratios.

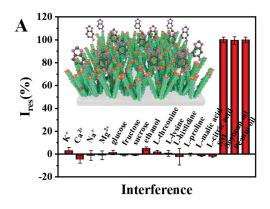
3.5. Detection Selectivity Stability and Reusability ofg-C₃N₄/CuS/TiO₂ PEC Sensing Platform

Sensor selectivity depends largely on the band structure of the composite. Specifically, an appropriate VB can oxidize target species with a lower redox potential, but does not react with the sugars and amino acids in food. As shown in Figure 5A, we tested the effect of 18 times of fructose, glucose, sucrose, L-malic acid, L-citric acid, ethanol, L-threonine, L-proline, L-lysine, and L-histidine, as well as 700 times of Na⁺, K⁺, Mg²⁺ and Ca²⁺ on photocurrent based on g-C₃N₄/CuS/TiO₂ PEC sensor with 123.46 μ mol L⁻¹ SM. No detectable signals (<10%) corresponding to the interfering agents were detected, suggesting that the PEC sensor hashigh SM detection selectivity. Furthermore, we had tested the photocurrent of real samples (soybean oil and peanut oil) in the same conditions of SM, there are insignificant change in the photocurrent response, which further verified the feasibility of the sensor in the analysis of antioxidant capacity of food.

Sensor stability and reusability are also important criterions to evaluate the performance of PEC sensor, especially for practical, long-term detection in real samples. As shown in Figure 5B, the g-C₃N₄/CuS/TiO₂ PEC sensor was monitored using amperometry with 500 s of "on–off–on" cycling in the presence of 1270.574 μ mol L⁻¹ SM, which showed no obvious decrease in the photoelectric signal. After 15 days, the photocurrent signal was reduced to 96.8% of the initial state, indicating that the sensor has excellent stability and reusability.

3.6. SM Detection in Soybean and Peanut Oils

Oils and fats naturally become rancid as a result of oxidation during storage or processing. The addition of synthetic or natural antioxidants can slow down lipid peroxidation and speed up the scavenging of free radicals [11,35,36]. In view of the high selectivity of the g- $C_3N_4/CuS/TiO_2$ PEC sensor platform, the applicability of the sensor for use with real oil samples was investigated using soybean and peanut oils containing various SM concentrations. The corresponding calibration curve is shown in Figure 3A. As listed in Table 1, both soybean and peanut oil samples exhibited high recoveries and low relative standard deviations (RSD; <3%) at different SM concentrations, indicating excellent detection precision and accuracy. Therefore, the natural antioxidant SM showed excellent antioxidant properties and could replace synthetic antioxidants as quality stabilizers for edible oils, which is important for food safety, quality control, and ingredient screening.



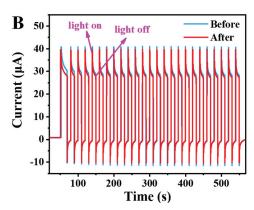


Figure 5. (A) Samples and interferences test in the presence of fructose, glucose, sucrose, L-malic acid, L-citric acid, ethanol, L-threonine, L-proline, L-lysine, and L-histidine, as well as Na⁺, K⁺, Mg²⁺, and Ca²⁺, on the photocurrent response of the g-C₃N₄/CuS/TiO₂-based PEC sensor. (B) Detection stability and reusability (after 15 days) of the g-C₃N₄/CuS/TiO₂-modified FTO electrode in the presence of 1270.574 μ mol L⁻¹ SM. The PEC sensor was operated at 0 V under 630 nm light excitation in 0.1 mol L⁻¹ PBS (pH = 7.4).

Table 1. SM Concentrations in Soybean and Peanut Oil Samples Detected Using the $g-C_3N_4/CuS/TiO_2$ -based PEC Sensor.

Sample	SM Added (µM)	SM Found (µM)	RSD (%)	Recovery (%)
Soybean oil	4	3.93 ± 0.01	0.31	98.28
	20	20.40 ± 0.20	0.98	102.00
Peanut oil	4	3.97 ± 0.09	2.30	99.16
	20	20.40 ± 0.53	2.59	102.00

4. Conclusions

We developed a novel red-light-mediated signal-on PEC sensor based on a g- $C_3N_4/$ CuS/TiO₂ ternary heterojunction for a highly sensitive analysis of the synergistic antioxidant effects of SM and other antioxidants. The g-C₃N₄/CuS/TiO₂ heterojunction forms a Schottky barrier that facilitates electron transfer. When SM and other antioxidants are present, they scavenge the holes in the VB of g-C₃N₄/CuS/TiO₂, suppressing the recombination of photogenerated charge carriers, thus increasing the PEC photocurrent. The sensor also allowed the rapid evaluation of the synergistic effects of SM and VE, BHA, BHT, TBHQ, and PG. The synergistic antioxidant reaction of multiple antioxidants is mainly driven by coupling oxidation, which enhances the total antioxidant performance by lowering the potential difference between the antioxidants. We found that SM reduces the free radicals of other antioxidants and promotes their regeneration because of its low redox potential. Overall, the g-C₃N₄/CuS/TiO₂ PEC sensor exhibited high sensitivity, good selectivity, and stability, and was successfully applied to the detection of SM in soybean and peanut oils. The reported sensor has applications in nutrition analysis and the development of nutritional additives, which could help improve the quality of life of consumers. In addition, the sensor has potential uses in food fraud analysis, as well as medical research.

Supplementary Materials: The following supporting information can be downloaded at: https://www.mdpi.com/article/10.3390/bios13090859/s1: Reagents; Apparatus; Figure S1: SEM images of TiO₂ and CuS/TiO₂; Figure S2: XP spectra of g-C₃N₄/CuS/TiO₂; Figure S3: Optimization of experimental conditions; Figure S4: Photocurrent response curves; Table S1: Electrochemical characteristics and synergistic effects; Table S2: Comparison of the SM detection performance of various methods. References [10,37–39] are cited in Supplementary Materials.

Author Contributions: Conceptualization, L.H.; methodology, L.H.; validation, Z.L., H.L. and Z.S.; formal analysis, L.H., J.Y. and R.L.; investigation, L.H. and L.N; data curation, J.Y., Z.S., Z.L., R.L. and H.L.; writing—original draft preparation, D.H.; writing—review and editing, L.H.; supervision, L.N. All authors have read and agreed to the published version of the manuscript.

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Article

Hierarchical Ti-MOF Microflowers for Synchronous Removal and Fluorescent Detection of Aluminum Ions

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Abstract: Bifunctional luminescence metal-organic frameworks with unique nanostructures have drawn ongoing attention for simultaneous determination and elimination of metal ions in the aqueous environment, but still remain a great challenge. In this work, three-dimensional hierarchical titanium metal-organic framework (Ti-MOF) microflowers were developed by a secondary hydrothermal method for not only highly sensitive and selective detection of Al(III), but also simultaneously efficient decontamination. The resulting Ti-MOF microflowers with a diameter of 5-6 µm consisted of nanorods with a diameter of \sim 200 nm and a length of 1–2 μ m, which provide abundant, surface active sites for determination and elimination of Al(III) ions. Because of their substantial specific surface area and superior fluorescence characteristics, Ti-MOF microflowers are used as fluorescence probes for quantitative determination of Al(III) in the aqueous environment. Importantly, the specific FL enhancement by Al(III) via a chelation-enhanced fluorescence mechanism can be utilized for selective and quantitative determination of Al(III). The Al(III) detection has a linear range of 0.4–15 μM and a detection limit as low as 75 nM. By introducing ascorbic acid, interference of Fe(III) can be avoided to achieve selective detection of Al(III) under various co-existing cations. It is noteworthy that the Ti-MOF microflowers exhibit excellent adsorption capacity for Al(III) with a high adsorption capacity of 25.85 mg g^{-1} . The rapid adsorption rate is consistent with a pseudo-second order kinetic model. Ti-MOF is a promising contender as an adsorbent and a fluorescent chemical sensor for simultaneous determination and elimination of Al(III) due to its exceptional water stability, high porosity, and intense luminescence.

Keywords: luminescence; metal-organic frameworks; hierarchical nanostructures; fluorescent sensors; metal ions

1. Introduction

Aluminum is a substance that is commonly present in the environment and in human activities, is thought to hinder the growth of plants, and is linked to diseases such as Alzheimer's and Parkinson's diseases [1]. The World Health Organization (WHO) recommends that the maximum aluminum content in drinking water should be limited to 7.4 μM , and the daily consumption of aluminum for humans should be between 3 and 10 mg [2]. Measurement and purification of Al(III) is essential for environmental monitoring and preserving human health, since aluminum often exists in the body and environment as the Al(III) ion [3]. Consequently, significant efforts are being undertaken to develop novel

approaches for the selective detection and removal of Al(III) because of its pervasive contamination and high toxicity. For sensitive detection of Al(III) pollution in the aqueous environment, a variety of atomic absorption or emission spectrometry methods, inductively coupled plasma mass spectrometry (ICP-MS), and fluorescence methods, have been utilized [4,5]. Due to its ease of use, cheap cost, convenience, sensitivity, real-time monitoring, and naked-eye visibility, fluorescence based chemosensors have been identified as promising tools for detecting Al(III). Nevertheless, sensitive and selective detection of Al(III) is still tedious, with several disadvantages due to its lack of spectroscopic features, sluggish coordination, and a great hydration capacity.

Toxic metals must be removed from polluted water. Many methods have been developed recently to remove harmful metals from diverse substances, including ion displacement, membrane filtration, redox co-precipitation, chemical deposition, absorption methods, and solid-phase extraction [6,7]. Among these technologies, adsorption is believed to be the most cost effective method to remove Al(III) from contaminated water because of its simplicity of use and the availability of a variety of adsorbents, such as Streptomyces rimosus biomass, Fe₃O₄/TEOS/AMEO/GA, activated carbon, and polyacrylonitrile beads [8]. Unfortunately, most of the previous studies either concentrated on functional adsorption materials to improve selectivity and sensitivity towards metal ion detection or centered on absorbing materials to boost their uptake capacity for metal ions, which seriously hindered their practical applications. Developing a novel bifunctional material that can both identify and eliminate metal ions from environmental samples is of ongoing interest, and would significantly increase the fluorescence signal due to effective adsorption. Qian et al. reported a covalently grafted naphthalimide derivative of 2,6-bis(aminomethyl)pyridine on the surface of silica particles to create a bifunctional fluorescence sensor for the simultaneous detection and separation of trace Hg²⁺ in polluted water samples [9]. Manos et al. developed a microporous metal-organic framework $(H_{16}[Zr_6O_{16}(H_2PATP)_4]Cl_8 \cdot xH_2O)$ that has an exceptional capacity to quickly collect and selectively detect hexavalent chromium in aqueous solution [10]. Our group also reported the selective detection and simple removal of arsenate from contaminated water using amino-functionalized iron-based metal organic framework (NH2-MIL-88(Fe)) nanooctahedra [11]. However, to the best of our knowledge, such bifunctional materials that are capable of detecting and removing Al(III) from environmental samples simultaneously have not been reported to date.

To meet this demand, we attempted development of a bifunctional material with the following characteristics: suitable recognition sites to provide specific receptor-target interactions for luminescence signal responses, strong metal chelating groups with high affinity toward specific target metal ions, ordered and high-density accessible binding sites, exceptional water and chemical stability to facilitate multiple uses, and high efficiency in converting light into electrical energy. To fulfill this goal, luminescence metal-organic framework (LMOF) materials are ideal candidates that offer a special scaffold and tailored functionality with luminescence in integrating two functions of sensing and adsorption into a single material, in contrast to conventional sensing and remediation materials [12-15]. Numerous experiments using MOF-based sensors have recently been published to detect Al(III) [16,17]. For instance, Zhan, et al. developed a water-resistant terbium-MOF sensor for the accurate, precise, and recyclable detection of Al (III) [17]. To ratiometrically detect Al³⁺ ions in an aqueous solution, Zheng et al. created a new zirconium-based MOF composite material called UiO-(OH)₂@RhB [18]. Design of an MOF-based sensor for simultaneous detection and removal of Al(III) in an aquatic environment, however, remains a formidable issue.

Adsorption efficiency can be significantly impacted by changing the geometry of adsorbents. This suggests that in order to reduce the diffusion barrier and increase adsorption efficiency, research on the design of adsorbents at the geometrical level should be pursued. The integration of nano-building blocks with the proper order and flexibility to create hierarchical nanostructures has proven to be a successful strategy for improving adsorption

performance [19]. Hierarchical nanostructures have been successfully created using a wide range of materials, including metals, metal oxides and sulfides, carbons, and certain organic polymers [20]. Metal-organic hierarchical nanostructures, however, are seldom explored. Even though MOF sensing applications have been of interest, it is often challenging to create a free-standing, three-dimensional (3D) hierarchical micro-/nano-structured MOF that has nanoscale thickness and microsized lateral dimensions. A collection of hierarchical metal-organic framework nanostructures has recently been published by Wang et al. [21,22]. These substances differ from their pure organic or inorganic equivalents in terms of functions, and they may find use in a variety of industries. Based on these advantages, we assumed that the integration of 3D hierarchical micro-/nano-structure and bifunctional LMOF would be highly desired and could improve fluorescence sensing and adsorption performance.

Following these criteria, we developed a novel approach for building three-dimensional (3D) hierarchical titanium-based MOF (Ti-MOF) microflowers by a secondary hydrothermal process for potential use dual-functional fluorescence sensors and adsorbents (Scheme 1). Using a variety of methods, including scanning electron microscopy (SEM), transmission electron microscopy (TEM), X-ray diffractometer (XRD), and fluorescence spectroscopy, the structural features of the obtained product were carefully examined. The sensing characteristics of the developed materials toward Al(III) were investigated, including sensitivity and selectivity in the presence of various interfering ions, to illustrate possible bifunctional applications. Meanwhile, the concomitant adsorption performance of Al(III) onto hierarchical Ti-MOF microflowers was carefully evaluated. Using X-ray photoelectron spectroscopy (XPS) and Fourier transform infrared (FT-IR) spectroscopic studies, the fundamental sensing and adsorption mechanism of hierarchical Ti-MOF microflowers towards Al(III) was investigated.



Scheme 1. Schematic illustration for manufacturing hierarchical Ti-MOF microflowers and Al(III) ions luminous sensing.

2. Experimental Section

2.1. Synthesis of Hierarchical Ti-MOF Microflowers

According to a previously published approach, with certain modifications, hierarchical Ti-MOFs microflowers were created by a secondary hydrothermal process [23,24]. First, 178 mg Ti(i-OPr)₄ was dissolved in 10 mL acetic acid with stirring until a milk white emulsion formed. Then, after being transferred, the mixture was heated at 150 °C for 48 h in a 30 mL Teflon-lined stainless-steel autoclave. The autoclave was allowed to cool naturally to ambient temperature before the solution was centrifuged for six minutes at 8000 rpm to remove the supernatant solution, and washed with deionized water three times. Secondly, the achieved white precipitate was re-dispersed in 10 mL of deionized water, and 124 mg $\rm H_4DOBDC$ was injected to the above solution with stirring. A 30 mL Teflon-lined stainless-steel autoclave was used to transfer the mixture, which was heated at 150 °C for 48 h. Finally, the dark red crystals were centrifuged at 8000 rpm for 6 min and rinsed with deionized water and ethanol three times.

To confirm that the synthesis of Ti-MOFs microflowers was achieved by employing the secondary hydrothermal method, direct hydrothermal reaction of Ti(i-OPr)₄ (178 mg) and H₄DOBDC (124 mg) in 10 mL acetic acid was carried out in a 48-h experiment in a 30 mL Teflon-lined stainless-steel autoclave. The produced samples were centrifuged at 8000 rpm for 15 min and rinsed with deionized water and ethanol three times.

2.2. Fluorescence Determination of Aluminium Ion

In a typical procedure, the dispersed solutions were made by combining 200 mL of deionized water with 83.5 mg of Ti-MOFs microflowers, followed by 30 min of ultrasonic agitation. Various concentrations of aluminum ion solutions were prepared in deionized water solution. The fluorescence spectrophotometer excitation wavelength was set at 355 nm, and the excitation and emission slit widths were both 2.0. A volume of 3 mL of the sample in a quartz cuvette was used for the fluorescence measurements. For fluorescence detection of Al(III), 60 μ L of Al(III) solutions at various concentrations were individually added to 2940 μ L stock solutions of Ti-MOFs microflowers, the final Al(III) concentration range being from 0.4 to 100 μ M.

To investigate the selectivity of the Ti-MOFs microflowers fluorescent probe towards Al(III) determination, various interference ions, such as Cu(II), Cd(II), Ag(I), Co(II), Ni(II), Ca(II), Mg(II), Fe(III), Fe(III), Hg(II), Mn(II), Zn(II), NH₄ $^+$, were separately added to the 2970 μL stock solutions of Ti-MOFs microflowers containing 30 μL of 1 mM Al(III). The concentration of all interference ions was 100 μM , i.e., 10-fold higher than that of Al(III), and the same detection conditions were employed as in the aforementioned Al(III) sensing tests.

The following figures and tables provide averages and standard deviations from all selectivity and selectivity studies, which were carried out in triplicate.

2.3. Adsorption of Aluminum Ion

For the kinetic adsorption experiment, 300 mg of hierarchical Ti-MOFs microflowers were mixed with 300 mL of 50 mg L $^{-1}$ Al(III) solution. The mixture was agitated at 150 rpm at 25 °C, and the pH of the solution was kept constant at 4.7 \pm 0.1 by adding tiny amounts of either HCl (0.1 M) or NaOH (0.1 M). At predetermined intervals, 3 mL samples of the suspension were taken. A 0.22 μm membrane filter was used to filter the samples before inductively coupled plasma optical emission spectrometry (ICP-OES) analysis (ICP-6000, Thermal Electron, Waltham, MA, USA). Thus, the ICP-OES could independently acquire the amount of aluminum present in the Al(III) containing solution.

Studies of isothermal adsorption were carried out at pH 4.7 \pm 0.1 and 25 °C. To attain adsorption equilibration, 20 mg of the as-prepared hierarchical Ti-MOF microflowers were introduced to 20 mL of Al(III) solution (2–100 mgL $^{-1}$) and shaken for 96 h. The aforementioned solution was filtered using a 0.22 μm membrane filter after the liquid and solid phases were separated by centrifugation. The concentration of Al(III) was analyzed by ICP-OES after filtration. The Al(III) adsorption amount was calculated by Equation (1):

$$q_e = \frac{(C_0 - C_e) \times V}{m} \tag{1}$$

where q_e is the amount of Al(III) adsorbed on adsorbent at equilibrium time (mg·g⁻¹), and C_0 and C_e represent the initial and equilibrium Al(III) concentrations (mg·L⁻¹) in the solution, respectively, V indicates the solution volume (mL) and m is the mass of the adsorbent (mg).

A pseudo-second-order kinetics model was used to analyze the adsorption kinetic data on the assumption that chemisorption was the rate-determining phase. The pseudo-second-order kinetics mode is presented in Equation (2):

$$\frac{t}{Q} = \frac{1}{Q_e}t + \frac{1}{k_2 Q_e^2} \tag{2}$$

where k_2 is the pseudo-second-order rate constant (g·mg⁻¹h⁻¹), Q_e is the equilibrium adsorption capacity (mg·g⁻¹), and Q is the amount of Al(III) ions adsorbed on the surface of Ti-MOF at time t (min), respectively. A linear plot of t/Q yields the values of Q_e and k_2 .

The capacities of absorbents were determined by equilibrium adsorption isotherms. The experimental data are fitted to Langmuir models.

$$Q_e = \frac{Q_m K_L C_e}{1 + K_L C_e} \tag{3}$$

where C_e is the equilibrium concentration of Al(III) in aqueous solution (mg·L⁻¹), Q_e represents the amount of Al(III) adsorbed on the surface of Ti-MOF (mg·g⁻¹), Q_m indicates the maximum amount of Al(III) adsorbed per unit weight of Ti-MOF to form a complete monolayer coverage on the surface, and K_L is the ratio of the rate constants of adsorption and desorption.

3. Results and Discussion

Monodisperse hierarchical Ti-MOF microflowers were produced by a quick, easy, and surfactant-free secondary hydrothermal technique. Through initial hydrothermal synthesis, the precursor of Ti-MOF was fabricated using Ti(i-OPr)₄ as a metal precursor, and acetic acid as the linker and solvent. Acetate groups (OAc⁻) interact with titanium alkoxides as complexing nucleophilic ligands. Acetate groups function as bidentate bridging ligands during the stoichiometric process, which also results in an increase in the coordination number of Ti from four to six, and the formation of oligomeric species $Ti(i-OPr)_{4-x}/(AcOH)_x$. Using SEM and TEM measurements, the shape and structure of the $Ti(i-OPr)_{4-x}/(AcOH)_x$ complex were identified. As illustrated in Figure 1A, the synthesized $Ti(i-OPr)_{4-x}/(AcOH)_x$ complex exhibited a great number of nanoscale petals combine to produce a monodispersed flower-like hierarchical structure, which typically had a diameter between 1–3 μm. The hierarchical microflowers are shown in a high-magnification TEM image in Figure 1B. It is evident that these microflowers are made up of different 2D nanosheets (petals), which grew from the flower center in all directions to construct 3D hierarchical nanostructures. The petals of the flower-like $Ti(i-OPr)_{4-x}/(AcOH)_x$ complex have a slightly curved compact structure and an average length of about 500 nm (Figure 1C). It is important to note that the production of the aforementioned $Ti(i-OPr)_{4-x}/(AcOH)_x$ complex with flower-like hierarchical structures depends critically on acetic acid. By contrast, by using pure deionized water to replace acetic acid, only highly aggregated, irregular particles were produced. Subsequently, the prepared flower-like $Ti(i-OPr)_{4-x}/(AcOH)_x$ complex was a precursor in a reaction with H₂DOBDC to produce Ti-MOF through a second hydrothermal synthesis in which AcOH was completely substituted by H₂DOBDC. The Ti-MOF completely inherited the original flower-like morphology from the $Ti(i-OPr)_{4-x}/(AcOH)_x$ complex. Interestingly, the petals of $Ti(i-OPr)_{4-x}/(AcOH)_x$ complex had smooth surfaces, while the petal building block of Ti-MOF, owing to the substitution of acetic acid by H₂DOBDC, had a perforated structure made of linked or stacked nanoparticles (10–20 nm in size) with a rough surface, as shown in Figure S1. Overall, the microflower morphologies were conserved during the hydrothermal process even if expanding pore channels affected the internal architecture. However, the direct hydrothermal reaction of Ti(i-OPr)₄ and H₄DOBDC in acetic acid resulted in pseudo-spherical Ti-MOF nanoparticles 200 ± 20 nm in size (Figure S2). The Ti-MOF microflowers precursor was Ti₂(HDOBDC)₂(H₂DOBDC) (NTU-9) [24]. This has two-dimensional (2D) hybrid layers that are aligned perpendicular to the c-axis in hexagonal prism crystals. Each layer has two-dimensional honeycomb-like layers and is built from the bond between Ti4+ and DOBDC ligands. The Ti4+ is octahedrally coordinated with six oxygen atoms from the hydroxide and carboxylate groups of the ligand DOBDC, with Ti-O bond lengths of 1.858 and 2.037 Å, respectively. Inside the layers, there are 11 \times 11 Å² one-dimensional hexagonal channels that all the uncoordinated oxygen atoms point to. Due to the full exposure of active sites and the low mass transfer diffusion barrier, such a

hierarchical, porous structure would, in theory, would be very advantageous for sensing and adsorption processes [19].

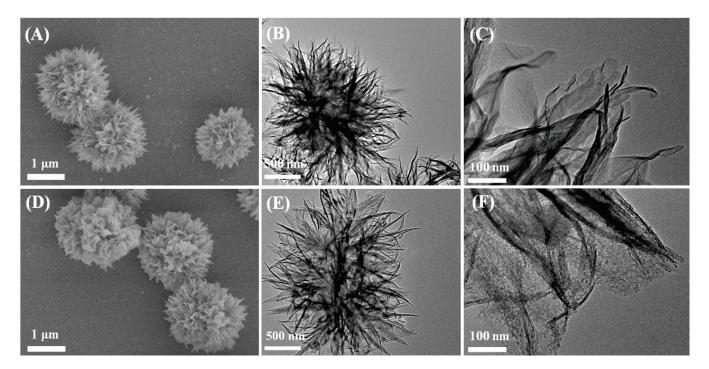


Figure 1. SEM images of Ti(i-OPr)_{4-x}/(AcOH)_x complex (**A**), and hierarchical Ti-MOF microflowers (**D**). TEM images of Ti(i-OPr)_{4-x}/(AcOH)_x complex (**B**,**C**), and hierarchical Ti-MOF microflowers (**E**,**F**).

The crystalline structure and surface analysis of hierarchical Ti-MOF microflowers were investigated by XRD and FT-IR. As illustrated in Figure 2A, the principal diffraction peaks in the experimental XRD pattern of the hierarchical Ti-MOF microflowers demonstrate a clearly defined crystal structure, and are in good agreement with the simulated patterns of the NTU-9 family [24]. The sharp peaks show the successful fabrication of Ti-MOF and indicate the exceptional crystallinity of the framework. Moreover, the FT-IR measurement was employed to confirm the surface functional groups of Ti-MOF. As illustrated in Figure 2B, Ti-MOF microflowers clearly displayed typical peaks at 3149–3435, 1650, 1496, 1429, 1360 and 1209 cm⁻¹ due to the vibrational stretching modes of hydroxyl groups, the C=O group, carbon atoms in the phenyl ring, the O-H group, in-plane O-H bending of phenolic hydroxy, and asymmetric stretching of the C-C-O group, in the same peak position as the H₄DOBDC ligand. The FT-IR analysis indicates that hierarchical Ti-MOF microflowers are inherently hydrophilic since the structures contain a large number of oxygen-rich functional groups.

A hierarchical structure with a porous structure is required for an adsorbent to function, and would be beneficial to mass transport. Thus, to characterize the BET surface area and internal pore superstructure of pseudo-spherical Ti-MOF NPs and hierarchical Ti-MOF microflowers, nitrogen adsorption-desorption measurements were performed. As shown in Figure 2C, the isotherm was classified as type IV, with a distinct hysteresis loop at relative pressures ranging from 0.5 to 1.0. The BET-specific surface area of hierarchical Ti-MOF microflowers was estimated to be approximately 264.47 m $^2 \cdot g^{-1}$, which is over 5-fold larger than that of pseudo-spherical Ti-MOF NPs (46.63 m $^2 \cdot g^{-1}$). Moreover, the hysteresis isotherms indicated that Ti-MOF microflowers have a micro/nano-porous structure, which is in line with SEM observations. The hysteresis loop in the Ti-MOF microflowers can be classified as Type H1, indicating a narrow pore size distribution (pore diameter on average: 1.8 nm) according to the Barrett-Joyner-Halenda (BJH) model, as illustrated in the inset of Figure 2C.

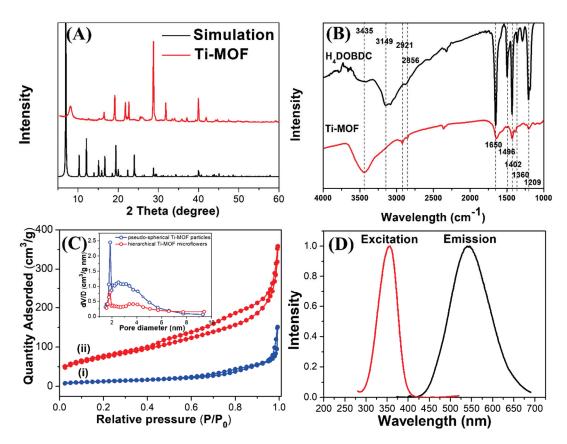


Figure 2. (**A**) Simulated and experimental powder XRD patterns for Ti-MOF microflowers. (**B**) FT-IR spectra of H₄DOBDC and Ti-MOF microflowers. (**C**) N₂ adsorption-desorption isotherms and the pore size distribution (inset image) of (i) pseudo-spherical Ti-MOF particles and (ii) hierarchical Ti-MOF microflowers. (**D**) Fluorescence excitation (red line) and emission (blue line) spectra of Ti-MOF microflowers.

To further investigate the optical properties of Ti-MOF microflowers, fluorescence spectrum studies were carried out. As shown in Figure 2D, the maximal excitation and emission peaks were located at 355 and 538 nm, respectively. A large Stokes shift of 187 nm occurred in the Ti-MOF microflowers to avoid crosstalk between excitation and emission signals. The quantum yield of Ti-MOF microflowers was estimated at 7.49%. Furthermore, due to their insufficient water solubility, most Al(III)-selective fluorescence chemosensors are studied in pure organic (such as DMSO and THF) or organic-water mixed solutions [25]. Therefore, the effects of various conditions on Ti-MOF microflowers should be investigated in order to determine fluorescence stability. As displayed in Figure S3, no distinct photobleaching under UV lamp irradiation over 5 h was observed, demonstrating the remarkable optical stability of Ti-MOF microflowers compared to traditional fluorescent dyes. The fluorescence properties of Ti-MOF microflowers were also investigated at different ionic strengths (modulated by 0-1000 mM NaNO₃) and different pH (4.0-9.0) conditions, as illustrated in Figure S4. There was no discernible change in fluorescence intensity in aqueous solutions under different ionic strengths and pH conditions, indicating that Ti-MOF microflowers are extremely stable even under high ionic strength and extreme pH conditions. Moreover, the Ti-MOF microflowers solution remained homogenous for over 1 month at room temperature (without any aggregation or color change). Their excellent water solubility and optical stability makes them potential candidates for a new class of fluorophores for determination of metal ions in aqueous solution without organic co-solvents.

Owing to their excellent fluorescence property and water stability, Ti-MOF microflowers are proposed to be used as fluorescent probes for quantitative determination of Al(III) ions. Because response rate is important in sensing performance, the time-response charac-

teristics of hierarchical Ti-MOF microflowers towards Al(III) were investigated, in which the fluorescence signals of the reaction mixture were collected immediately and consecutively upon the addition of 10 M Al(III) into the Ti-MOF microflowers suspension. Figure S5 shows the difference in FL intensity before and after addition of 10 M Al(III). The fluorescence intensity of Ti-MOF microflowers significantly increased within 2 min after Al(III) addition and remained essentially constant afterward, indicating that the Al(III)-induced fluorescence enhancement reaction is quick. The FL intensity-Al(III) concentration ([Al(III)]) relationship was then established. Figure 3 shows the change of fluorescence intensity of Ti-MOF microflowers with gradually increasing Al(III) concentration from 0 to 100 μ M for evaluating their fluorescence sensitivity. Figure 3A shows that the maximum emission peak of Ti-MOF microflowers appears at 538 nm, and is blue-shifted to 510 nm with remarkable fluorescence enhancement with increasing concentrations of Al(III). Figure 3B shows a plot of fluorescence enhancement rate (FER, derived from the peak intensity data in Figure 3A) against Al(III) concentration. There is an excellent linear correlation (Y = 0.114X - 0.024, $R^2 = 0.999$) between the FER and [Al(III)] in the low concentration range of 0.4–15 μ M. The real limit of detection (LOD) is $0.4 \mu M$, based on Figure 3B. Furthermore, the maximum FER could reach 2.34 as the concentration of Al(III) increased to 100 μM. On the basis of a signal-to-noise ratio of $3\sigma/\kappa$ (κ : slope; σ : standard error), the theoretical LOD was estimated to be 75 nM, which is approximately 100-fold lower than the permitted level (7.4 µM) of aluminum in drinking water by the WHO. Ti-MOF analytical performance is comparable to, or better than, that of previously reported FL sensors for Al(III) detection, as shown in Table S1. The poor coordination ability and strong hydration ability of Al(III) has hampered the development of a suitable fluorescence sensor [26]. A Job's plots experiment was carried out to demonstrate the coordination of Al(III) and Ti-MOF microflowers. The maximum fluorescence intensity as a function of molar fraction of Al(III) was used in the Job's method, and the total concentration of Ti-MOF and Al(III) ion was 20 μM. As shown in Figure S6, 3:2 stoichiometric complexations between Ti-MOF microflowers and Al(III) was confirmed.

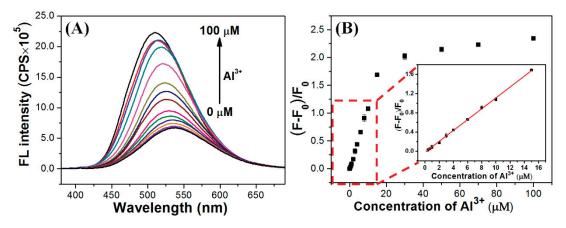


Figure 3. (A) Fluorescence response of Ti-MOF microflowers with various concentration of Al(III) (from bottom to top: 0, 0.4, 0.7, 1, 2, 3, 4, 6, 8, 10, 15, 30, 50, 70, 100 μ M). (B) Fluorescence enhancement rate ((F–F₀)/F₀ at 510 nm) of Ti-MOF microflowers to Al(III) at various concentrations (0 to 100 μ M). The linear detection range for 0.4–15 μ M. Al(III) is shown in the inset image. The error bars represent standard deviations from three independent measurements.

Because of the similar electron configuration in the recognition process, other trivalent ions such as Fe^{3+} and Cr^{3+} frequently interfere with Al(III) detection [27]. Thus, Al(III)-selective fluorescent probes are still in high demand. The Al(III) ion selective detection ability of Ti-MOF microflowers was evaluated towards different interference ions including Cu(II), Cd(II), Cr(III), Co(II), Ni(II), Ca(II), Mg(II), Fe(III), Fe(III), Hg(II), Mn(II), Zn(II) and NH_4^+ under similar conditions. The fluorescence intensity ratios (F/F₀) of Ti-MOF microflowers in the absence and presence of interference metal ions are shown in Figure 4.

As expected, the addition of 10 µM Al(III) resulted in over a two-fold FL enhancement compared with that of the blank. Except for Fe³⁺ ions, there was no significant change in FL intensity based on emission at 510 nm for all interference ions at 10-fold higher concentrations than Al(III). Subsequently, the fluctuation of the F/F₀ was recorded after 10 μM of Al(III) was injected into the above solution to form a competing mixture. The obtained value of the F/F_0 after the solution was added with Al(III), as shown by the green bars in Figure 4, was consistent with the control group, in which Ti-MOF microflowers were simply mixed with Al(III). Furthermore, Fe(III) ions could reduce the fluorescence of Ti-MOF microflowers, while Fe(II) ions resulted in no change in FL intensity of Ti-MOF microflowers. Thus, we employed reductant (e.g., ascorbic acid, AA) to reduce Fe(III) ions to shield the interference of Fe(III) ions towards FL detection of Al(III) ions by Ti-MOF microflowers. As illustrated in Figure S7A, the fluorescence response of Ti-MOF microflowers to Al(III) ions was unaffected by AA. Due to the reduction of Fe(III) ions by AA, the fluoresce intensity of Ti-MOF with Fe(III) ions in the presence of AA showed no significant difference compared to the control group (Figure S7B). All in all, the Al(III) induced fluorescence enhancement of Ti-MOF was not obviously influenced by the interference ions, and further illustrates that Ti-MOF microflowers display an extremely selective response to Al(III) ions.

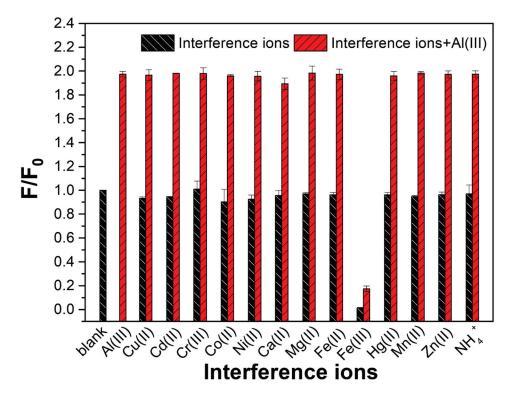


Figure 4. Ion selectivity ability of Ti-MOF microflowers towards various cations. The blank and Al(III) bar represent the Ti-MOF microflowers and the Ti-MOF microflowers with 10 μ M Al(III), respectively. The black bars represent the Ti-MOF microflowers with a variety of single interference ions (100 μ M). The red bars indicate the Ti-MOF microflowers with mixtures containing 10 μ M Al(III) and 100 μ M various interference ions. F₀ and F denote FL intensity of the blank in the absence and presence of the target ions, respectively.

3.1. Removal of Al(III) by Ti-MOF Microflowers

In addition to accurate and dependable detection, effective removal of Al(III) from polluted water is critical in wastewater treatment to reduce Al(III) accumulation below levels that pose a serious threat to humans. Time-dependent Al(III) adsorption kinetics experiments were performed to gain insight into the adsorption behavior of Ti-MOF microflowers towards Al(III), with an initial concentration of Al(III) set at 50 mgL^{-1} and a sorbent dosage

of 1 g·L $^{-1}$. To reveal the outstanding characters of the hierarchical structure, the adsorption performance of Ti-MOF microflowers was contrasted with that of pseudo-spherical Ti-MOF nanoparticles synthesized by a direct hydrothermal reaction. As illustrated in Figure 5, it is clear that the concentration of Al(III) decreased rapidly at the beginning (2 h), and then approached a constant value after a long contact time (approximately 60 h) for both types of Ti-MOF. In the case of hierarchical Ti-MOF microflowers, the adsorption of Al(III) was as high as 21.78 mg·g⁻¹ within just 12 h. With increasing treatment time, the adsorption of Al(III) on the hierarchical Ti-MOF microflowers was increased to 25.5 $\text{mg} \cdot \text{g}^{-1}$, which is higher than that of pseudo-spherical Ti-MOF nanoparticles (20.8 $\text{mg} \cdot \text{g}^{-1}$). A pseudo-second-order kinetic model was used to describe the kinetic characteristics of the adsorption sites, which was proportional to the square of the number of unoccupied sites, according to Equation (2), and as shown in Figure 5B. The parameters Q_e and k_2 for Al(III) adsorption were estimated by using linear fitting with high correlation coefficients $(R^2 = 0.998)$, as shown in Table 1. The low value of k_2 indicates that the removal rates of Al(III) by both type of Ti-MOF were fast. Furthermore, $h=k_2Q_e^2$ was calculated to quantitatively describe the initial removal rate using Equation (1). The initial removal rate (h) of Al(III) on the hierarchical Ti-MOF microflowers was larger than that of pseudo-spherical Ti-MOF nanoparticles, demonstrating a higher removal rate of Al(III) by hierarchical Ti-MOF microflowers than that of pseudo-spherical Ti-MOF nanoparticles. These results might be attributed to the distinctive hierarchical structure of Ti-MOF microflowers, the thin nanosheets reducing the diffusion resistance of Al(III) ions and allowing for faster kinetics.

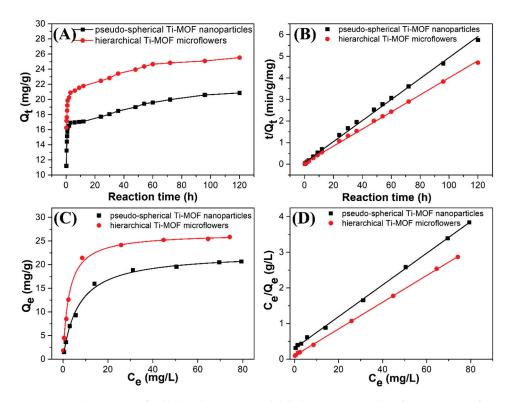


Figure 5. (**A**) Kinetics of Al(III) adsorption and (**B**) the corresponding fitting curves for pseudospherical Ti-MOF particles and hierarchical Ti-MOF microflowers via a pseudo-second-order kinetic model (50 mg·L⁻¹ of the initial concentration of Al(III), 1 g·L⁻¹ of adsorbents with initial solution pH 4.7 \pm 0.1 at 25 °C). (**C**) Isothermal adsorption of Al(III), and (**D**) the corresponding plots of C_e/Q_e versus C_e for pseudo-spherical Ti-MOF nanoparticles and hierarchical Ti-MOF microflowers (2–100 mg·L⁻¹ of initial Al(III) concentration, 1 g·L⁻¹ of adsorbents, 96 h of equilibrium time at 25 °C and initial pH 4.7 \pm 0.1).

Table 1. Adsorption parameters of Al(III) on both type of Ti-MOF.

Adsorbents	Kinetic Adsorption (Pseudo-Second-Order Model)		Isotherms (Langmuir Model)		
	K_2 $(g \cdot mg^{-1} \cdot min^{-1})$	$Q_e \ (ext{mg} \cdot ext{g}^{-1})$	$h \ (\text{mg}\cdot\text{g}^{-1} \ \cdot \text{min}^{-1})$	$q_m \; (mg {\cdot} g^{-1})$	K_L (L·mg ⁻¹)
pseudo-spherical Ti-MOF NPs	0.0324	20.86	13.67	20.66	0.1468
nierarchical Ti-MOF microflowers	0.0343	25.53	21.92	25.85	0.3831

Isothermal adsorption experiments were carried out based on kinetic adsorption behavior, and a reaction time of 96 h was used for the equilibrium adsorption time. Figure 5C shows the adsorption isotherm of Al(III) on pseudo-spherical Ti-MOF nanoparticles and hierarchical Ti-MOF microflowers by varying initial Al(III) concentrations between 2–100 mg·L $^{-1}$ under initial solution pH of 4.7 \pm 0.1. Adsorption of Al(III) increased with increasing concentrations of Al(III) in aqueous solution for both types of Ti-MOF. Table 1 shows the fitted results and calculated parameters for the Langmuir model, which was used to further analyze the Al(III) adsorption isotherms. Obviously, the Al(III) adsorption results for the two adsorbents with a high correlation coefficient, could be well fitted by the Langmuir model ($R^2 = 0.999$), which applies to monolayer adsorption on a surface with a finite number of identical active sites. Figure 5D shows the corresponding plots, which indicate good linear relationships for both types of Ti-MOF. The substantial maximum Al(III) removal capacities (q_m) were up to 20.66 mg·g⁻¹ for pseudo-spherical Ti-MOF nanoparticles and 25.85 mg \cdot g $^{-1}$ for hierarchical Ti-MOF microflowers. Based on the Langmuir model these outperformed the adsorption capacities of other adsorbents listed in Table S2. The superior adsorption capacity of hierarchical Ti-MOF microflowers towards Al(III) can be attributed to their higher specific surface area in comparison to pseudo-spherical Ti-MOF nanoparticles, as well as their unique hierarchically flower-like morphology composed of ultrathin nanosheets that provide abundant adsorption sites for Al(III) capture.

3.2. Possible Mechanism for Selective Determination of Al(III) Ions

According to previous research [28], the effect of metal cations on luminescent MOFs can be attributed to three factors: (1) the collapse of the MOFs framework, (2) cation exchange between the MOFs' central cations and the target cations, and (3) interactions between the metal cations and organic ligands. XRD and SEM measurements were used to confirm whether the crystal structure of the original framework was affected in order to further understand and elucidate the possible sensing mechanism of the phenomenon, in which Al(III)-enhances the FL of Ti-MOF microflowers. The XRD patterns (Figure S8) and SEM images (Figure S9) revealed that the sample obtained by immersing Ti-MOF microflowers in Al(III) ions was nearly identical to the original Ti-MOF, indicating that the Al(III) ions do not cause the Ti-MOF framework to collapse or change.

To verify whether the cation exchange between the central cations of Ti-MOFs and Al(III) occurred, ICP-OES was used to determine the concentrations of Ti(IV) and Al(III) in the solid MOF and supernatant after Ti-MOF treatment with various concentrations of Al(III) ions. As illustrated in Table S3, there were no Ti(IV) in all supernatants after Ti-MOF treatment with different concentration of Al(III) ions, which precluded the possibility of cation exchange between Ti(IV) and Al(III). However, Al(III) was found in all solid Ti-MOF after treatment with different concentrations of Al(III) ions. The summation of Al(III) ions in the supernatant and solid Ti-MOF were equal to the initial concentration of Al(III) ions, which is consistent with the results of the adsorption experiment. This verifies that Ti-MOF microflowers can form a complex with Al(III) ions.

Because of the interaction between metal cations and organic ligands, we propose that the Al(III)-induced fluorescence enhancement of Ti-MOF microflowers can be attributed to the chelation-enhanced fluorescence (CHEF) mechanism [29]. As previously stated, FT-IR

spectral analysis shows that the prepared Ti-MOF microflowers have phenolic hydroxyl and carboxyl groups. Metal ions can be chelated by these functional groups to form stable coordination complexes. The fluorescent selectivity of Ti-MOF microflowers towards Al(III) may be attributed to the smaller ionic radius (0.5 Å) of Al(III), which allows for a suitable coordination geometry with Ti-MOF, and a higher charge density, causing strong coordination between Al(III) and Ti-MOF. After adding Al(III), the carbonyl O, and hydroxyl O on DOBDC of Ti-MOF microflowers can coordinate with the Al(III) center atom, which increases the energies of the $n-\pi^*$ transitions compared to the corresponding π – π * transitions, and the photoinduced electron transfer process is interrupted, and the FL enhanced [30,31]. As illustrated in Figure 3A, the Al(III) ions induced a blue-shift with remarkable fluorescence enhancement of Ti-MOF microflowers. The wavelength shift indicates that the structure of DOBDC ligands in the surface of Ti-MOF microflowers may be changed in the presence of Al(III), which could result from phenolic hydroxy and carboxyl of DOBDC strongly bound with Al(III). Figure 6A shows the UV-vis absorption spectra of Ti-MOF microflowers in the present of Al(III) ions. Ti-MOF microflowers have four absorption peaks at 248, 290, 350 and 458 nm, while all of four absorption peaks are red-shifted (253, 291.5, 368.5 and 486 nm) after adding Al(III) ions. It is well known that the binding of an auxochrome to a chromophore causes increased absorption and a red shift of the chromophore. This further illustrates that Ti-MOF microflowers can form a complex with Al(III) ions.

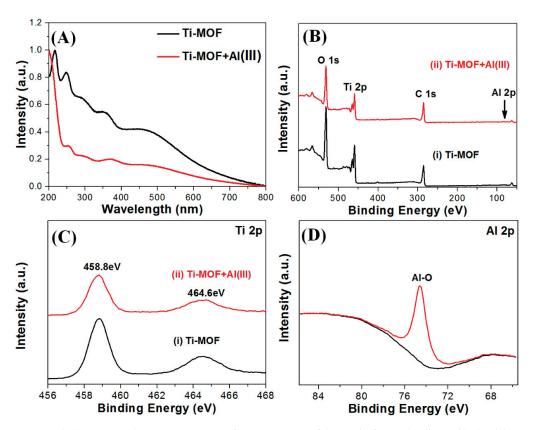


Figure 6. (A) UV-vis absorption spectra of Ti-MOF microflowers before and after Al(III) addition. (B) XPS spectra of Ti-MOF microflowers before and after Al(III) addition. (C) Ti2p spectra of Ti-MOF microflowers before and after adding Al(III). (D) Al2p spectra of Ti-MOF microflowers after addition of Al(III).

The coordination interaction between the Al(III) and the carbonyl O, and hydroxyl O on DOBDC of Ti-MOF microflowers was further validated by XPS. Figure 6B–D shows the survey spectra, Ti 2p and Al 2p spectra of Ti-MOF microflowers in the absence and present of Al(III). According to the XPS spectra, the binding energies around 286, 459 and 531 eV

corresponded to the C 1s, Ti 2p and O 1s, which are the basic elements of the Ti-MOF microflowers (Figure 6B). After the interaction with Al(III), the emergence of the Al 2p peak at 74.6 eV (Al-O) verified that Al(III) is loaded into the framework of Ti-MOF microflowers (Figure 6C).

4. Conclusions

A secondary hydrothermal method was used to successfully prepare three-dimensional hierarchical titanium-based metal-organic framework (Ti-MOF) microflowers for the simultaneous detection and removal of Al(III) ions in aqueous solution. The as-prepared Ti-MOF microflowers are made up of flowerlike micro/nanostructured particles with sizes ranging from 5 to 6 μ m. The particles are made of 200 nm nanorods with a length of 1–2 μ m and a porous structure with a pore diameter of 1.8 nm and a specific surface area of 264.47 m 2 ·g $^{-1}$, which provides abundant surface active sites for detection and adsorption of Al(III) ions. The resultant Ti-MOF microflowers exhibit excellent stability, high water dispensability, and dramatical fluorescence characteristic, making them useful as fluorescence probes for quantitative detection of Al(III) in aqueous solution. Importantly, specific FL enhancement by Al(III) via the chelation-enhanced fluorescence (CHEF) mechanism can be utilized to selectively and accurately determine Al(III) due to the coordination interaction between the Al(III) and the carbonyl O, and hydroxyl O on DOBDC of Ti-MOF microflowers. Linear concentration ranges of up to 15 µM and a detection limit of 75 nM are easily achievable, the latter being approximately 100-fold lower than the WHO permitted level of aluminum in drinking water (7.4 μM). By introducing ascorbic acid, the interference of Fe(III) can be shielded to achieve selective detection of Al(III) in the presence of various cations. Ti-MOF microflowers have an excellent adsorption capacity for Al(III), with a maximum adsorption capacity of 25.85 mg·g $^{-1}$. The adsorption rate is rapid and corresponds to a pseudo-secondorder kinetic model. The Ti-MOF is an excellent candidate as a fluorescent chemical sensor and adsorbent for aqueous contaminants due to its impressive water stability, high porosity, and strong luminescence.

Supplementary Materials: The following supporting information can be downloaded at: https://www.mdpi.com/article/10.3390/bios12110935/s1, Figure S1: HRTEM images of hierarchical Ti-MOF nanoflowers through secondary hydrothermal synthesis; Figure S2: SEM image of pseudospherical Ti-MOF by direct hydrothermal reaction; Figure S3: Photostability of Ti-MOF nanoflowers under UV lamp irradiation; Figure S4: Relative fluorescence intensity (F/F₀) of Ti-MOF nanoflowers at different ionic strengths (0–1000 mM NaNO₃) and pH values (4.0–9.0); Figure S5: Fluorescence response—time profile of Ti-MOF nanoflowers before and after adding 10 μ M Al(III); Figure S6: Job's plot for stoichiometric determination of Ti-MOF and Al(III); Figure S7: Fluorescence spectra of Ti-MOF in the presence of Al³⁺ and AA, and Fluorescence spectra of Ti-MOF in the presence of Fe²⁺, Fe³⁺ and AA; Figure S8: PXRD pattern of Ti-MOF nanoflowers after treatment with Al³⁺ ions; Figure S9: SEM image of Ti-MOF nanoflowers after treatment with Al³⁺ ions; Table S1: Comparison of different fluorescent material sensors for Al(III) determination [32–44]; Table S2: Comparison of the Al(III) adsorption capacities for different adsorbent materials [35,45–51]; Table S3: ICP-OES analysis of the content of Ti⁴⁺ and Al³⁺ in the solid MOF and supernatant after Ti-MOF treatment with different concentration of Al³⁺ ions.

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Article

Terahertz Fingerprint Metasurface Sensor Based on Temperature Variation for Trace Molecules

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Abstract: Terahertz (THz) spectroscopy has demonstrated significant potential for substance detection due to its low destructiveness and due to the abundance of molecular fingerprint absorption signatures that it contains. However, there is limited research on the fingerprint detection of substances at different temperatures. Here, we propose a THz metamaterial slit array sensor that exploits localized surface plasmons to enhance the electric field within the slit. The transmission peak frequency can be modulated via temperature adjustments. This method enables the detection of molecular absorption characteristics at multiple spectral frequency points, thereby achieving a specific and highly sensitive detection of characteristic analyte fingerprint spectra. Additionally, the sensor supports the detection of substances at multiple temperatures and sensitively identifies changes in their absorption properties as a function of temperature. Our research has employed temperature variation to achieve a highly sensitive and specific detection of trace analytes, offering a new solution for THz molecular detection.

Keywords: Terahertz; metasurface; sensing; temperature variation; plasmons

1. Introduction

Terahertz (THz) waves possess ultra-low photon energy, which eliminates the risk of causing damaging ionization in analytes [1]. Many chemical molecules exhibit resonance and rotational frequencies within the THz range, resulting in distinct absorption resonances and intricate characteristic fingerprint spectra. Analyzing the interaction between incident THz waves and molecules provides valuable insights into molecular configuration, conformation, and environmental effects [2], thus enabling specific detection capabilities [3]. THz sensing holds significant promise for the non-destructive, rapid, and precise identification of various chemical molecules, with wide-ranging applications spanning across physics, chemistry, biomedicine, and related fields [4]. Nevertheless, a significant challenge arises from the mismatch between the THz wavelength, typically at the micrometer scale, and the absorption cross-section of molecules, which operates at the nanometer scale [5]. Most molecules exhibit weak interactions with THz waves. Traditional THz spectroscopy methods typically require large analyte quantities to observe their characteristic finger-print spectra, leading to low sensing sensitivity [6,7]. Enhancing the sensing sensitivity of analytes in practical applications has become imperative.

The application and sensing of high-energy explosive has been one of the international research hotspots. Hexogen (RDX), also known as cyclonite, is a powerful military explosive. It is a white crystalline powder and is widely used in various fields such as military, aerospace, mining, and chemical industries due to its high explosive speed and power [8]. It is almost impossible to detect through X-ray imaging, necessitating the development of new non-contact and non-destructive techniques for effective detection. THz waves, due to their non-destructive nature, offer unique advantages for RDX detection. The broad frequency range and weak absorption of RDX fingerprints limit the practical applications of THz sensing. Enhancing the sensitivity of THz sensing for RDX is the primary focus of

current research. HNIW (CL-20), chemically known as 2,4,6,8,10,12-hexanitro-2,4,6,8,10,12-hexaazaisowurtzitane, is a highly energetic and potent non-nuclear elementary explosive. It has wide-ranging applications in military and civilian sectors, garnering significant attention from researchers. CL-20 undergoes an irreversible ε – γ phase transition induced by temperature, a process that is closely linked to its explosive behavior. The precise detection of CL-20 and its phase transition are crucial to ensure its safe utilization [9,10]. Unfortunately, achieving a rapid and highly sensitive sensing of CL-20 and its phase transition process remains challenging. Specifically, detecting fingerprints at different temperatures before and after the ε – γ phase transition of CL-20 presents significant difficulties.

In previous studies, researchers have proposed various artificially designed metasurface structures, including open-ring resonators [11], metal gratings [12], photonic crystal structures [13], waveguide configurations [14], and structures based on bound states in the continuum [15,16], to achieve heightened sensitivity in material detection [17–19]. The majority of these structures are limited to detecting material absorption characteristics at individual frequency points via resonance peaks. They cannot accurately analyze the absorption spectra curve within a specific frequency range to identify particular substances [20]. In recent years, researchers have employed strategies such as structural scanning, angle scanning, stretchable flexible materials, and adjustable graphene to modulate the resonance peak frequency [21]. This advancement allows for the detection of characteristic absorption curves within a defined frequency band, facilitating material analysis and specific substance identification [22,23]. Nonetheless, limitations persist, including narrow sensing frequency bands and the inconvenience of frequency modulation strategies. Current research on the spectral response and sensing capabilities of metasurfaces at varying temperatures remains inadequate.

In this paper, we propose a THz metamaterial nano-slit array sensor based on temperature variation to enhance the detection of molecular fingerprint spectra. The sensor utilizes the temperature-sensitive semiconductor material InSb to form the nano-slit array structure, successfully generating localized surface plasmons (LSPs) [24]. When a surface plasmon is confined to a particle of a size comparable with the wavelength of light, the particle's free electrons participate in the collective oscillation, and it is termed an LSP. LSPs significantly enhance the electric field near the particle's surface, with the maximum enhancement at the surface rapidly diminishing with distance. The particle's optical extinction peaks at the plasmon resonance frequency. This extinction peak depends on the refractive index of the surrounding medium [25]. The sensor achieves localized electric field enhancement based on LSP, significantly improving sensing sensitivity [26,27]. InSb is a frequently utilized temperature-sensitive material. As the temperature increases, both the real and imaginary components of the InSb dielectric constant in the terahertz band significantly and continuously increase. This leads to alterations in the metasurface resonance frequency [28]. When the temperature varies from 275 K to 520 K, the resonance peak frequency gradually shifts from 0.78 THz to 1.6 THz, with a change in amplitude of 0.82 THz. By continuously changing temperatures, the resonance peak can be located at different frequency points. This allows for the sensitive sensing of molecular absorption features at multiple frequencies and the formation of an envelope curve depicting characteristic absorption spectra. Taking RDX detection as an example, this sensor has a detection limit of 1.61 μ g/cm². The proposed sensor can specifically identify trace substances. It provides a reliable approach for trace molecule sensing using THz technology [29]. In addition, the sensor can operate at multiple temperatures. By altering the slit length to modify the resonance frequency at a constant temperature, characteristic absorption spectra at multiple temperatures can be detected. By coating the sensor surface with CL-20, distinct characteristic absorption curves of CL-20 were detected at 298 K and 453 K. This enables the observation of changes in CL-20 absorption features before and after the ε - γ phase transition. The sensor can detect changes in molecular fingerprint spectra with temperature. This advances the application of THz sensing in monitoring temperature-dependent absorption characteristics of substances.

2. Structure and Design

Figure 1a illustrates the structural diagram of the THz metamaterial nano-slit array sensor. THz waves, polarized along the x-axis, propagate vertically into the metamaterial. The substrate material utilized is silicon dioxide. The relative permittivity of SiO₂ shows no significant differences at various operating temperatures. We approximate that the relative permittivity of SiO₂ remains constant at 1.73 across all operating temperatures [30–32]. The nano-slit array structure comprises the temperature-sensitive InSb, whose dielectric constant can be approximated by the Drude model as [33].

$$\varepsilon(\omega) = \varepsilon_{\infty} - \frac{\omega_{\rm p}^2}{\omega^2 + i\gamma\omega} \tag{1}$$

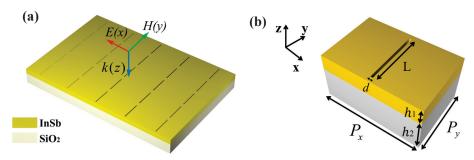


Figure 1. (a) Structure diagram of the metasurface sensor. (b) Metasurface cell and its structural parameters.

In Equation (1), ε_{∞} = 15.68 denotes the dielectric constant at infinite angular frequency, ω represents the angular frequency, $\omega_p = \sqrt{Ne^2/\varepsilon_0 m^*}$ denotes the plasma frequency, $\gamma = 0.1\pi$ THz signifies the damped vibration frequency, N denotes the intrinsic carrier concentration, ε_0 denotes the vacuum dielectric constant, e represents the electron charge, and $m^*=0.015m_e$ (m_e denotes the free electron mass) represents the effective mass of the free carriers [34].

N can be expressed by the following relation [35]:

$$N = 5.76 \times 10^{20} T^{1.5} exp(-0.26/2k_B T)$$
 (2)

In Equation (2), k_B denotes the Boltzmann constant, and T represents the temperature in Kelvin. Hence, the intrinsic carrier concentration N of InSb is influenced by the external temperature, subsequently affecting its dielectric constant. Therefore, InSb serves as a temperature-sensitive semiconductor material, facilitating the convenient modulation of the sensor's resonance peak via changes in the external environment, ultimately leading to the formation of a broadband transmission envelope.

Figure 1b depicts the unit cell of the metamaterial. The geometric parameters of the unit cell are set as follows: period P_x = 120 μ m and P_y = 80 μ m; slit width d = 3 μ m and length L = 68 μ m; thickness of InSb and slit h₁= 20 μ m; and thickness of the SiO₂ substrate h₂ = 40 μ m. To investigate the performance of the sensor, the optical characteristics of the metamaterial are simulated using three-dimensional finite-difference time-domain (FDTD) software. In the simulation, perfect matching layers are applied along the z-direction, while periodic boundary conditions are applied along the x- and y-directions.

3. Result and Discussions

Figure 2a depicts the transmission spectrum of the proposed metasurface sensor when THz waves, polarized along the x-axis, vertically impinge upon the metasurface at an ambient temperature of 310 K. As observed from Figure 2a, a prominent resonance transmission peak is evident at 1 THz. Meanwhile, Figure 2b illustrates the distribution of electric and magnetic fields at the resonance frequency. The electric field (x-z plane) is

entirely localized at the slit and is significantly enhanced. The magnetic field is localized at the ends of the slit [36]. Such a field distribution markedly enhances the interaction between THz waves and analytes, thereby enhancing the sensitivity of the sensor [37].

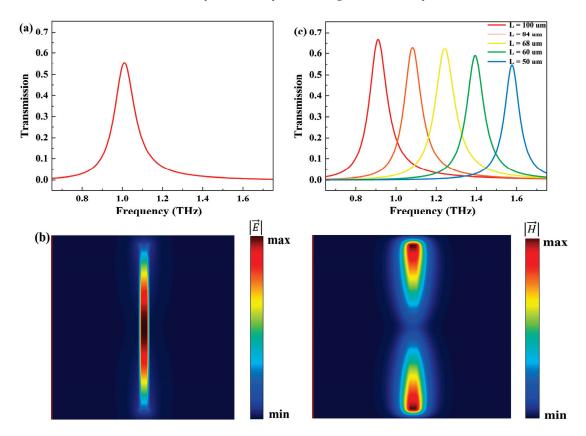


Figure 2. (a) Transmission spectrum of the metasurface at T = 310 K. (b) Electric and magnetic field distributions at f = 1 THz. (c) Transmission spectra for different slit lengths L at T = 360 K.

To comprehensively understand the impact of slit length L on the resonance transmission peak, the transmission spectra of sensors with different L ranging from 50 μ m to 100 μ m were analyzed. Simultaneously, while adjusting the slit length L, the distance between slits (the period P_y) was adjusted to prevent adjacent slits from contacting each other. Specifically, when L = 50, 60, 68, 84, 100 μ m, the corresponding periods P_y were 70, 80, 80, 110, and 120 μ m, respectively. As depicted in Figure 2c, with increasing L, the resonance peak frequency gradually decreases, while the transmission peak value slightly increases [38]. The slit length L can influence the resonance frequency and alter the position of the resonance peak. The change in transmission peak value can be explained by the dielectric properties of InSb. As the frequency decreases, the extinction coefficient k of InSb gradually increases. This significantly suppresses THz transmission, thereby reducing the peak transmission.

To characterize the spectral response of the sensor, the relationship between the structural transmission curve and temperature T was investigated. Figure 3a illustrates the variation of the sensor's transmission spectra with temperature T ranging from 275 K to 520 K. As the temperature increases, the resonance angle experiences a blueshift, and the transmission peak value continuously increases. Under slit length $L=68~\mu m$, the resonance peak can shift from 0.78 THz to 1.6 THz, achieving an ultra-wideband shift of 0.82 THz. Figure 3b demonstrates the relationship between the dielectric constant of InSb and temperature, explaining the variation in resonance frequency and transmission peak value [39]. Figure 3c shows the electric field distribution of the metasurface at 1 THz from 300 K to 330 K. The electric field within the slits is significantly enhanced, with the maximum electric field intensity occurring at 310 K. Both increasing and decreasing the

temperature lead to a reduction in electric field intensity, corresponding to changes in transmittance at 1 THz for the metasurface. We selected a room temperature range of 260 K to 340 K for molecular fingerprint sensing [40]. Figure 3d presents the transmission spectra (from 260 K to 340 K, with a step size of 2 K) and corresponding metasurface envelope curves. The resonance angles of the transmission spectra at different temperatures were extracted and further fitted into the envelope curves through interpolation (as shown by the red curve in the figure) [41]. With a temperature variation of only 80 K, this envelope curve can cover a broad frequency range from 0.68 THz to 1.16 THz, which is advantageous for the characteristic fingerprint sensing of substances. In this scenario, when the analyte is placed at the slit of the metasurface, an absorption peak matching the fingerprint of the analyte will be observed in the envelope curve [42,43].

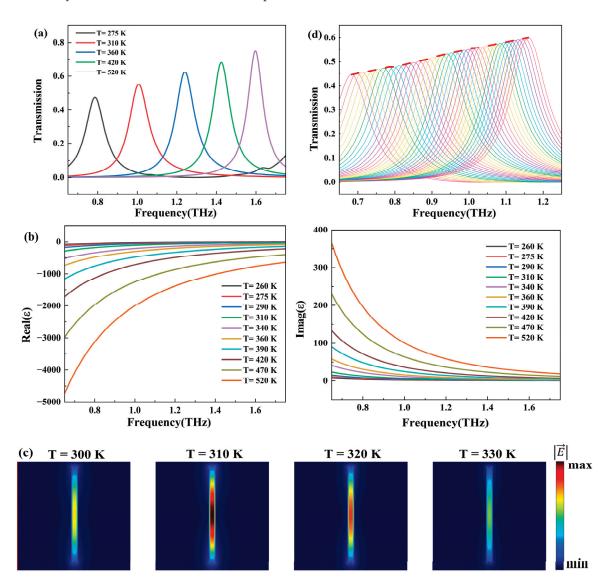


Figure 3. (a) Transmission spectra of the metasurface within the ambient temperature range of 275 K to 520 K. (b) Dielectric constant of InSb at different ambient temperatures. (c) The electric field distribution of the metasurface at 1 THz from 300 K to 330 K. (d) Transmission spectra of the metasurface and its envelope curves (the red curve) within the ambient temperature range of 260 K to 340 K (increment of 2 K).

To study its sensing performance, we covered the sensor surface with RDX. Based on the existing THz time-domain spectroscopy analysis results of RDX, we obtained the refractive index n and extinction coefficient k of RDX in the THz range [44], as shown in

Figure 4a. Due to weak intermolecular interactions and lattice vibrations, RDX exhibits a weak absorption peak at 0.88 THz [45]. Additionally, the absorption peak of RDX is relatively broad, necessitating a sufficiently wide sensing frequency bandwidth. Figure 4b illustrates the transmission spectra of the metasurface covered with 8.03 µg/cm² RDX as a function of temperature (260-340 K) and the corresponding envelope curves. Because RDX exhibits excellent chemical stability at room temperature and can stably exist for long periods below 400 K, detection at temperatures ranging from 260 K to 340 K is reasonable. With increasing temperature, a significant blueshift in the resonance angle is observed, and a distinct absorption peak appears at 0.88 THz in all spectra. The maximum transmission rate at the same frequency point for different temperatures was extracted, and the results were further fitted into envelope curves through interpolation. It can be observed that the envelope curve exhibits the minimum transmission rate at 0.88 THz and forms an envelope peak at this frequency, corresponding to the fingerprint spectrum of RDX [46]. Let the envelope curve of the sensor without RDX coverage be denoted as T0 and the envelope curve of the sensor covered with 8.03 µg/cm² RDX be denoted as Ts. The change in transmission rate was normalized to Tn = (T0 - Ts)/T0. The normalized envelope curve Tn shown in Figure 4c closely follows the trend of the extinction coefficient k of RDX. This indicates that our proposed nano-slit array sensor can accurately amplify and detect the characteristic absorption fingerprint of RDX [47,48], achieving a specific and highly sensitive sensing of RDX.

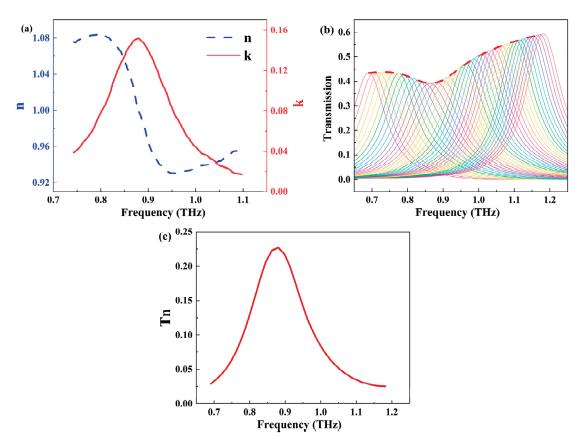


Figure 4. (a) Refractive index n and extinction coefficient k of RDX in the THz range. (b) Transmission spectra of the metasurface covered with 8.03 μ g/cm² RDX as a function of ambient temperature (260–340 K) and the corresponding envelope curves (the red curve). (c) Envelope curves of the transmission spectra of the metasurface after normalization.

To further investigate the system's capability to detect trace analytes, we covered the metasurface with different concentrations of RDX. Figure 5a illustrates the normalized envelope curves deposited with varying concentrations of RDX on the metasurface. As

the RDX concentration increases from 1.61 $\mu g/cm^2$ to 40.16 $\mu g/cm^2$, the corresponding normalized transmission rate change also gradually increases. Figure 5b demonstrates the relationship between the normalized transmission rate change at 0.88 THz and the RDX concentration in the range of 1.61 $\mu g/cm^2$ to 40.16 $\mu g/cm^2$. The linear fitting equation of the curve in the figure is approximately y=0.01519x+0.08182, where y represents the normalized transmission rate change at 0.88 THz and x denotes the RDX concentration in $\mu g/cm^2$. The correlation coefficient of this fitting curve is 0.98, indicating a good linear correlation between the normalized transmission rate change at 0.88 THz and the RDX concentration. In our experiments, the minimum RDX concentration featured a clearly observable envelope absorption peak, and the detection limit of our sensor was 1.61 $\mu g/cm^2$. These studies suggest that RDX content can be detected based on the normalized transmission rate change of the envelope peak, enabling a quantitative sensing of RDX with a very low detection limit.

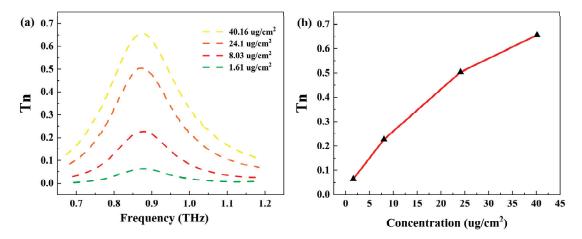


Figure 5. (a) Normalized envelope curves deposited with different concentrations of RDX on the metasurface. (b) Relationship between the normalized transmission rate change at 0.88 THz and the RDX concentration.

In addition, the sensor can track the characteristic absorption spectra changes of CL-20 before and after the temperature-induced phase transition. Previous studies have demonstrated that CL-20 can undergo an irreversible ε-γ phase transition triggered by temperature variation. At 298 K, the ε -CL-20 phase exhibits two absorption peaks at 1.31 THz and 0.99 THz. As the temperature increases to 453 K, CL-20 transitions into the γ phase. The characteristic absorption peaks at 1.31 THz and 0.99 THz diminish, while a strong absorption occurs at 1.53 THz [9,10]. When the ambient temperature of the metasurface is maintained at 298 K, altering the slit length L results in a redshift in resonance frequency and a slight increase in transmission peak value. Different slit lengths are employed across various regions of the metasurface to encompass the resonance inclination angles of different frequency positions, forming corresponding envelope curves [22]. Figure 6a illustrates the transmission spectra and their envelope curve (depicted in red) for slit lengths ranging from 100 μ m to 36 μ m. This envelope curve spans the 0.72–1.45 THz range. As the slit length L varies, the period P_{ν} changes with it, ensuring that the distance between the edges of two adjacent slits in the y-direction remains at 40 μm. Figure 6b displays the transmission spectra and their envelope curve (depicted in red) following the introduction of trace CL-20. Notably, distinct absorption peaks emerge at 1.31 THz and 0.99 THz, aligning with the position of the characteristic absorption peak of ε -CL-20. Figure 6c presents the normalized envelope curve, effectively describing the absorption characteristics of ε-CL-20 at 298 K. As the ambient temperature rises to 453 K, Figure 6d showcases the transmission spectra and their envelope curve, with only the transmission spectra from 1 THz to 1.8 THz being displayed, omitting the formant at higher frequency positions. Figure 6e exhibits the transmission spectra and its envelope curve following the introduction of an equal content of CL-20, while Figure 6f represents the normalized envelope curve. At 453 K, the envelope curve exhibits strong absorption at 1.53 THz, while the absorption peak at 1.31 THz disappears. This is consistent with the absorption characteristics of γ -CL-20, effectively describing the absorption features of CL-20 at 453 K. The comparative analysis of the transmission spectra and envelope curves at 298 K and 453 K, covering trace amounts of CL-20, shows a shift in the characteristic absorption peak frequencies from 0.99 THz and 1.31 THz to 1.53 THz. This allows for a clear observation of CL-20's absorption characteristics before and after the phase transition, facilitating the sensing and study of its phase transition mechanism. These findings validate the system's capability to perform highly sensitive sensor detection of material feature absorption fingerprints at multiple temperatures and track changes in material feature absorption fingerprints with temperature.

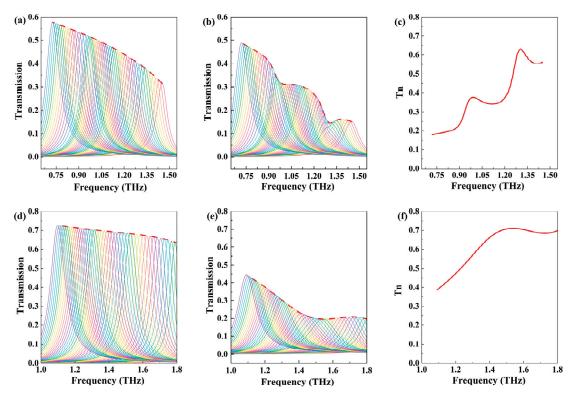


Figure 6. (a) Metasurface transmittance spectra and envelope curve of slit length L from 36 μ m to 100 μ m at 298 K. (b) Metasurface transmittance spectra and envelope curves after covering trace CL-20 at 298 K. (c) The envelope curve after normalization at 298 K. (d) Metasurface transmittance spectra and envelope curve at 453 K. (e) Metasurface transmittance spectra and envelope curves after covering trace CL-20 at 453 K. (f) The envelope curve after normalization at 453 K.

4. Conclusions

In summary, we propose a nano-slit array sensor based on temperature variation that is capable of a highly sensitive and specific sensing of molecular fingerprints. By integrating structural parameter scanning, this sensor allows for the detection of substance characteristics at various temperatures. The metasurface consists of periodically arranged InSb slits, with SiO₂ as the substrate. When THz waves are vertically incident on the metasurface, LSPs are successfully excited, leading to the localized enhancement of the electric field at the slits. This enhanced electric field significantly amplifies the interaction between the analyte and THz waves, effectively improving sensing sensitivity. With the variation in temperature, the dielectric properties of the temperature-sensitive InSb continuously vary, causing the transmission resonance angle to shift, thereby generating transmission envelope curves covering a wide frequency range. Based on the temperature-scanning strategy, we successfully delineate the characteristic fingerprint spectra of RDX,

enabling qualitative and quantitative sensing of RDX with a detection limit of 1.61 $\mu g/cm^2$. Simultaneously, employing the structural parameter scanning method and adjusting the slit length L, we successfully depicted the characteristic fingerprint spectra of ϵ -CL-20 at 298 K and γ -CL-20 at 453 K. This approach achieved a high sensitivity detection of CL-20 absorption features at different temperatures, allowing for the observation of the variation in CL-20 ϵ - γ phase transition absorption characteristics. Our research indicates that the proposed temperature-scanning metasurface sensor holds significant potential for the wideband fingerprint sensing of trace analytes, thus promoting the development of THz sensing in various fields, including civilian and military security, biomedicine, and beyond.

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Article

Fiber-Optic Distributed Sensing Network for Thermal Mapping of Gold Nanoparticles-Mediated Radiofrequency Ablation

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Abstract: In this work, we report the design of an optical fiber distributed sensing network for the 2-dimensional (2D) in situ thermal mapping of advanced methods for radiofrequency thermal ablation. The sensing system is based on six high-scattering MgO-doped optical fibers, interleaved by a scattering-level spatial multiplexing approach that allows simultaneous detection of each fiber location, in a 40×20 mm grid (7.8 mm² pixel size). Radiofrequency ablation (RFA) was performed on bovine phantom, using a pristine approach and methods mediated by agarose and gold nanoparticles in order to enhance the ablation properties. The 2D sensors allow the detection of spatiotemporal patterns, evaluating the heating properties and investigating the repeatability. We observe that agarose-based ablation yields the widest ablated area in the best-case scenario, while gold nanoparticles-mediated ablation provides the best trade-off between the ablated area (53.0–65.1 mm², 61.5 mm² mean value) and repeatability.

Keywords: radiofrequency ablation; fiber-optic shape sensors; optical fiber sensor; distributed sensors; gold nanoparticles; biomedical sensors

1. Introduction

The role of minimally invasive cancer thermotherapies is substantially extended in the latest years, as clinical trends show strong progress toward methods with limited invasiveness and consistent repeatability in the replacement of more invasive methods such as surgical resection [1–3]. Methods based on thermal ablation find consistent progress as the technological advances in miniaturization of the devices [4], sensing [5], and real-time imaging [6] allow for achieving substantial clinical improvements while maintaining minimal invasiveness and percutaneous insertion of the surgical device.

Thermal ablation procedures accomplish a successful cancer treatment by transferring energy in the form of electromagnetic waves from a source placed outside of the patient into an applicator, percutaneously inserted into the tumor location [7] or externally placed [8]; the energy dissipated into the applicator is converted into thermal energy, resulting in a localized rise of temperature within the tumoral tissue that spreads from the applicator to the peripheral side. Temperature values over 60 $^{\circ}$ C result in a nearly instantaneous mortality of cancer cells, while cytotoxic effects are recorded for temperature values over 42 $^{\circ}$ C [9]; the overall effect is evaluated by the thermal dosimetry [10,11], which integrates the temperature over the exposure time.

Thermal ablation methods differ on the frequency f of the electromagnetic waves, and consequently on the type of applicator involved in the energy delivery [5]. The four main methods that involve thermal heating are based on: (1) radiofrequency ablation (RFA, f~450 kHz), which uses an electrical RF generator and a miniature applicator with an electrode mounted on a single tip [12] or multiple tips [13]; (2) high-intensity focused ultrasound (HIFU, f~1.2 MHz), which uses an ultrasound generator and an external transducer that focuses the incoming waves into the tissue, with a non-contact applicator that ablates the tissue by a combination of thermal effects and cavitation [14]; (3) microwave ablation (MWA, f~2.4 GHz), which uses a microwave generator and an applicator shaped as a transmission line and operating as a near-field antenna [15]; and (4) laser ablation (LA, f~300 THz), which uses a mid-power laser source either externally firing into the tissue surface [9] or coupled into a large-core optical fiber for in situ delivery [16].

Among these procedures, RFA is highly versatile, and achieves growing success rates in the treatment of solid tumors such as hepatocellular carcinoma [2,17], and spinal tumors [18]. An RFA needle can integrate multiple miniature electrodes [19], or a single electrode with a high contact surface for rapid thermotherapies [20]; some applicators incorporate micro-thermocouples for sensing [21]. Outside of cancer care, RFA can be applied in the treatment of cardiac arrhythmia [22], and in interventional pain management [23,24].

The main research trends aimed at the technological improvement of RFA point in two different directions. The first area of interest involves the use of the advances in material sciences to improve the heating efficacy. A critical factor for the cancer treatment via RFA relies on the changes of electrical impedance of the tissues as the temperature approaches the 100 °C value [5]; the vaporization of the inner part of the tissue causes the impedance to abruptly rise, and therefore only a small portion of the electrical power is dissipated onto the tissue. Under this condition, medical generators enter a "safe mode", discontinuing the power supply and interrupting the RFA procedure [25]. In order to improve over this effect, both agarose and chitosan gels [26] have been employed in order to reduce the tissue impedance. In addition, gold and silver nanomaterials [27] have been used in combination with agarose with the purpose of improving the thermal delivery, and ultimately increasing the ablated region.

The second area of interest is the use of advanced biosensors for the detection of physical parameters in situ prior to the ablation (for diagnostic purposes) and during the thermal treatment. Optical fiber sensors serve this purpose much better than electrical or mechanical sensors [28], as they can incorporate biological [29] and biophysical sensing [30], as well as owning key properties in terms of biocompatibility, miniaturization of the footprint, disposability, response time, and spatial distribution.

Several optical fiber biosensors have been reported for this purpose. Evers et al. [31] reported fiber-optic biosensors for liver tissue identification, discriminating the healthy and tumoral tissues with a spectroscopic probe. Loyez et al. [29] reported a plasmonic tilted fiber Bragg grating with the capability of in situ detection of cytokeratin biomarkers for the identification of tumoral cells for following treatments. Tosi et al. [32] reported a fiber-optic dual pressure and temperature sensor for measuring the pressure in the proximity of the RFA tip during an ex vivo procedure. Several authors also reported the use of fiber Bragg gratings (FBGs) [12,33] and distributed sensors [34,35], for the temperature detection during hyperthermia, exploiting the spatial resolution of optical fiber sensors.

From an application standpoint, and from the point of view of controlling the thermal ablation process in real-time [36], temperature sensing plays the largest role in measuring the effectiveness of the procedure since thermal damage is largely depending upon the instantaneous temperature and also the variability of the electrical and thermal properties of the tissues prevents from obtaining ablation patterns with high repeatability in the experimental conditions. Densely arrayed FBG sensors [37], chirped FBGs [30], long period gratings [38,39], and distributed sensors, particularly when arranged into a multi-fiber grid-shaped arrangement [34], have the possibility of sub-centimeter spatially resolved thermal sensing. However, while distributed sensors interrogated via optical backscatter

reflectometry [33] can use inexpensive single-mode fibers or high scattering fibers as sensors without any additional manufacturing, FBG sensors are more expensive as they require inscription of multiple devices into the fiber, making the system less compatible with a disposable use.

Prior works were performed on RFA using ferromagnetic nanoparticles [26], and silver nanoparticles with a green-oriented synthesis method [40], showing an increase of the treated region when nanoparticles are inserted in situ in a solution of agarose or chitosan gel. However, gold nanoparticles (AuNP) represent the most popular method to extend the performance of thermotherapies, as AuNP have been demonstrated in radiofrequency [41] and laser [42] ablation. AuNP have been investigated as a method for improving the thermal treatments, since they combine their properties of ease of synthesis and biocompatibility [43] with their advantageous electrothermal effects: lowering the impedance of the tissue at the electrode contact point, hence extending the duration of RFA [44], and improving the heating process at the peripheral side of the tumor [45,46], hence targeting a wider region. In addition, AuNP improves the drug delivery systems within cancer tissues, as shown in previous works [27,47,48].

In this work, we consolidate the design of a distributed fiber-optic sensing network for real-time, mini-invasive, and spatially resolved thermal detection for AuNP-mediated RFA; results are presented by comparing the thermal response of pristine ablation with AuNP-mediated ablation with two different density levels (1 and 4 mg/mL), and agarose-mediated RFA. The biosensing system is based on a network of six optical fibers having high scattering, arranged in a spatial division multiplexing setup. The positioning of the fibers in the tissue allows sensing over a 40×20 mm grid, with 7.8 mm² pixel size. RFA experiments have been performed ex vivo on bovine phantom, in pristine mode, mediated by agarose, and finally mediated by Au nanoparticles in different densities. The resulting 2-dimensional (2D) thermal maps allow for the recording of the different heating patterns and spatiotemporal trends, and investigate the efficacy of each ablation process.

2. Materials and Methods

2.1. Experimental Setup

The experimental setup of the thermal ablation procedure shown in Figure 1 comprises the following parts: (a) optical backscatter reflectometer (OBR, Luna 4600, Roanoke, VA, USA) with a computer used to collect and process the data; (b) distributed MgO nanoparticle-doped optical fibers spliced to single-mode fibers; (c) commercially obtained bovine liver; (d) RF applicator, with cylindrical shape (160 mm length, 3 mm diameter) and an active brass electrode on the conical tip of ~10 mm length; and (e) RF/MWA Hybrid Generator (LEANFA S.r.l., Ruvo di Puglia, Italy) that launches a 450-kHz radiofrequency signal to the applicator. The OBR is used in distributed sensing mode, measuring the frequency shift of each fiber signature within the network with millimeter-level spatial resolution [49]. Thermal ablation experiments employed six separate high-scattering optical fibers with different lengths plugged with the OBR Luna 4600 to measure temperature change during radiofrequency ablation.

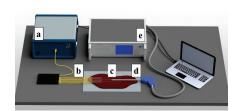




Figure 1. Schematic (**left**) and photographic view (**right**) of the radiofrequency ablation setup of the parenchymal tissue. The setup consists of: (a) OBR Luna 4600 with the computer used for data acquisition, (b) optical fibers, (c) bovine liver, (d) the RFA applicator, and (e) a hybrid RF/MWA generator used in RF mode.

2.2. Fiber Calibration

The presented work applied six Mg-silicate NPs fibers tagged M01 containing erbium in the core. The fabrication of M01 fiber includes solvents such as erbium (III) chloride hexahydrate and magnesium chloride in volumes of 10-4~mol/L and 0.1~mol/L, respectively [50]. Six MgO nanoparticle-doped optical fibers with the core diameter of $10~\mu m$ and the cladding diameter of $125~\mu m$, matching the size and compound of single-mode glass fibers, have been used for the detection of the temperature change. Figure 2 shows the photographs of the used fibers.

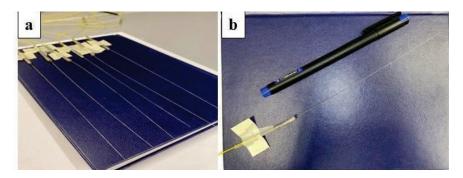


Figure 2. Photographic view of the MgO-doped fibers employed for temperature sensing. (a) Photos of the fibers displayed on a grid; (b) view of a single nanoparticle-doped optical fiber.

We report in Figure 3 the thermal calibration of the fiber, obtained by exposing the fiber in a water bath heated by a thermal plate (IKA magnetic stirrer hot-plate, IKA-Werke Gmbh, Staufen, Germany), at temperatures ranging from 21 to 57 °C, and recording the temperature with a commercial fiber Bragg grating (Technica Optical Components LLC, Beijing, China) interrogated by an FBG (Fiber Bragg Grating) interrogator (si255 model, Micron Optics/Luna, Roanoke, VA, United States). The thermal coefficient is estimated as 9.18 pm/°C, which is very similar to FBG sensors working at 1550 nm (~10 pm/°C [32]).

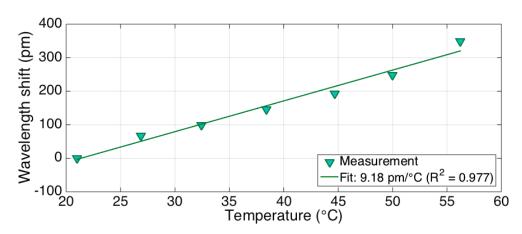


Figure 3. Calibration of the temperature coefficient of the optical fibers used for sensing.

2.3. Gold Nanoparticle Preparation

The presented study uses AuNPs of 15–20 nm size and in different densities introduced directly into the phantom. The synthesis of AuNPs was conducted using the citrate reduction method as presented by Turkevic et al. [51]. The 20 nm AuNPs were achieved by adding 2.0 mL of 34 mM trisodium citrate solution into a boiling solution, containing 0.5 mL of 1% hydrogen tetrachloroaurate (III) trihydrate and 50 mL of deionized water. The color of the solvent changed from bright yellow to dark violet in a few minutes and turned to ruby red after 15 min of stirring. The obtained solution was cooled at room temperature for the next 20 min and then cleaned by deionized water using a centrifuge operating at

15,000 rpm in 2-mL tubes (Figure 4). The size and shape of the synthesized GNPs were characterized using transmission electron microscopy (TEM) (Figure 5).

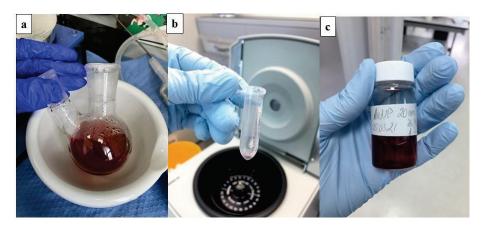


Figure 4. Gold nanoparticle cooling process (a) and after centrifuge (b,c).

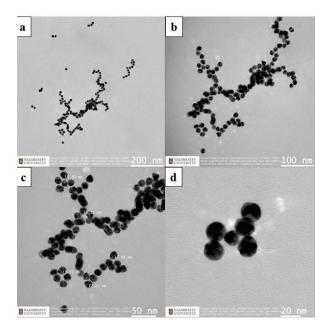


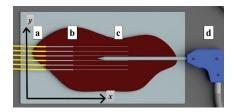
Figure 5. TEM images of AuNPs in different zooms: 200 nm (**a**), 100 nm (**b**), 50 nm (**c**), and 20 nm (**d**). The shape and size of AuNPs equivalent to 15–20 nm spheres in different zooms.

Agarose gel was prepared by mixing 1 g of agarose powder with Tris Acetate EDTA (TAE buffer) in a 50 mL volume and heating in the microwave until it dissolved completely. In order to obtain different densities of nanoparticles starting from 1 mg/mL to 4 mg/mL, AuNPs were dissolved in 0.2% agarose solution at a 1:1 to 1:4 ratio, respectively.

2.4. RF Ablation Experiments

Gold nanoparticles with different densities were prepared before thermal ablation experiments. The calibration of MgO-doped optical fibers was conducted before the experiments. Gold nanoparticles at different densities such as 1 mg/mL and 4 mg/mL were introduced ex vivo in the parenchymal tissue surface and around the applicator. The hybrid generator and reflectometer turned on at the same time: during the first phase, the tissue was ablated; then, the generator was automatically turned off by reaching the safe mode impedance value at $800~\Omega$, while the measurement was continued for the next 50~s. The applicator and fibers were positioned on opposite sides to each other on the *y*-axis; the distance between fibers was 4 mm on the *x*-axis and the tip of the applicator was placed

between the third and fourth fibers (Figure 6). Thermal ablation experiments for each condition were repeated four times: pristine, pure agarose, gold nanoparticles with the density starting at 1 mg/mL, and gold nanoparticles with a density of 4 mg/mL to avoid contamination of fibers.



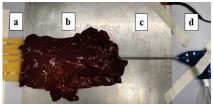


Figure 6. The scheme and photographic view fibers (a), bovine liver (b), negatively charged plate (c), and the applicator (d) in the vicinity.

In the experiments, the applicator and fibers have been positioned manually by the operator in order to maintain the relative distance between the electrode tip and each fiber of the grid; agarose and agarose/AuNP were introduced after the fiber insertion, through a syringe. In order to preserve the fidelity of fiber insertion through the tissue, the fibers were inserted through the phantom, from one side; Figure 6a shows the cross-section of the tissue, while in Figure 6b, the photograph shows that in the actual experiments, the liver tissue was positioned around the fibers. Post-ablation photographs are displayed upon cutting the tissue in the cross-section.

MgO-doped optical fibers were spliced to single-mode fibers using a standard telecom splicer, resulting in a sensing system that detects the temperature change during radiofrequency ablation with 1-s speed. Figure 7 shows the resulting amplitude trace recorded on the OBR when all six fibers are connected, forming a scattering-level multiplexing network that allows simultaneous temperature recording on each location. The scattering gain of each nanoparticle-doped fiber is about 40 dB. The received data were first processed using a threshold algorithm which can identify sudden dips in signal amplitude and, thus, exact locations of all six fibers' ends were discovered. Then, the peak temperature reading, which is in correspondence of the vertical coordinate of the applicator tip, was determined for each fiber and 20 data points around it were used for interpolation (using spline functions). Finally, interpolated domains were converted to thermal maps to accurately monitor the temperature change area.

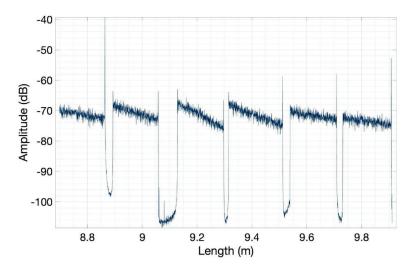


Figure 7. Backscattering trace, recording the Rayleigh backscattering intensity at every point along the fiber sensing network. The six regions with high intensity correspond to each nanoparticle-doped fiber span, while the low-intensity regions correspond to single-mode fibers used to distribute the signal to each sensing fiber.

3. Experimental Results

3.1. Thermal Maps

We report in Figure 8 an example of a thermal map, recorded with the previously described setup for RFA ablation carried out with AuNP with a density of 4 mg/mL. Data are recorded on a xy grid with size 40 mm \times 20 mm; the pixel of the grid is 2.5 mm (OBR spatial resolution after processing) \times 5.0 mm (distance between each MgO-doped fiber), for a total of 102 sensing points (one sensor per each 7.8 mm²). Data reported in the figure show the isothermal curves, spacing the temperature data by 10 °C.

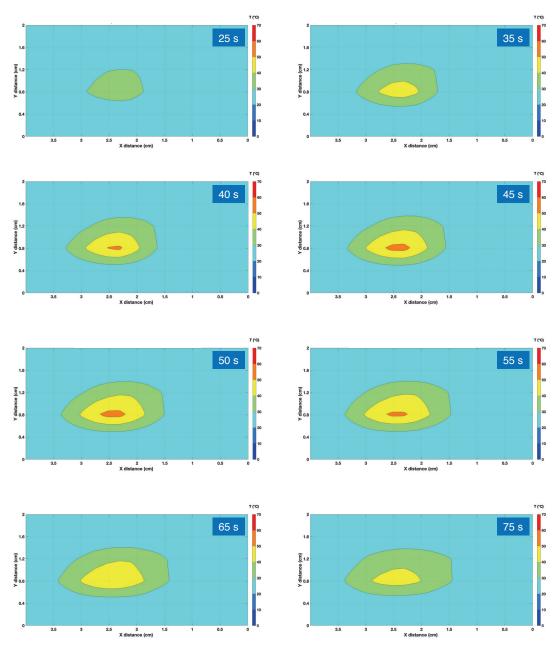


Figure 8. Thermal map recorded for an RFA experiment, using AuNP with a density of 4 mg/mL. The chart reports the data recorded on the xy plane (x = direction parallel to the RFA applicator and to the sensing fibers). The colorimetric map reported the isothermal curves, with 10 °C separation between each layer. Eight different maps are reported, during the heating phase (with times 25, 35, 40, 45, and 50 s from the RF generator power on), and during the cooling phase when the RF power is discontinued (55, 65, and 75 s elapsed time). The horizontal lines on each chart report the position of the six fibers.

During the first part of the ablation, we observe a progressive heating as the temperature rises, peaking in correspondence to the active electrode. After 40 s, the inner temperature overcomes the $60\,^{\circ}\text{C}$ threshold, reaching the maximum extension after 50 s. At this point, the RF generator gets discontinued, causing the temperature to drop.

In Figure 9, we compare the thermal map observed at the peak temperature, in correspondence to the RF power discontinued from the generator. Thermal maps are showed in comparison with the pictures of the ablated tissue, for each different experiment: pristine ablation, agarose-mediated ablation, and AuNP-mediated ablation with densities of 1 mg/mL and 4 mg/mL. We observe that the experiments return a similar result in terms of the shape of the ablated region; AuNP-mediated RFA, however, results in a slightly elongated pattern of the isothermal curves.

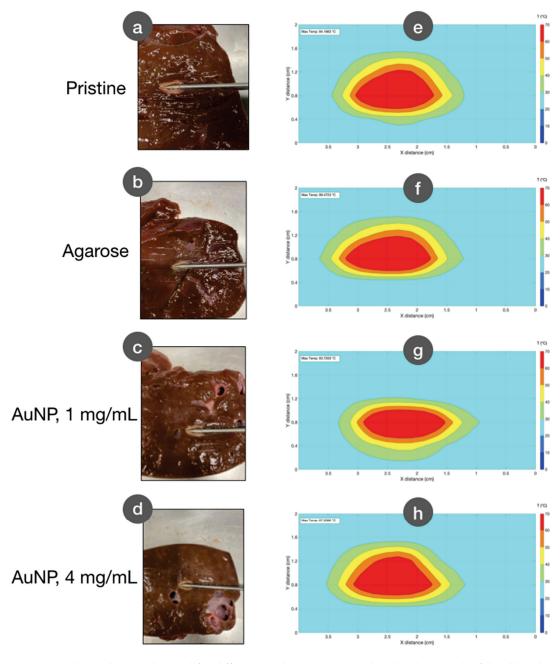


Figure 9. Thermal maps observed for different conditions, compared to the photograph of the ablated tissue. (\mathbf{a} – \mathbf{d}) Photographs of the tissue after RFA procedure; (\mathbf{e} – \mathbf{h}) thermal maps observed at the peak temperature condition. Experiments have been performed in pristine condition (\mathbf{a} , \mathbf{e}), with agarose gel inserted in the tissue (\mathbf{b} , \mathbf{f}), and with AuNP with densities of 1 mg/mL (\mathbf{c} , \mathbf{g}) and 4 mg/mL (\mathbf{d} , \mathbf{h}).

The experiments carried out during RFA are subjected to a high variability of results [27]. While the distribution of the nanoparticles in the tissue appears to play a role [40], the major source of variability is due to the different properties of the tissues: this is both attributable to the differences in electrical and thermal properties of the parenchyma, which also tend to increase in tumoral tissues [31], and also to the presence of blood vessels that prevent an efficient heat distribution [52]. In order to account for the repeatability of the ablation process, we performed four experiments for each ablation type, all using a similar phantom. The results are shown in Figure 10, where the thermal maps at the temperature peak are reported. We observe a difference in the shape and extension of the thermal zones: in particular, the AuNP-mediated isothermal curves appear to have a rounder contour, more similar to an elliptical shape; conversely, pristine ablation returns isothermal shapes characterized by more irregularities.

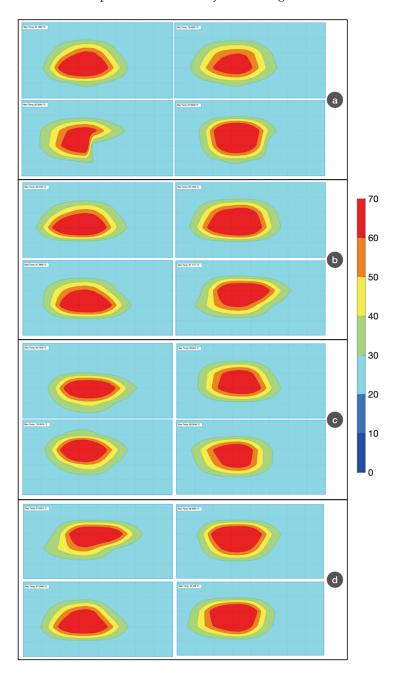


Figure 10. Evaluation of all thermal maps for four experiments, under each condition: (a) pristine; (b) agarose; (c) AuNP, 1 mg/mL; (d) AuNP, 4 mg/mL. The chart reports isothermal curves, according to the color bar on the right. Horizontal axis: x direction (4 cm range); vertical axis: y direction (2 cm range).

3.2. Maximum Temperature

The maximum temperature is an important indicator of the heating process occurring during RFA. According to Rhim et al. [53], temperature values exceeding 120 °C should be avoided, while the 80-100 °C peak temperature ensures the most efficient ablation process. Figure 11 shows the peak temperature recorded for all experiments, reporting both the average and the standard deviation. Pristine ablation shows an average temperature of 89.0 °C (the lowest value); whole agarose-mediated RFA shows 95.3 °C average (the highest value); and AuNP-mediated RFA returns similar temperature values (92.2 and 92.3 °C, respectively). AuNP-mediated RFA with 4 mg/mL shows the best repeatability of peak temperature (standard deviation = 4.0 °C), while AuNP-mediated ablation shows the worst repeatability (standard deviation = 13.7 °C).

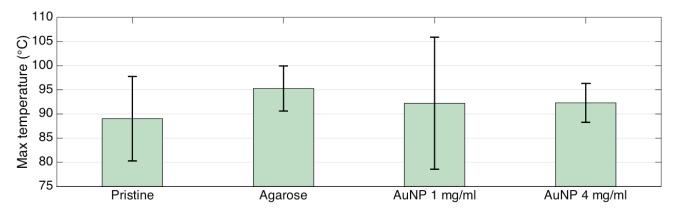


Figure 11. Peak temperature recorded during the RFA experiments under each condition; bars = average of four experiments; error bars = \pm standard deviation.

The temporal evolution of the peak temperature shows insights on the heating process and the duration of the ablation procedure [54]. In Figure 12, the temporal trend of the peak temperature is reported for each experimental condition, for the first 50 s of the ablation process. We observe that the duration of the ablation process is similar for all conditions, about 37 s with a larger variability for the AuNP-mediated condition, 1 mg/mL and pristine ablation, and a smaller extension of the standard deviation region for the AuNP-mediated ablation, 4 mg/mL and agarose-mediated. Looking at the average trend, the agarose-mediated ablation reaches the 60 °C threshold in the fastest time (15 s), about 2 s faster than the pristine ablation.

3.3. Evaluation of Thermal Damage Regions

Thermal dosimetry estimates the amount of thermal damage induced by the ablation process [10]. The mortality rate of cancer cells is nearly instantaneous for temperature values higher than 60 °C, while no damage is induced below 42-44 °C depending on the type of tumor. For intermediate values, the mortality rate is a function of the temperature and the exposure time, and a common estimate is 52 °C for one minute of exposure. However, since RFA is a faster process with a duration of ~35 s, we can highlight three regions: (1) maximum thermal damage (temperature >60 °C), where the mortality rate is almost ideal; (2) cytotoxic region (temperature between 42 °C and 60 °C), where temperature induces a partial damage of the tissue; and (3) safe region (temperature below 42 °C), unaffected by RFA. The possibility of drawing isothermal curves at the moment of maximum heating determines the width of each region.

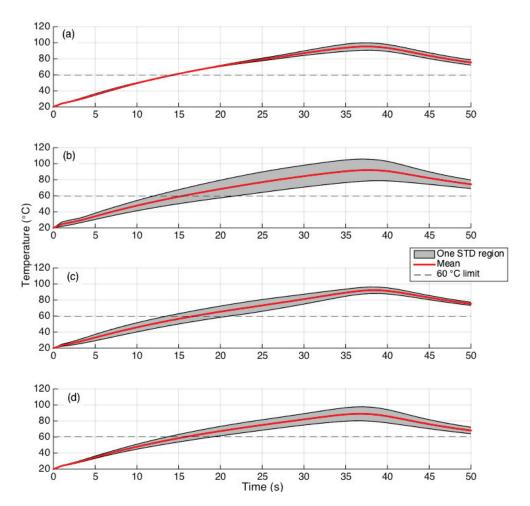


Figure 12. Temporal evolution of the maximum temperature over time, for each RFA ablation. The chart reports the results of four experiments, displaying the mean value (red curve) and the range containing \pm one standard deviation (grey interval). RFA conditions: (a) agarose; (b) AuNP, 1 mg/mL; (c) AuNP, 4 mg/mL; (d) pristine.

Figure 13 shows the quantification of the thermal damage condition. Pristine ablation shows an ablated surface of $55.0~\rm mm^2$ (mean value), and has the widest extension (33.9– $66.4~\rm mm^2$). Agarose-mediated RFA shows the highest extension of the ablated region ($67.3~\rm mm^2$ mean value, $74.7~\rm mm^2$ maximum). On the other side, the AuNP-mediated ablation with $4~\rm mg/mL$ density shows the best trade-off between amount of ablated tissue ($61.5~\rm mm^2$ mean value) and repeatability (53.0– $65.1~\rm mm^2$ minimum—maximum values). The percentual indicators displayed on the right chart show the quantification of the increase or decrease of ablated tissue with respect to pristine ablation (mean value): it shows that pristine ablation has the worst amount of ablated tissue and the widest repeatability range, while agarose-mediated RFA can ablate up to 19.8% more of tissue.

In Figure 14, the cytotoxicity regions are shown, reporting the areas of the tissue exposed to temperatures >42 $^{\circ}$ C. We observe a similar trend with respect to the thermal damage results, with agarose performing the widest region of damage (mean value 132.3 mm²) but with the widest range, and AuNP-mediated nanoparticles with the higher density to represent the best trade-off between amount of thermal damage (mean value 127.9 mm², +8.5% with respect to pristine conditions) and repeatability (116.7–141.1 mm² minimum—maximum values).

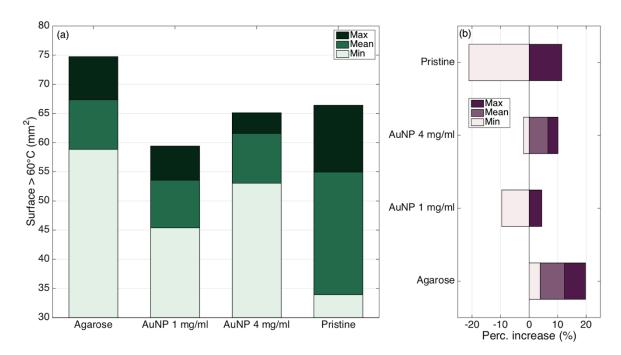


Figure 13. Quantification of the thermal damage region (temperature >60 °C) for each experimental condition. (a) Evaluation of the surface exposed to temperature exceeding 60 °C at the maximum ablation temperature; bar charts show the maximum (dark), minimum (bright), and mean (intermediate) values of the areas, over four experiments for each RFA condition. (b) Percentual increase or decrease for the ablated surface, with respect to the reference condition (pristine RFA, mean surface).

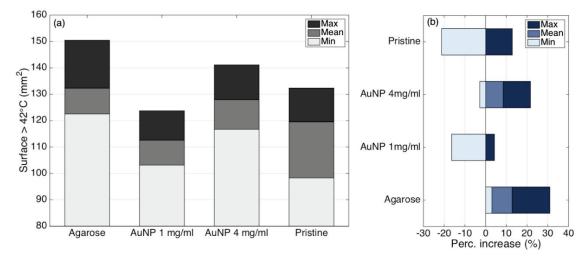


Figure 14. Quantification of the cytotoxic region (temperatures >42 °C) for each experimental condition. (a) Evaluation of the surface exposed to temperatures exceeding 42 °C at the maximum ablation temperature; bar charts show the maximum (dark), minimum (bright), and mean (intermediate) values of the areas, over four experiments for each RFA condition. (b) Percentual increase or decrease for the cytotoxic surface, with respect to the reference condition (pristine RFA, mean surface).

4. Discussion

The display of results shows the effectiveness of the ablation method, and the importance of the in-situ sensing device for the real-time detection of temperature patterns. A first point of discussion involves the efficacy of temperature sensing in the context of thermal ablation. As highlighted in Figure 10, the lack of repeatability of RFA requires a precise thermal control in situ: distributed sensing in this case is a very effective technology, as it ensures not only rapid sensing with accurate detection (about 1 pm uncertainty,

corresponding to ~0.1 °C), but the entire sensing network is designed for the possible in vivo use, since the fibers have miniature form factor, biocompatibility, and the sensing mechanism is robust with respect to the catheterization options [28] and possible strains occurring during the insertion [55]. The proposed sensing network has the potential of being much more cost-effective than FBG arrays, in terms of disposable use. It is in fact essential to ensure that the cost of the whole sensing device that complements the RFA device is a fraction of the applicator itself, in order to ensure an affordable treatment; while FBG arrays have a cost of a few hundreds of dollars per array, the MgO-doped fiber can be fabricated with the same technology in terms of doping, preforming, and drawing of SMF fibers that cost a few dollar cents per meter (about \$0.08/m in current markets). Additionally, MgO-doped fibers can be spooled and spliced to SMF fibers using standard splicers, without the need to develop new splicing methods. In addition, FBG sensors encode the sensing in a specific location, and therefore they are both limited by the grating length and misalignment of each grating position [12]. Distributed sensing, on the other hand, can improve the spatial resolution down to 10 μm (theoretical value of the OBR [56]), and since the entire fiber acts as a sensor, it is possible to align thermal maps using signal processing methods or other artefacts [57]. On the other hand, distributed sensing provides a much more detailed detection with respect to thermocouples, even in miniaturized formats [58], or fluorescence-based detectors [59], as these methods are intended for single-point sensing while the proposed method can resolve several tens of sensors per each cm².

Regarding the thermal ablation process, the analysis of cytotoxicity and thermal damage levels shows that the RFA process improves when using agarose and Au nanoparticles with the appropriate density. While agarose-mediated RFA achieves the best results in terms of heat delivery, guaranteeing the largest thermal ablation results, the use of AuNP shows the best compromise between the width of the ablated area and the need for a repeatable process. As validated by experiments, by using 4 mg/mL density, the best results can be achieved. The possibility of extending the ablated region is an important asset in advancing RFA to treat larger types of cancer, particularly for hepatic tumors [60].

5. Conclusions

In conclusion, we reported the use of a distributed fiber-optic sensing network optimized for the detection of thermal patterns in RFA procedure, mediated by the use of agarose and gold nanoparticles for the improvement of the heat delivery. The fiber-optic sensing network is formed by a set of six MgO-NP fibers, deployed over a grid of 40×20 mm, forming 102 sensing points spanning the spatial resolution of the interrogator on the horizontal axis and the spacing between each fiber on the vertical axis.

Several experiments of RFA were performed, in pristine condition as well as mediated by agarose and AuNP (1 and 4 mg/mL density). Thermal maps and isothermal representation allow precise identification of the ablated region, determining the area exposed to the ideal thermal damage (60 $^{\circ}$ C and above) and cytotoxicity levels for partial mortality (between 42 and 60 $^{\circ}$ C). Experimental results show that agarose-mediated thermal ablation yields the widest treated area, but with low repeatability; on the other hand, AuNP-mediated ablation using 4 mg/mL density provides the best trade-off between efficacy of ablation and repeatability.

Future work will further consolidate the RFA applicator, the on-board sensors for real-time detection, and the nanoparticle delivery into a single device, scaling up the possibility of using this arrangement in clinical settlements.

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Review

Cost-Effective Fiber Optic Solutions for Biosensing

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Abstract: In the last years, optical fiber sensors have proven to be a reliable and versatile biosensing tool. Optical fiber biosensors (OFBs) are analytical devices that use optical fibers as transducers, with the advantages of being easily coated and biofunctionalized, allowing the monitorization of all functionalization and detection in real-time, as well as being small in size and geometrically flexible, thus allowing device miniaturization and portability for point-of-care (POC) testing. Knowing the potential of such biosensing tools, this paper reviews the reported OFBs which are, at the moment, the most cost-effective. Different fiber configurations are highlighted, namely, end-face reflected, unclad, D- and U-shaped, tips, ball resonators, tapered, light-diffusing, and specialty fibers. Packaging techniques to enhance OFBs' application in the medical field, namely for implementing in subcutaneous, percutaneous, and endoscopic operations as well as in wearable structures, are presented and discussed. Interrogation approaches of OFBs using smartphones' hardware are a great way to obtain cost-effective sensing approaches. In this review paper, different architectures of such interrogation methods and their respective applications are presented. Finally, the application of OFBs in monitoring three crucial fields of human life and wellbeing are reported: detection of cancer biomarkers, detection of cardiovascular biomarkers, and environmental monitoring.

Keywords: POC monitoring; smartphone optical biosensors; optical interrogation methods; cancer biomarkers; cardiovascular biomarkers; environmental monitoring

1. Introduction

Biosensors are devices that quantify (or semi-quantify) a biological or chemical analyte by generating a measurable signal proportional to its concentration [1]. They are usually applied in different fields, from biomedical to environmental, allowing monitoring of specific disease biomarkers in body fluids (blood, urine, saliva, and sweat) [2] and detection of micro-organisms [3] and pollutants [4] in the environment, among other things. In the case of optical biosensors, the signal generated can be coded in wavelength, phase, and signal intensity, among other optical features [5]. Optical absorbance, fluorescence, bioluminescence, interferometry, ellipsometry, reflectometric interference spectroscopy, and surface-enhanced Raman scattering are also examples of highly applied optical techniques to biosensing [6].

In relation optical fiber biosensors, usually, the biosensing process relies on the interaction of the evanescent wave with the fiber's surroundings. Therefore, in order to turn the fiber sensitive to the environmental refractive index (RI) changes, fibers can have gratings inscribed or can be altered in different geometries. Moreover, noble metal coatings with thicknesses around 30–70 nm, enabling the surface plasmon resonance (SPR) phenomena,

or nanoparticles, and thus generating a localized SPR (LSPR), can be applied. The use of optical fibers brings important advantages, namely: small dimensions—allowing them to be used in biomedical applications in invasive and non-invasively approaches; dielectric nature—making them electrically safe; and low signal attenuation [7]. The main disadvantages to optical fiber biosensors are the fragility and high-cost interrogation approaches [8]. However, as research on optical fiber sensors (OFS) has been evolving, a large number of studies reporting robust and cost-effective optical fiber biosensing (OFB) approaches has been published in the past few years. The sensing principles that will be described are mainly based on interrogation in the visible spectrum due to the well-known cost-effectiveness of such an approach. However, it should not be forgotten that, in the last few years, efforts that have been made in improving interrogation costs of infrared and gratings spectral analysis. Currently, this is a hot topic in this research field with excellent outputs [9–11].

This review focuses on the use of OFB approaches that are cost-effective. As can be verified in Figure 1, different fiber geometries and structures will be addressed, namely, U-bent, D-shaped, unclad fibers, tapered fibers, ball resonators, end-face reflection, and light-diffusing fibers (LDFs). Different OFB packaging methods for robustness improvement will be also presented, as well as examples of sensors coupled with smartphones. Cost-effective biofunctionalization methods, mainly based on aptamers and molecularly imprinted polymers (MIPs), will also be discussed. Different applications on cancer and cardiovascular biomarkers' detection and environmental monitoring will be presented, and finally, a critical analysis of the state-of-the-art and prospects will be made.

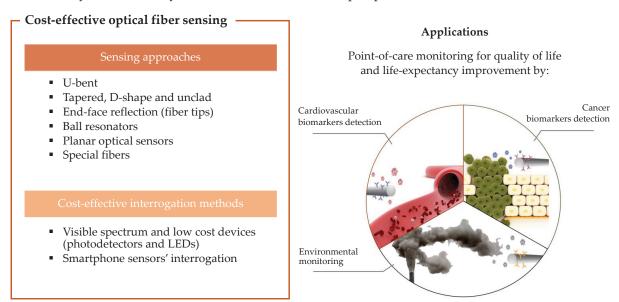


Figure 1. Schematic illustration of the main contents of this review.

2. Cost-Effective Optical Fiber Configurations for Biosensing

Starting from the optical fibers that are usually applied in cost-effective biosensing approaches, different kinds of fibers can be used, from regular silica single-mode fibers (SMFs) [12] to multimode fibers (MMFs) such as plastic-clad silica (PCS) [13] and plastic optical fibers (POFs) [14]. The different kinds of optical fibers, conventional silica-based, microstructured, or specialty fibers in terms of materials and designs offer several useful advantages in developing sensors/biosensors. For instance, when telecommunication wavelength range and setups are used, the silica-based fibers are more appropriate, such as to exploit the monomodal characteristic for several sensing approaches. In contrast, POFs are especially advantageous due to their excellent flexibility, lower attenuation in the visible range, easy handling, multimodal characteristic, and the fact that plastic can withstand smaller bend radii than glass [15].

As regards the optical signal trajectory, as shown in Figure 2, the different sensing configurations, which will be described as follows, can work on transmission (a) or in reflection (b). Additionally, OFBs can be used as intrinsic and extrinsic OFS. In particular, when the optical fiber interacts with the analyzed medium directly, it is defined as an intrinsic OFS, whereas, when it is used as an optical waveguide allowing the launch of the light to the sensing region and its collection, it is defined as an extrinsic OFS.

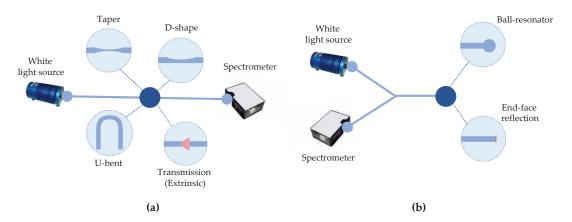


Figure 2. Schematics of OFB working in (a) transmission and (b) reflection.

The interrogation methods in the scope of this review are mainly based on the analysis of the visible spectrum, with data codification in wavelength, especially when using SPR and LSPR, and in intensity variations. Although both SMF and MMF can be suitable for wavelength interrogation, as concerns intensity-based approaches, OFBs find great advantage in the use of fibers, having a large core, which implies a high numerical aperture and strongly multimodal operation. The possibility of applying MMFs enables the use of inexpensive light-emitting diodes (LEDs) [16] and simple photodetectors such as charge-coupled devices (CCDs) or photodetectors (PDs) [17], as light can be easily coupled into the fibers using basic optics and cheap interconnections. In contrast, SMFs require hardware with higher complexity and costs.

While the use of a large numerical aperture and strongly multicore fibers is typically a disadvantage in the design of physical sensors, several biosensing structures exploit the favorable properties of MMFs as part of the detection systems. From a light propagation standpoint, in a multimode fiber, modes with a lower index are strongly guided, while higher-order modes are weakly guided and become more sensitive to geometrical modifications of the fiber, such as thinning the fiber core [17], removing the cladding in a D-shaped form [18] or completely, through an etching [19,20] or mechanical [21] process, or bending the fiber in curved shapes [22]. The two most notable approaches for low-cost biosensors based on geometrically modified MMFs are based on U-bent fibers [22,23] and tapered fibers [24].

2.1. U-Bent Biosensors

The U-bent fiber is a transmission device operating in the visible spectral range that can be made of a SMF [25] or MMF [15], in which the fiber is curved with a specific bend radius, as highlighted in Figure 3a. During the bending process to obtain the U-shape, the fiber experiences mechanical stress, which causes a variation in its optical properties due to the photoelastic effect. The induced stress will yield a non-uniform RI of the core, with a higher RI at the inner bend and lower at the outer bend of the core. As the bend radius decreases, the RI profile changes more significantly due to the increased stress. Lower order modes are well confined within the fiber and propagate through with minimal losses; higher-order modes, due to the increased angle of incidence at the interface between core and cladding, do not satisfy the total internal condition anymore and are strongly attenuated. Consequently, the intensity of the field increases at the outer region of the core

and, at the core-cladding interface, the Fresnel transmission coefficient approximates 1, leading to greater radiation loss. As regards the inner interface of the bent region, more light satisfies the total internal reflection principle, with less light lost in that region. In U-shape fibers, there is a critical radius of curvature, below which the fiber no longer transmits light. The critical radius of curvature decreases for a large difference of RI or large NA values. Larger NA fibers can bend with a smaller bend radius without significant loss of light given the same wavelength of light [26].

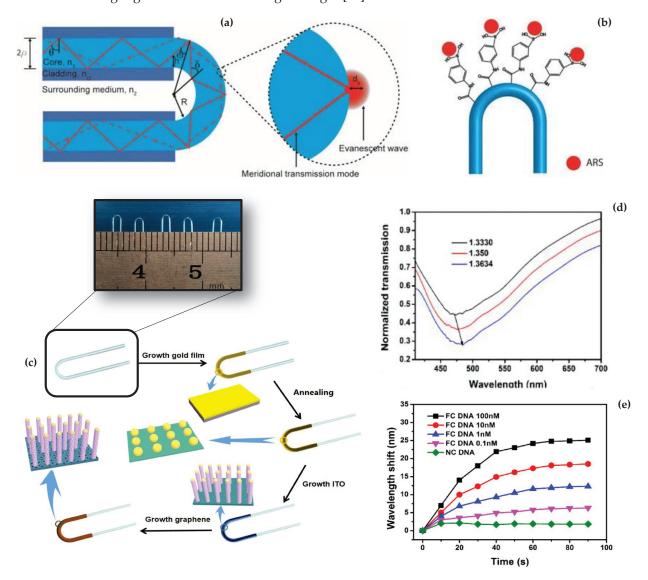


Figure 3. Schematic and implementation of U-shape fiber biosensors: schematic of a U-bent fiber section with light propagation within the curved portion of the fiber (a) with inset showing the surface-penetrating meridional rays providing RI sensitivity, and functionalization of the biosensor for glucose detection (b) (ARS—Alizarin Red S) (images reproduced from [27]); (c) U-bent LSPR sensor proposed by Yang et al. with the respective production steps, namely, U-shape production and photographs, and formation of ITO coating for plasmonic-based sensing, with results of the change in transmission spectrum for three different values of RI (d) and the response of the sensor, measuring the wavelength shift over time, for different DNA concentrations (e) (images adapted from [28]).

In order to transform this structure into an efficient RI sensor capable of sensing the surrounding environment, the design [22] makes the use of fiber with an exposed core (hence, with no cladding or depleted cladding) in U-shape. This sensing structure is shown in Figure 3a and was optimized for biosensing according to Azkune et al. [27]. The bent

portion of the fiber has no cladding, depleted through wet-etching, dry etching, or by simply fabricating a cladding-less fiber. POF fibers in this regard have been substantially used in U-bent fiber sensors [27,29–31] as they are more stretchable, having a lower Young's modulus, and millions of modes are propagated, hence, guaranteeing a measurable change of transmission when the RI in the surrounding medium changes.

In addition, the U-bent fiber structure shows the best versatility when packaging the device, as the fiber can be tightly folded within a narrow space. Chen et al. [32], among others, reported a U-bent fiber sensor folded in a 1.1 mm tube structure, which is compatible with most medical needles, to fabricate a device with a high level of miniaturization.

The possible functionalization of a U-bent fiber section, as reported in Figure 3b [27], relies on immobilization of bioreceptors on the bent portion of the fiber, in correspondence with the region with the maximum sensitivity. The authors propose a glucose detection system based on the release of ARS (Alizarin Red S) occurring when glucose is bound through competitive binding to phenylboronic acid.

U-bent fibers are often used in conjunction with SPR effects to improve detection. While a U-bent fiber provides a significant transmission change, the wavelength shift observed due to the plasmonic resonance can be detected in a more robust way. This principle was implemented by Yang et al. [28], who reported a DNA biosensor by a thin-film U-bent plasmonic graphene/ITO (indium tin oxide) biosensor. The main achievements reported by Yang et al. are also shown in Figure 3. The Figure 3c shows the graphene/ITO layer deposition around the core-exposed section of a POF, folded around a 1.3 mm bending diameter in the inner region. The transmission spectrum through the bent section of the fiber shows both an intensity change and a wavelength shift due to the combined effect of the fiber losses of meridional rays and localized SPR at the outer interface (Figure 3d. By measuring the wavelength shift, Yang et al. estimated the sensor response to different full complementary DNA and non-complementary DNA at 0.1–100 nM concentrations, as displayed in Figure 3e.

Several biosensors have been reported as U-bent fiber probes based on MMFs and inexpensive spectrometers and hardware based on these principles. George et al. [29] reported an immunobiosensor for the detection of chikungunya (an infectious disease) through the non-structural protein 3 (CHIKV-nsP3) biomarker, achieving a limit-of-detection (LOD) of 0.52 ng/mL. Bandaru et al. [33] reported the production of a plasmonic U-bent fiber biosensor capable of detecting human immunoglobulin G with ultra-low LOD (\sim 7 aM). Manoharan et al. [34] described a similar architecture for detecting bacterial endotoxins such as lipopolysaccharides with a LOD of 0.4 ng/mL.

2.2. Tapered Fiber Biosensors

Tapered fibers have been used to implement fiber-optic interferometers [35] or multiparametric sensors [36] in SMFs. In MMFs, tapers are used to induce losses in the fiber without the need for bending, which allows for easier packaging of the sensor in medical devices [37]. The fabrication of fiber tapers is implemented by modern CO₂ or fiber laser splicers [38] as a rapid process; it is, therefore, compatible with large volume manufacturing. Figure 4 shows a tapered silica SMF with LSPR for biosensing, developed by C. Huang et al. [39]. As can be seen, in a tapered fiber, the diameter of the fiber is progressively thinned until it reaches a minimum diameter in the so-called waist region, that corresponds to the minimum thickness (region A), with an intermediate region with variable diameter (region B). Unlike the etching process, tapering maintains the proportion between core and cladding diameter. When exposed to a thickness change, the lower radius of the core filters out the higher-order modes as the V number of the fiber decreases, reaching its minimum in the waist region [37]. For the case of an LSPR-based taper, it can be seen in Figure 4d,e that its response changes in wavelength and relative intensity/transmittance in the visible region.

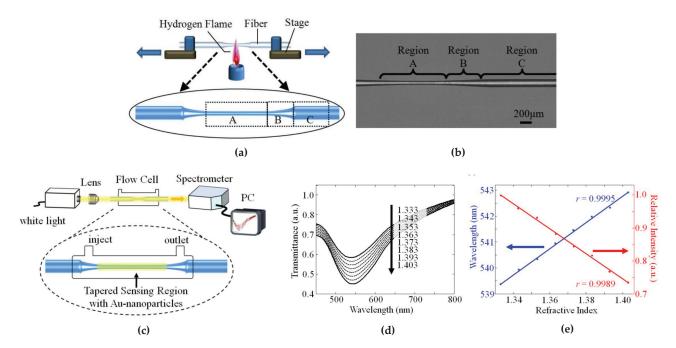


Figure 4. Schematic and implementation of tapered OFB: (a) taper production highlighting the respective fiber taper regions and (b) optical microscopy image of the tapered fiber showing the three regions; (c) taper coated with Au NPs and respective interrogation scheme, and (d) results showing the variation in wavelength and transmittance with RI variation with respective linear fittings (e) (images reproduced from [39]).

POF are also good candidates for low-cost tapered fiber structures, as the fiber has very thin cladding and, therefore, allows the core to be exposed to the outer environment. Dash et al. [38] reported a fiber taper in a POF having an initial diameter of 250 μ m tapered to 175 μ m and coated with a layer of gold nanoparticles (Au NPs) over a 4 mm length. By measuring the output power after the tapered section, a drop is observed when the outer RI increases; a power change of about 5% has been observed when measuring up to 0.2 mM of bovine serum albumin (BSA) concentration. A similar system was reported by Rahman et al. [40] for water salinity detection.

2.3. D-Shaped and Unclad Biosensors

As U-bent and tapered OF sensing structures, D-shaping or uncladding an optical fiber aims to increase the interaction of the evanescent field with the surrounding medium. The D-shaped OFS, or side-polished, can be made in SMF and MMF, usually by polishing. As can be seen in Figure 5, the polishing is made until the core surface is exposed [41], but, in other cases, the core can be completely removed in that region [42]. The distance from the polished surface to the fiber surface is usually referred to as the residual thickness and is of utmost importance in the strength of the interaction of the evanescent wave fields with the external medium. For this specific sensor, developed by Zheng et al. [42] and represented in Figure 5a, the optimal RI sensitivity was achieved by a residual thickness of 34.09 μ m, attaining a value of 10,243 nm/RIU (RI units). In this kind of structure, the surface roughness also seems to have a crucial impact on the sensor's response [43]. A great feature that D-shaped OFS have is the flat surface that can be achieved, facilitating deposition of different materials [44] and, also, biosensing testing with only a sample drop [45].

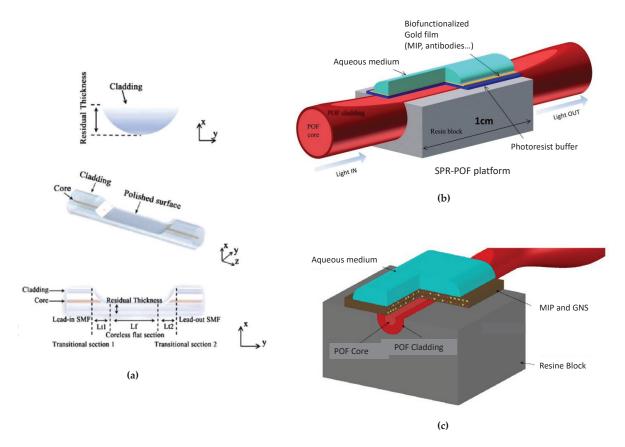


Figure 5. Schematics of D-shaped OFBs: (a) core-less D-shape in different perspectives (images adapted from [42]); (b) SPR-POF D-shaped biosensing platform (image adapted from [46]. Copyright ©2021, with permission from Elsevier); (c) LSPR D-shaped tapered-POF biosensor based on gold nanostars (GNS) and MIPs (image reprinted from [47]. Copyright ©2015, with permission from Elsevier).

Cennamo et al. [14,18] have presented SPR sensors produced by side removing the cladding and part of the core of the POF, spinning a thin layer on the exposed core with a RI major of the core, and finally sputtering a gold nano-film. The buffer layer under the gold nano-film improves the performance of the SPR sensor and the gold film's adhesion on the platform. Figure 5b outlines these SPR–POF platforms, in which the gold layer can be functionalized with different bioreceptors, namely MIPs [48], aptamers [18], and/or antibodies [49]. Moreover, D-shaped tapered-POFs have also been used to accomplish different kinds of plasmonic bio/chemical sensor configurations, as depicted in Figure 5c, for the detection of trinitrotoluene (TNT) [47]. On the other hand, unclad OFBs have the advantage of a higher interaction of the evanescent wave with the surrounding environment. The removal of the cladding can be achieved by polishing using sandpapers in rotative movements [21] or, in the case of silica fibers, by chemical etching [19].

2.4. End-Face Reflection/Optical Fiber Tips

Optical fiber tips are sensors that work on reflection, usually by having the fiber end coated with a reflective/plasmonic material such as silver or gold. In its simpler form, fiber tips are flat at the end and are usually created in MMFs, for instance, PCSs [50], which are interrogated using bifurcated optical fibers connected to the sensor by means of a bare fiber terminator, such as the BFT1 from Thorlabs. To produce flat tips, the fibers are usually cleaved by precision cutters, since a flat surface is crucial for assuring the collected signal spectra reflectivity [51]. Nevertheless, a wide range of tip forms and functionalities can be achieved employing different fabrication technologies, such as mechanical processing,

chemical etching, laser processing, self-assembly, and chemical or physical vapor deposition processing, as well as material transfer methods [52].

Fiber tips-based biosensors are expected to be an important tool in POC monitoring due to their high RI sensitivity, compactness, and ease of operation. For example, for flat fiber tips coated with gold by sputtering, it is possible to achieve RI sensitivities between 700 and 2000 nm/RIU depending on the gold film thickness [50]. An alternative approach was reported by Lee et al. [53], who developed a tip aptasensor based on LSPR created with gold nanorods (GNR). The sensor structure and response in the detection of ochratoxin A (OTA) can be seen in Figure 6.

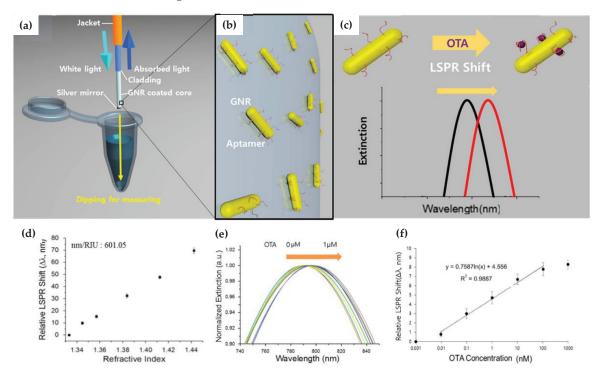


Figure 6. Schematic of a tip OFB: (a) components and measurement procedure; (b) magnification of the sensor surface biofunctionalized with aptamer-modified GNR; (c) schematization of LSPR wavelength shift when the analyte, in this case, OTA, is recognized by the specific aptamer; (d) GNR-coated tip OFS response to RI; (e) LSPR wavelength shift with different concentrations of OTA, and (f) experimental response of the biosensor to different concentrations of OTA with respective linear fitting (images adapted with permission from ref. [53]. Copyright ©2018, with permission from Elsevier).

2.5. Ball Resonators

A ball resonator is a device where a circularly shaped microsphere is fabricated on the tip of an optical fiber of approximately a hundred micrometers in diameter. The RI inside the ball resonator is higher than the environment around the spherical structure. In this case, the light that goes through optical fiber and reaches the edges of the microsphere is reflected inside the ball by total internal reflection and, therefore, cannot escape to the surrounding environment [54]. Nevertheless, due to the fact that the surface of the ball resonator has a curved interface, its total internal reflection is never complete; therefore, bending losses could occur, which can lead to the light leakage out of the ball resonator. This leads to an enhanced interaction of samples outside the ball resonator with light thanks to the confinement and rebounding of light inside the circular cavity [55,56]. These evanescent waves are able to interact with both the surface of the microsphere and the environment that surrounds it. Each waveguide of optical field circulation allows it to contact the surrounding environment, which subsequently increases the sensitivity as the amount of circulating rounds rises. The photon lifetime in the microspheric cavity is able to measure

it, which is called the quality (Q) factor. The higher the lifetime of the photon, the lower the limit of detection. As previously stated, evanescent waves enable some of the light inside the ball resonator to escape into the outer environment. This allows a microsphere to acquire an increased sensitivity, which makes it useful for a host of applications. Interaction of the surface of the ball resonator with an analyte outside it results in changes in the effective RI. The experience of the analytes of various refractive indexes results in spectral changes such as wavelength and frequency shift. These favorable properties make ball resonators promising as biosensors for the detection of various analytes.

The advantage of microspheres is that they can be easily fabricated in the laboratory from various materials, organic and inorganic; they are simple in handling and have a high-quality factor [56]. The ball resonator is a device that requires a quick manufacturing process. A CO_2 laser splicer allows it to be manufactured in a controllable way in a single step by fusing two single-mode fibers and then tearing them off to make a spherical tip (Figure 7) [57]. Fabrication of ball resonators is advantageous over the fabrication of other optical devices thanks to its duration, which takes from ~ 60 s up to a couple of minutes depending on the speed of the motors within the machine. Moreover, the fabricated device demonstrates robustness, stability, accuracy in measurements, and higher tensile strength compared to etched fibers, which easily break [57–59].

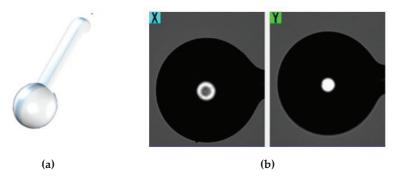


Figure 7. Image of the ball resonator: (a) 3D picture of optical ball resonator; (b) photographs taken by X and Y cameras of CO₂ laser splicer. Image reproduced from [57].

In comparison to previously studied optical fiber-based biosensors such as etched Fiber Bragg grating (eFBG), etched tilted Fiber Bragg grating (eTFBG), and etched MgO nanoparticles (NPs) doped fiber, ball resonators show a rather pronounced cost-effectiveness [58–60]. The use of a CO₂ laser splicer enables the achievement of ball resonators with an accurate diameter and low alignment error in a highly repeatable and fast manner within a few seconds. However, there are some limitations in the use of optical ball resonators. Low reflectivity is one of them, requiring the use of an optical backscatter reflectometer which has excellent sensitivity. Moreover, the weak interference phenomenon occurring in the micro-sphere between the multiple reflective paths results in poor visibility of the spectral fringes.

2.6. Planar Optical Sensor Chips Monitored via POFs

POFs can be used to implement simple and highly sensitive optical fiber sensors by exploiting their advantages, such as the excellent flexibility, the easy manipulation and modification, the large diameter, the fact that plastic can withstand smaller bend radii than glass, and the great numerical aperture (NA). In particular, this last characteristic is very useful in comprising extrinsic optical fiber sensor configurations, for instance, for light injection in plasmonic thin bacterial cellulose slab waveguides [61] and InkJet-printed optical waveguides [62].

Figure 8 shows several plasmonic schemes based on extrinsic POF sensors. A sensing approach based on a nano-plasmonic sensor chip monitored via a custom 3D-printed holder combined with POFs is reported in Figure 8a. The nano-plasmonic chip is based on a gold nanograting (GNG) fabricated on a poly(methyl methacrylate) (PMMA) sub-

strate by electron beam lithography (EBL) [63]. As shown, the light source is connected through a POF optical coupler (50:50) to two POF patches, one illuminating the sensor with nanograting and the other illuminating the reference sensor (a PMMA chip with the same gold film but without the nanograting). At the output of the holder, there are two POF patches used to collect the transmitted light and send it into two similar spectrometers. In this configuration, the PMMA substrate is considered a transparent substrate.

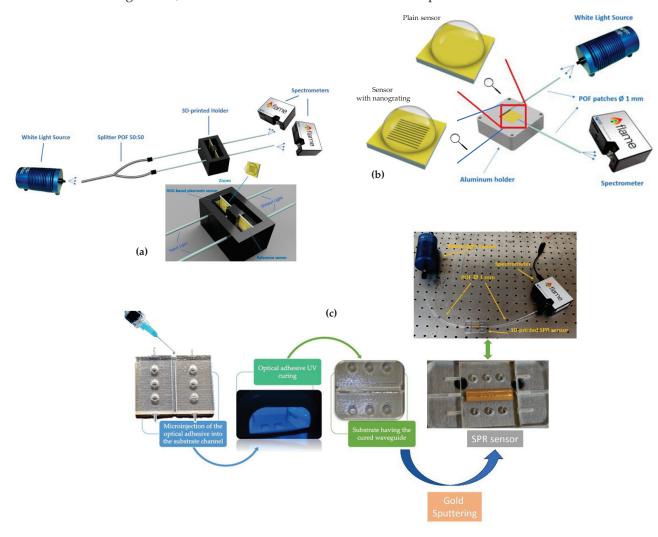


Figure 8. Plasmonic sensor configurations based on extrinsic POF sensors: (a) PMMA-based nanoplasmonic sensor chip monitored by exploiting the transparency of the substrate (image adapted from [63]); (b) PMMA-based plasmonic chips monitored by a custom setup produced to excite the plasmonic phenomena via the multimode slab waveguide (image adapted from [64]); (c) 3D printed SPR sensor chip with the respective setup (image adapted from [62]).

Alternatively, the PMMA substrate can be regarded as a multimode slab waveguide, as shown in Figure 8b, similar to other plasmonic sensor chips [65,66]. In particular, a gold nano-film can be deposited on the PMMA chip, obtaining an SPR sensor chip [67,68] or producing a gold nanograting on the PMMA chip surface via EBL, generating a nano-plasmonic sensor chip [69]. By using the sensor setup illustrated in Figure 8b, when a gold nano-film is present, the sensor's performance is similar to that obtained by an SPR D-shaped POF sensor [64], whereas, when a nano-plasmonic chip is used, the orientation of the nanostripes forming the grating pattern, with respect to the direction of the input light (longitudinal or orthogonal), influences the biosensing performances [69].

Figure 8c shows another low-cost 3D printed SPR sensor chip monitored via POFs. This SPR sensor chip has been designed as a disassembled component composed of four

different parts by using Autodesk[®] Fusion 360, and then the STL files were generated [62]. The sensor's production was performed by a 3D printing method. The material used was a liquid photopolymer ink. Once the SPR sensor parts' construction was completed, the waveguide core of the 3D printed optical device was fabricated. Thus, a UV photopolymer adhesive was microinjected into the sensor channel and cured for 10 min by means of a lamp bulb with UVA emission at 365 nm. Finally, on the cured core, a gold nanofilm was deposited by a sputtering. The thickness of the sputtered gold was about 60 nm.

2.7. Special Fibers

LDF-Based SPR Sensors

Recently, an attractive kind of special optical fibers, named Light-Diffusing Fibers (LDF), has been used to generate highly sensitive SPR sensors. In LDF, the light is not confined to the core; on the contrary, it is diffused on the external medium all along their length. Thus, Cennamo et al. have used this aspect to excite the plasmonic phenomenon on metal nano-films deposited on the LDF [70-74]. In particular, this type of SPR sensor can be produced using silica [70,71] or polymer [72] by coating the fiber with metal nanofilms, such as gold [70,72] or silver [71], by sputtering. This aspect allows an easier fabrication procedure since only a metal deposition step is required to build the SPR sensor. Moreover, SPR sensors based on silica LDF have already demonstrated high performances in bio/chemical sensing applications, as reported in [73], when the SPR-LDF probe is combined with a bioreceptor (e.g., antibody). For instance, Figure 9a shows an outline of a SPR sensor system based on a silica LDF covered by gold nano-film connected to a white light source and a spectrometer [70]. SPR spectra obtained studying the sensor RI sensitivity, using different concentrations of water-glycerin solutions, are represented in Figure 9b. The experimental results reveal a high sensitivity of the SPR wavelength to the outer medium's refractive index (bulk sensitivity), with values ranging from 1500 to about 4000 nm/RIU in the analyzed range. Moreover, it was seen in [71] that these bulk sensitivity values can be improved by using silver nano-films instead of the gold.

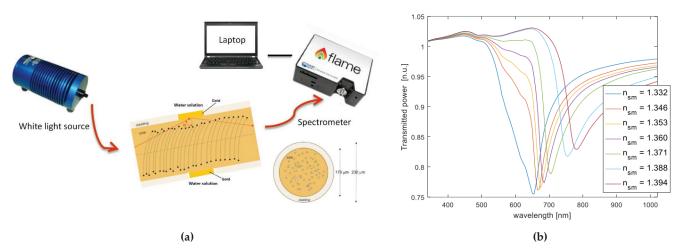


Figure 9. Silica LDF-based SPR sensor system: (a) data acquisition and working scheme; (b) SPR spectra obtained at different refractive indices from 1.332 to 1.394. Image reproduced from [70].

Finally, concerning these types of silica SPR–LDF sensors, Cennamo et al. have demonstrated that the tapering process produces a significant worsening of the bulk sensitivity and a slight decrease in the full width at half maximum (FWHM) of the SPR spectra [74]. Therefore, in this SPR–LDF sensor, a possible alternative consists in using the tapered LDF as a modal filter (after the sensitive region) by determining a trade-off between the loss in the sensitivity and the FWHM decrease [74], in a similar way to other configurations, based on different modal filters [75].

Furthermore, low-cost polymeric LDF have been also explored to produce SPR-based sensors [72]. In particular, by sputtering a gold nano-film on the PMMA-based LDF, the obtained bulk sensitivity ranges from 1000 to almost 3000 nm/RIU in the refractive index range from 1.332 to 1.392. In the same work, a novel kind of modal filter has been shown, accomplished by covering the PMMA-based LDF with an aqueous solution to improve the signal-to-noise ratio (SNR) of the SPR sensor, thanks to the filtering of the higher modes, without losing its sensitivity.

3. Packaging of Optical Fiber Biosensors

In order to move from the optical biosensing system to practical devices, miniaturization, portability, and easy handling are required. Therefore, strategies for the encapsulation and packaging of the optical components should be considered in order to meet the demands of the final device, namely by preserving the sensing properties. Fiber-optic biosensors are designed to perform real-time detection in medical applications; the possibility to design packaging structures and embed the sensing elements within medical devices allows performing subcutaneous, percutaneous, or endoscopic operations [76]. In addition, wearable technologies designed with optical fiber biosensors embed the fibers within fabric, forming permanent sensing structures [77]. Figure 10 displays some of the options for packaging, targeting different in situ sensing applications.

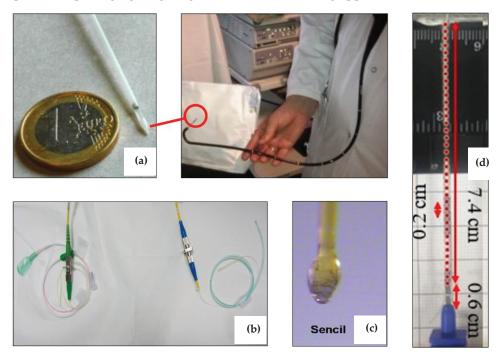


Figure 10. Example of proposed packaging for fiber optic biosensing applications: (a) Endoscopic packaging for in situ detection of cancer biomarkers (image reproduced from [78]); (b) Urologic packaging, with pressure sensors embodied in a rectal catheter (left fiber) and an abdominal catheter (right fiber) (image reproduced from [79]); (c) SencilTM (sensory cilia) packaging for subcutaneous fluorescence sensing (image reproduced from [80]); (d) Epidural needle for percutaneous insertion embodying a network of fiber sensors (image reproduced from [81]).

Endoscopic packaging has been proposed by Guo et al. [78] (Figure 10a) and Loyez et al. [82] for in situ detection of cytokeratin-7 cancer biomarkers. The proposed design integrates the plasmonic biosensor used for detection into a commercially available endoscope that maintains its original functions and allows the precise positioning of the device into the target zone. The work reported in [82] integrates the sensor in a 1.2/1.6 mm (inner/outer diameter) hollow endoscope made with a thermoplastic that maintains a semi-rigid profile. The catheter is terminated by a conical section that allows penetration

into the tissue. The active part of the sensor is exposed to the tissue through a fence that allows direct contact between the sensor and its coating and the detection area. The in situ detection reported in this work allows for obtaining discrimination between healthy and tumoral cells.

An example of packaging suited for the urologic diagnostic is shown in Figure 10b [79]; other notable packaging options have been reported by Poeggel et al. [83] and Guo et al. [84]. Urologic diagnostic requires the measurement of biological [84] or biophysical [83] parameters within the bladder and requires additional elements to compensate for external factors, usually requiring a rectal catheter. Urologic devices are flexible silicon rubber hollow catheters that are positioned by the surgeon within the bladder through the urinary tract. The sensors can detect the surrounding environment through a set of side holes in the proximity of the tip.

A form of less invasive packaging, suitable for subcutaneous sensing (underneath the skin), has been reported by Liao et al. in [80] (as shown in Figure 10c) and [85]. The packaging form proposed in these works is labelled as SencilTM, or sensory cilia. The sensor is based on quantum-dot fluorescence made with multimode fiber. The low-cost packaging is based on immersing the sensing tip in a hydrogel: by using a polyethylene glycol acrylate precursor, the fiber forms a thicker tip that can sustain the penetration underneath the skin, reaching a subcutaneous position where the fluorescence-based glucose measurement can be performed.

Percutaneous packaging formats are in use in biomedical applications that require deep-seated diagnostic or therapeutic devices [76]; in this scenario, a rigid needle is used to pierce through the skin for several centimeters, reaching the target area by a straight penetration. Epidural sensors, such as the implementation shown in Figure 10d, proposed by Issatayeva et al. [81], and reported in [86], make use of a Tuohy needle bent around the tip and with fibers mounted on the outer surface.

The packages proposed in [76–86] integrate the sensing tasks with additional purposes, such as delivering a therapy [76], tracing the needle position through X-ray imaging [77], or performing differential diagnostic [83], to form multifunctional, theranostic medical devices.

The integration of sensors in wearable technologies [77], on the other hand, poses the challenge of integrating one or multiple sensors on a device that is worn for a long term. Wearable biosensors based on optical fibers follow two main approaches for fiber packaging [87]. A first method relies on the integration of the fibers into an external element, to be worn separately from the apparel: Issatayeva et al. [81] integrated an array of sensors in two wearable bands, designed to transfer the strain due to respiratory movements to the sensors.

However, the main trend in wearable devices is to integrate the fibers directly on clothing, where they can continuously monitor vital parameters during sports activities or specific actions. Lo Presti et al. [88] proposed a sensory system based on several gratings integrated into a t-shirt for the measurement of physical parameters during daily activities. Li et al. [89] performed a study of the integration of sensors into textiles by sewing an optical fiber through a piece of fabric and allowing a highly efficient thermal transfer of the body temperature to the sensor. In addition, Esmaeilzadeh et al. [90] reported a biosensor based on the surface plasmon resonance, having the sensing element integrated within a textile for contact sensing.

4. Smartphone Fiber Optic Sensors

Smartphones are ubiquitous pocket devices that offer unprecedented diagnostic opportunities for the real-time detection of vital parameters [91]. Modern off-the-shelf phones have several instruments that can be used for sensing purposes [92], and mobile apps can be developed to use each component of the phone to create sensing systems. In addition, apps can use features such as geolocalization, internet of things, or access to external machine learning methods [93] to enhance the potential for diagnostic and biomedical devices.

The potential of smartphones has been exploited within a number of optical sensors [94–98]; the main purpose of the research carried out by several groups in the last few years is driven by replicating diagnostic methods based on fluorescence spectroscopy [94], fluorescence resonance energy transfer [95], colorimetry [96], or spectrophotometry [97], among others, but by replacing the hardware with the existing components on the smartphone, with the minimal addition of external hardware [98].

Overall, optical biosensors based on smartphone platforms are an excellent diagnostic tool and use mounts that can be used as phone accessories in order to integrate the pre-existing devices of the smartphones (such as the flashlight LED or the camera) with additional components such as lenses, spatial filters, diffraction gratings, or additional laser sources [99,100].

Within optical sensors for smartphones, a new generation of optical fiber sensors for smartphones has been developed and reported within the latest years, and it is consistently being applied in the detection of biological analytes, biohazards, and biophysical parameters. While smartphone optical sensors can analyze previously collected samples or can inspect samples positioned in the proximity of the phone's camera, optical fiber sensors for smartphones introduce the remote sensing possibility. Thanks to the low losses of optical fibers and the availability of multimode fibers that operate in the visible wavelength range, it is possible to design systems that implement the remote sensing feature that is typical of an optical fiber sensor, but where the hardware is partially or entirely replaced by the smartphone itself.

Figure 11 illustrates the main implementations and applications of smartphone optical fiber sensors.

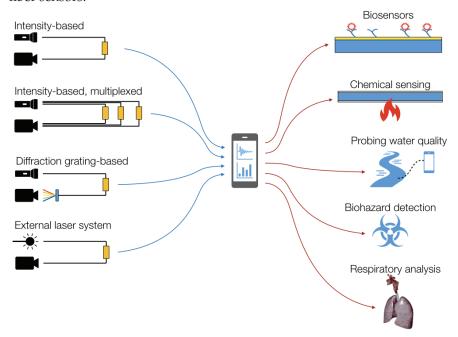


Figure 11. Illustration of the main architectures (**left**) and applications (**right**) of optical fiber biosensors based on smartphones.

4.1. Intensity-Based All-Fiber Smartphone Sensors

Intensity-based sensors represent the simplest implementation, as they can be defined as all-fiber sensors; in these architectures, reported in [101–103], the sensing element modulates the light transmitted through the fiber, and both the light source (flashlight LED off) and the light detector (camera) belong to the phone. The only external hardware is a 3D-printed connector that holds the fibers in place and performs light coupling from the flashlight to the input fiber and from the output fiber to the camera.

All systems based on intensity-varying sensors operate with POF fibers, with a core/cladding diameter of 0.98/1 mm and numerical aperture (NA) of 0.47; this wide

acceptance angle allows easy coupling of the light in and out of the fibers, delivering a sufficient amount of light to the detector.

The first system was reported by Sultangazin et al. [101] for biochemical detection of hydrogen sulfide (H_2S), which is a potentially dangerous biohazard for oil and gas and mining workers; the system uses the flashlight LED from the phone (a phosphor white-light LED illuminating the whole visible wavelength range) and the camera as a pixel-based intensity detector, which converts the output beam from the fiber into greyscale intensity level, and integrates this on the whole set of available pixels. The working principle is based on the change of attenuation in an Ag-coated POF when exposed to hydrogen sulfide; the formation of silver sulfide (Ag_2S) decreases the wall reflectivity of the outer fiber, increasing evanescent losses and, therefore, reducing the transmission. By integrating, via an app developed in Java, the power incident on the camera, Sultangazin et al. reported an intensity drop of 0.54%/min to 4.9%/min when the sensor is exposed to hydrogen sulfide.

A similar concept was reported by Aitkulov and Tosi [102], extending the sensing mechanism to biophysical sensing by measuring the breathing rate pattern. The sensing element is constituted by a pair of POF fibers, cleaved and spaced by a set distance; when fibers are aligned, they exhibit the maximum transmission, while during breathing movements, the angle between the transmitting and receiving fiber increases the coupling losses, decreasing the detected power. The sensing element is then embodied in an elastic band and worn on the chest for a wearable application. By means of a Fourier analysis of the obtained intensity time signals, the respiratory rate is determined.

The first attempt at a multiplexed fiber-optic sensing system for smartphones was reported in [103]; by merely changing the connector type from the prior system [102], the authors demonstrated a three-fold increase in sensing points. The principle of operation relies on splitting the LED light into multiple fibers and collecting each output into a different portion of the camera. This setup, labelled camera-division multiplexing, allows for isolating each light spot collected on the camera by a tessellation, with minimal crosstalk.

All-smartphone-POF systems [101–103] make full use of the phone's optical circuitry and application programming interfaces for each component. By setting the exposure time, controlled by the ISO of the camera, and disabling the automated routines for the flashlight, it is possible to detect the intensity through photo or video acquisition [102,103] or through a separate mobile app [101]. In this case, no further components are needed to implement the sensing mechanism, as the software can set ISO and calculate the intensity at saturation point, adapting to the power level at the detector. The only external device serves the purpose of mating the phone to fiber and is a phone accessory that can be 3D-printed to be customized to each smartphone model.

4.2. Diffraction-Grating Assisted Smartphone Optical Fiber Sensors

The main alternative to power detection via smartphone is to use the camera of the phone as a spectrometer and detect the wavelength shift of spectrally sensitive elements such as surface plasmon resonance (SPR) or fiber Bragg gratings (FBGs). In order to convert the camera to a spectrometer, a diffractive element, such as a grating, or a low-cost element, for instance, a piece of DVD disc, can be used. Thanks to this arrangement, it is then possible to measure the wavelength shift of transmission spectra, making a sensing system that is insensitive to power level and, therefore, potentially more robust to the fiber coupling.

Bremer et al. [104] presented the first fiber-optic RI sensor adapted for smartphones by means of the SPR effect. The system proposed in this work converts the phone camera to a spectrometer by means of a diffractive element, a holographic grating with 1200 lines/mm mounted at a 45° angle with respect to the camera lens. Then, by detecting the output spectrum sampled by each array of pixels, the spectrum of the SPR sensor was detected, obtaining a sensitivity of 1678 pixel/RIU after optimizing the resolution of the phone's camera.

The system proposed by Bremer and Roth was adapted in [105] to the detection of SPR in a planar waveguide configuration. Here, the fiber sensor is replaced by a tapered waveguide geometry with an initial thickness of 1200 μ m, reduced to 200 μ m in the proximity of the sensor. Two SPR sensors have been designed in the waveguide system, enabling multiplexed detection; the waveguide is coated with a thin film of gold in order to immobilize aptamers, which are used as specific bioreceptors for the detection of 25-hydroxyvitamin D. The waveguide system demonstrates the possibility to detect this vitamin D with a device mounted as an external accessory for the phone.

A different approach was proposed by Markvart et al. [11,106], who showed the possibility to interrogate traditional wavelength-selective fiber-optic sensors, such as interferometers [106] or gratings [11], using an architecture based on the smartphone. The first system, shown in Figure 12, is designed for the interrogation of Fabry–Perot (FP) interferometers. The system uses a diffraction element in front of the camera, collecting the light at a 60° angle. The spectrometer is used as a colorimeter through a signal processing method that converts the output image into the corresponding transmission spectrum by means of signal processing.

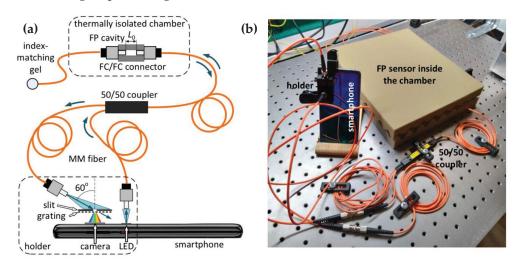


Figure 12. Schematic representation (**a**) and photograph (**b**) of the smartphone optical fiber sensing system for Fabry–Perot sensor interrogation (image reproduced from [106]).

A similar system was reported in [11] for the interrogation of chirped fiber Bragg grating sensors, inscribed on a multimode fiber for strain and temperature sensing. The chirped grating, with bandwidth ~ 8 nm at the red wavelength, exhibits a wavelength shift that can be detected by estimating the pixel-based spectrum via signal processing. Furthermore, Pan et al. [107] reported a temperature sensor based on a multimode fiber, detected in reflection through a smartphone system similar to the previous ones.

4.3. Smartphone Optical Fiber Sensors with External Sources

While the previous architectures makes use of both the light source and detector of the phone, other systems proposed in [107–111] maintain a fiber-optic probe but use a different light source. In these cases, the system is still fiber-coupled and compatible with a smartphone, but significant external hardware is required.

Kamizi et al. [107] developed a fiber-optic sensing network that monitors walking patterns. To do so, a series of single-mode fibers have been placed underneath insoles, with sensors acting as macro-bending units that increase the propagation losses when pressed during the walking motion. The system relies on external hardware that incorporates the LED light source and all the coupling mechanisms, while the output beam is detected through the camera of the phone.

Liu et al. [109] reported a fluorescence fiber-optic sensing system based on a smartphone assisted by an external laser source. The system uses a red laser in order to excite fluorescence in a multimode glass fiber functionalized with quantum dots. The reflected light is collected by the phone's camera by means of an external module customized for the smartphone's size and camera position. An Android app is used to calibrate the system and collect the data. Here, the proposed application is the detection of mercury cations (Hg^{2+}) in the 1–1000 nM range. A similar system was described in [110] for pH detection and in [111] for dual-channel fluorescence monitoring.

The different main implementation strategies of smartphone-based optical fiber sensors are summarized in Table 1.

Table 1. Review of the fiber optic smartphone sensors and their applications.

Ref.	Application	Smartphone	Operative System	Optical Fiber	Internal Devices	External Devices
[101]	Remote detection of hydrogen sulfide	Samsung Galaxy S6 Edge	Android 7.1.1	POF, Ø0.98/1 mm, NA 0.47	Flashlight, camera (set ISO)	3D-printed connector
[102]	Breathing pattern detection	Redmi Note 4	Android 7	POF, Ø0.98/1 mm, NA 0.47	Flashlight, camera (set ISO)	3D-printed connector
[103]	Multiplexed Breathing pattern detection	Redmi Note 4	Android 7	3xPOF, Ø0.98/1 mm, NA 0.47	Flashlight, camera (set ISO)	3D-printed connector
[104]	SPR sensor for refractive index	Huawei Ascend Y300	Android 4.1.1	Thorlabs BFL48, Ø400 µm, NA = 0.48	Flashlight, camera	Diffraction grating, couplers
[105]	Vitamin D Detection	Apple iPhone 6s	iOS	Tapered waveguide, 200–1200 µm	Flashlight, camera	Diffraction grating, external mount
[106]	Interrogation of Fabry–Perot sensor	Huawei P20 Pro	Android	Graded-index MMF, core Ø62.5 μm	Flashlight, camera	Slit, grating, 2× FC/PC connectors
[11]	Interrogation of chirped fiber Bragg grating	Huawei P20 P50	Android	Graded-index MMF, core Ø62.5 μm	Flashlight, camera	Slit, grating, 2× FC/PC connectors
[107]	Temperature Sensing			Multimode glass fiber, Ø300 μm	Flashlight, camera	Diffraction grating, pedestals, couplers
[108]	Identification of walking pattern	Motorola, Moto G3 Turbo	Android	Glass single-mode fiber	Camera	External phone mount, LED
[109]	Detection of mercuric cations	Nubia Z17 Mini	Android	Multimode glass fiber, Ø105/125 μm	Camera	Red laser, connecting module
[110]	On-site pH detection	Nubia Z17 Mini	Android	Multimode glass fiber, Ø105/125 μm	Camera	Red laser, connecting module
[111]	Dual-channel fluorescence detection	Nubia Z17 Mini	Android	Multimode glass fiber, Ø105/125 μm	Camera	Red laser, connecting module

5. Biosensing Applications

As introduced before, due to their unique characteristics, OFBs have enormous potential to be used as devices for POC monitoring of biochemical components in diverse fields. This review is focused on three areas that are of utmost importance for the health and

wellbeing of the population in general: cancer and cardiovascular biomarkers detection and environmental monitoring.

Real-time, label-free, and high specificity towards a biomarker of interest are a few of the fundamental features of the modern biosensor [6]. Among other essential traits are the miniaturized size of the device and high sensitivity [6]. In the sensing technology, optical fibers demonstrate adaptability for the recognition of numerous analytes and allow the combination of diverse detection techniques. They can monitor refractive index change in the media from basic buffers and water to a complex serum, urine, and blood when detecting proteins, small biomolecules, cells, and other biomarkers [112]. Thanks to strong evanescent fields, optical fiber biosensors modified with biorecognition elements are a promising tool for sensing target molecules [113]. With an increasing demand for enhanced selectivity, specificity, and sensitivity, it is necessary to continue improving optical systems for biosensing applications.

5.1. Cancer Biomarkers Detection

Detection of cancer biomarkers plays a significant role in the prognosis and treatment of various cancers [114]. At the early stage of illness, the concentration of biomarkers is extremely low, and this requires highly sensitive diagnostics methods with very low LODs. Although there is ongoing progress in the standard detection and treatment methods, the mortality due to advancement and metastasis of cancer is still high [115].

Conventional detection procedures for identifying cancer biomarkers require patients to provide a biopsy of the tumor. Subsequently, the tested sample needs to undergo either fluorescence in situ hybridization (FISH) or immunohistochemistry (IHC) [116,117]. Both testing procedures involve the use of fluorescent tags to label chromosomes or DNA sequences [118] in the former and proteins [119] in the latter case. Since the above testing methods are expensive, take long processing times, and, above all, involve invasive procedures [120], alternative measures need to be called for. The research in the field of cancer detection using optical fiber sensors has made considerable progress where non-invasive, real-time, label-free methodologies are proposed [50,121–124]. Different optical fiber based systems have been designed and developed to specifically detect cancer biomarkers, aiming towards a fast diagnostic method for POC testing, as will be presented in the following sections.

5.1.1. Human Epidermal Growth Factor Receptor 2 (HER2)

HER2 (185 kDa) is part of the human epidermal growth factor receptor family, which participates in cell growth and differentiation. However, its overexpression plays an active role in the inception and proliferation of aggressive forms of breast cancer [125]. Patients who undergo costly trastuzumab treatment have a high chance of cardiac toxicity. Thus, taking into account the risk associated with the above therapy, detection of the biomarker is crucial for the prediction and therapy of tumors [126].

In 2017, Sun et al. [124] reported the detection of HER2 on a taper interferometer embedded in FBG. The device was sensitive to RI changes, and wavelength shift was recorded while detecting HER2. The lowest concentration of HER2 detected was at 2 ng/mL. Moreover, the FBG part of the compact device was suggested to be used as an independent thermometer during the sensing process.

An optical fiber-based surface plasmon resonance (OF-SPR) optrode for detection of HER2 was proposed by Loyez et al. [50]. In the experiment, a multimode optical fiber with a diameter of 400 μ m was sputtered with gold (thickness 45 nm), followed by silanization and attachment of anti-HER2 aptamers. HER2 biomarkers were detected in phosphate buffered saline (PBS) solution presenting a LOD of 0.6 μ g/mL. CK17 protein was tested as a negative control, and no significant shift was observed. Moreover, a sandwich assay experiment using anti-HER2 antibody further improved the LOD to 9.3 ng/mL (~86 pM).

5.1.2. Cluster of Differentiation 44 (CD44)

CD44 is a transmembrane protein with a molecular weight of 85–200 kDa present in human cells, embryonic stem cells, cancer cells, and cancer stem cells [127]. Its overexpression promotes the migration, proliferation, metastasis, and reoccurrence of cancer [128]. Numerous variant forms of CD44 (CD44v) were confirmed to be related with several malignant tumors and metastasis [129]. Additionally, there is a correlation between the role of CD44 in circulating stem cells and the impact it has on the development of cancer and therapeutic outcome [130].

In a recent work from Bekmurzayeva et al. [131], a spherical fiber optic tip fabricated using a $\rm CO_2$ laser splicer was investigated. The fiber surface was pre-treated with (3-aminopropyl)trimethoxysilane (APTMS), following several steps to functionalize it with anti-CD44 antibodies, as shown in Figure 13, aiming to detect various levels of CD44. The determined LOD was of 17 pM. In order to demonstrate the specificity of the studies, thrombin and interleukin 4 were used as controls, reporting a minor change in signal.

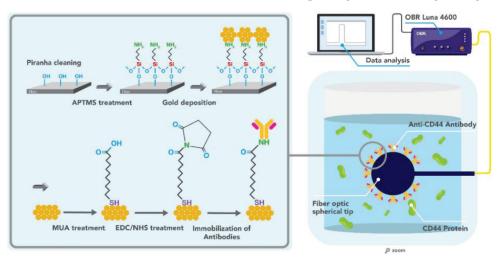


Figure 13. Schematic of functionalization of spherical tip-biosensor used for detection of CD44 biomarker (image reproduced from [131]).

5.1.3. Thyroglobulin (Tg)

Differentiated thyroid cancer (DTC) is one of the most common types of endocrine cancer [132]. While the progression rate and death incidence for DTC are indolent, the relapse cases after primary therapy with the need for thyroidectomy are high, which are related to the increased instances of long-term morbidity [133–135]. Hence, a lifetime check-up is advised for timely diagnosis to avoid further progression and subsequent death [136]. Thyroglobulin (Tg) is a dimeric precursor glycoprotein weighing 660 kDa, made by the thyroid follicular cells, and is necessary for thyroid hormone synthesis (T_3 and T_4). Its levels can be measured in the serum of a healthy person up to 40 ng/mL. The elevated amount of Tg, which is a highly specific biomarker to DTC, is an indication of the latter [137].

In work performed by Kim et al. [138], a fiber-optic LSPR sensor combined with a microfluidic channel was proposed to detect Tg. One of the preparation steps for optical fibers involved treating them with Au NPs (average diameter of 50 nm). The device was exposed to a range of concentrations of Tg, giving the LOD of 93.11 fg/mL. Comparatively, the response time was 10 min, whereas conventional techniques such as chemiluminescence assay (CLIA), immunoenzymometric assay (IEMA), and immunoradiometric assay (IRMA) take a day or more [138].

5.1.4. Cytokeratin 7 (CK7)

There is a discrepancy between the biomarkers expressed by tumors at the early and metastatic stages. Cytokeratin 7 (CK7), which is usually expressed by epithelial cells coating

the cavities of internal organs, can assist in the diagnosis of carcinomas. Cancer cells of the lung, breast, thyroid and salivary glands, and female reproductive organs will express CK7. Besides finding the biomarker in primary lung cancer, fragments of cytokeratin can be released into the bloodstream by malignant cells [139].

Ribaut et al. [123], detected CK7 protein (78 kDa) and CK7 (2.6 kDa) peptide using an SPR optical biosensor based on TFBG. In this work, TFBGs with tilt angles between 7° and 9° were coated with gold (thickness of 50 nm) using a high vacuum sputter followed by surface pre-treatment and immobilization of anti-CK7 antibodies. Detection of both CK7 protein and CK7 peptide attained a LOD of 1 pM in PBS and 0.4 nM in serum, respectively.

5.1.5. Cytokeratin 17 (CK17)

Nearly a third of cases of cancer mortality worldwide are caused by late diagnosis and aggressive growth of lung cancer [140]. The average span of life after confirmation of disease and its subsequent treatment with chemotherapy is less than a year [141]. To enable early-stage diagnosis and to enhance recovery and survival rate, the identification of biomarkers is highly in demand, since various types of lung cancer require various treatment procedures. Generally, an overexpression of cytokeratins is associated with carcinogenesis [142,143]. Among them, Cytokeratin 17 (CK17) promotes cancer progression and epithelial proliferation [144].

Ribaut et al. [121], in 2017, reported a gold-coated (50 nm) TFBG immunosensor with 7° for detection of the CK17 biomarker. The sensor surface was functionalized with a self-assembled monolayer of S_2 -PEG₆-COOH followed by immobilization of the anti-CK17 antibody. Furthermore, the sensor was packaged into a hollow cylindrical needle, and a window was cut with a laser to expose the fiber to the surrounding environment. The packaged biosensor was used to detect biomarkers in PBS and CK17-encapsulated gel matrix, giving a LOD of 1 pM and 0.1 ng/mL, respectively. In addition, the recognition of analyte was conducted ex vivo in healthy and tumoral tissue samples. Compared to healthy tissue, the sensor's response to protein showed a linear trend in a biopsy sample of CK17+ positive patients.

Similarly, Loyez et al. [122] detected CK17 comparing four biofunctionalized TFBG based sensors (without gold—using silanization, gold-coated via sputtering, gold-coated with electroless deposition (ELP), and hybrid gold-coated—a mix of sputtering and ELP) that underwent various pre-treatment procedures. Sensing of analyte showed a linear trend in all four TFBGs. Fiber prepared using silanization showed the lowest sensitivity comparing to the others. However, it was stated that it could be used without taking into account polarization effects. The LOD for the gold-sputtered TFBG was close to 1 pg/mL, and it was the most sensitive device. The sensor was further inserted in the tumorous part of the lung biopsy, showing a positive trend for CK17. Fibers were gold-coated using ELP and hybrid methods and were both sensitive to the analyte, with the former showing sensitivity down to 1ng/mL and the latter being as sensitive as the gold-sputtered TFBG.

5.2. Cardiovascular Biomarkers Detection

Cardiovascular diseases (CVDs) are the prime cause of death in the world. It is estimated that, in 2019, more than 19 million people died from CVDs, corresponding to 32% of all global deaths [145]. Most CVDs are related to lifestyle and modifiable factors [146]. Therefore, prevention is the greatest weapon for reducing CVDs and improving wellbeing. Factors such as stress, sedentary lifestyle, obesity, metabolic diseases, and hypertension are preponderant in the development of such epidemy [147]. Hence, their monitoring and control play a fundamental role in the decreasing of CVDs. Population-level interventions always have to consider the cost-effectiveness of the applied measures and monitoring techniques [147]. In a disease with such a large prevalence, this aspect is of utmost importance.

In CVD management, each person is different, and individual assessment should be personalized according to age, co-morbidities, and lifestyle. Moreover, monitoring should be cost-effective and, when possible, at point-of-care [148]. Recently, different cost-effective optical fiber sensors were developed to monitor different cardiovascular biomarkers, namely, blood pressure [149], cholesterol levels [150] and diabetes mellitus (glucose levels) [151,152], heart acute myocardial infarction (AMI) by Troponins quantification [151,153], heart and kidney failure [154], and stress [21].

5.2.1. Detection of Cholesterol and Glucose

Dyslipidemia, which is the presence of abnormal lipidic levels in the blood, and glucose metabolism disorders, such as diabetes, are independent risk factors of CVDs [147]. Therefore, monitoring of such analytes in POC locations is of outmost importance in CVD prevention. Kumar et al. [150] developed a label-free reflectance-based cholesterol biosensor consisting of SMF with a hollow core fiber working in reflection, coated with Au NPs and with diameters of around 11 nm, for sensitivity enhancement using LSPR phenomenon. The biorecognition molecule was cholesterol oxidase enzyme. With this approach, a LOD of 25.5 nM and a sensitivity of 16.15 nm/µM were attained. In 2019, an optical tapered fiber structure coated with Au NPs, working on transmission, was developed [155]. The taper, created on SMF fiber, achieved a minimum diameter of 40 µm, a total length of 4 mm, and a 5 mm transition stretching region on each side of the taper. The biorecognition molecule was glucose oxidase (GOx) enzyme. This sensing approach was studied in a detection range up to 10 mM, achieving a LOD of 322 μM and a sensitivity of 0.93 nm/mM. In 2020, Zheng et al. [156] described a highly-sensitive glucose sensor, based on back-reflection configuration, using plastic cladding fiber of 600 µm. The fiber tip was coated with gold by sputtering and was then biofunctionalized with GOx enzyme. A measurement range of 0–0.5 mg/mL and a resolution of 0.0004 mg/mL were achieved.

5.2.2. Detection of Acute Myocardial Infarction Biomarkers

Acute myocardial infarction (AMI) refers to the death of myocardial cells due to ischemia or the imbalance between the blood supply and demand within the coronary arteries as a result of an acute thrombotic process [157]. As a life-threatening situation, the fast diagnosis of AMI is vital for the early initiation of appropriate therapeutic measures. The main tools to diagnose AMI are electrocardiography (ECG) and cardiac troponins detection [158]. Cardiac troponins are regulatory proteins specific to the myocardium that are released into the circulation when myocytes are being damaged [157]. There are three kinds of troponins (Tn): I, T, and C, with troponin T and I being cardiac-specific, while troponin C is expressed by both cardiac and skeletal muscle.

An ultrasensitive label-free optical microfiber coupler biosensor based on interference turning point effect was developed to detect cardiac troponin I. The interrogation setup consisted of a halogen light source, with a microscope objective to focus the light into the fiber. The sensor consisted of a microfiber coupler fixed in a fluid cell integrated into a polydimethylsiloxane (PDMS) chamber for sample solutions delivery. A visible spectrometer was used as a signal analyzer. With this sensing architecture, an RI sensitivity of 91,777.9 nm/RIU was achieved, and after biofunctionalization with anti-cTnI antibodies and testing in cTnI solutions, it reached a LOD of 2 fg/mL [151]. Krupin and Berini [152] developed a long-range surface plasmon-polariton (LRSPP) waveguide biosensor to detect troponin I. The sensors consisted of gold stripes of about 35 nm thickness, embedded in a low-index optical-grade fluoropolymer (CYTOPTM) with fluidic channels etched to the surface of the gold strips. LRSPPs were excited by butt-coupling a polarization-maintaining single-mode fiber (PM-SMF) to the input facet. The optical interrogation included an LED peaking at 1310 nm connected to the PM-SMF. The transmitted light through the LRSPP passes to a $25 \times$ objective lens and a 50:50 beam splitter. One part of the split beam is sent to an infrared camera for visual monitoring and alignment, and the other part is sent to a photodetector for evaluation of power variations in time during the detection procedure. The gold waveguide was biofunctionalized with anti-cTnI antibody via Protein G and passivated with BSA. The sensor was tested in a direct and sandwich detection mode, reaching a LOD of 430 pg/mL and 28 pg/mL, respectively.

5.2.3. Heart and Kidney Failure

Heart failure is a cardiovascular disease characterized by a disorder of the heart, that can be structural and/or functional, which can cause increased intracardiac pressure and/or inadequate cardiac output [147]. In heart failure patients, kidney function has to be closely examined, since renal dysfunction due to diuretics, particularly in patients with heart failure, is a common cause of hospitalization [159]. Botewad et al. [154] developed a biosensor for urea detection, a biomarker of kidney function [160]. This study used as transducer a 450 µm diameter core plastic-clad fiber (PCF) uncladded throughout a 2 cm portion. The uncladded region of the fiber was modified with a composite of polyaniline (PANI) with ZnO and biofunctionalized with urease enzyme, which catalyzes the hydrolysis of urea. This intrinsic sensor worked in transmission using a halogen lamp and a spectrophotometer as light source and detector, respectively. When testing in urea solutions, this biosensor reached a LOD of 10 nM.

In 2021, Li et al. [161] reported a biosensing approach comprised by SMFs, a multicore fiber (MCF), and MMF fibers, in the structure SMF–MCF–MMF–SMF, for the detection of creatinine in the human body. This structure, produced by a fusion splicer machine, was etched to a 90 μ m diameter and coated with graphene oxide (GO), Au NPs, and molybdenum disulfide NPs (MoS₂-NPs) and then biofunctionalized with creatinase enzyme. When tested in creatinine solutions, the sensor presented a LOD of 128.4 μ M in a linear range of 0–2000 μ M.

In heart failure, higher enzymatic activity with maintained oxygen consumption contributes to the physiopathology of myocardial insufficiency and appears to be an indicator of oxidative stress [162]. In 2021, Ortega-Gomez et al. [51] developed a plasmonic tip biosensor for the detection of reduced cytochrome c, which is a multifunctional enzyme with a crucial role in electron transfer in the mitochondrial transport chain. This biosensor consisted of a PCF MMF uncladded tip, that was coated with Au NP biofunctionalized with cytochrome c as the biorecognition molecule, achieving a LOD of 60 nM.

5.2.4. Stress

Stress conditions (depending on the degree, duration, and individual response) often lead to maladaptive physiological responses and are associated with CVDs in their acute and chronic forms [163]. Stress also increases the probability of developing CVD risk factors such as hypertension, diabetes, and obesity [164]. Prolonged high levels of cortisol, known as the stress hormone, are related to physical and psychological disorders [165]. Therefore, the development of an easy and low-cost POC monitoring method to detect this stress biomarker is a current hot topic. In 2021, C. Leitão et al. [21] reported a POF uncladded approach for the detection of cortisol. The biosensor used SPR as a sensitivity enhancer by coating the unclad part of the fiber with a gold/palladium (AuPd) alloy by the sputtering technique. The surface was biofunctionalized with anti-cortisol antibodies using cysteamine as an intermediated linker. The final surface was passivated with BSA. When tested in cortisol solutions of concentrations between 0.005 and 10 ng/mL, the proposed sensor had a total of 15 nm wavelength shift. The attained sensitivity and LOD were 3.56 ± 0.20 nm/(log(ng/mL)) and 1 pg/mL, respectively. In this research, control tests were also performed in a sensor functionalized with antibodies for human chorionic gonadotropin (anti-hCG antibodies) in which the variance of the resonance wavelength was only 1 nm, much lower than the sensor modified with anti-cortisol antibodies.

5.3. Environmental Monitoring

Today, more than ever, there is a growing concern about emerging contaminants and pollutants which severely affect the environment and human health. These contaminants and pollutants, which include insecticides, gases, stimulants, antibiotics, anti-inflammatory

drugs, and pesticides, among others, are biologically dynamic and highly resistant. As a result, they continue in the environment, instigating harmful effects to non-target organisms and humans, making it imperative to find approaches capable of correctly detecting these substances in order to protect the environment and our lives.

In agricultural and fisheries engineering, a variety of environmental monitoring sensors (quality of air, soil, and water) allow us to understand the environmental impact of agricultural activities and recognize whether the farm conditions are suitable to cultivate. Many studies have focused on the monitoring of farm soil and aquaculture water, since soil fertility is the key to farmers looking for improvement of crop yields and agricultural productivity, and well-controlled fish water tanks are key to the fishery sector. Further, oil/gas industry pieces of equipment (pumps, pipes, and joints, among others) are required to be monitored. In case any of these become damaged due to natural disasters or human intervention, not only are their properties lost but also the surrounding environment can become polluted. The equipment is submerged at the sea bottom or into a downhole, which are high-pressure and environments with corrosive saltwater. In the case of delivery pipes, these are spread over several kilometers and buried underground. This means that oil and gas industries need non-invasive and low-cost sensing solutions to monitor many critical parameters such as gas detection and/or leakage and oil quality, among other things.

5.3.1. Phenolic Compounds

Phenolic compounds are mainly produced by burning wood and coal apart from their existence as sewage and industrial by-products [166]. With water being the most valuable natural resource for mankind, its pollution by organic and inorganic compounds in the current industrial age has been a serious concern and has placed it as one of the world's leading health risks [167]. Phenol wastewater is one of the industrial sewages which causes harm worldwide, with it also being one of the worst sources of environmental water pollution [168]. Drinking water with extremely high concentrations of phenolic compounds can cause muscle convulsions, difficulty in walking, and even death [169]. These phenolic compounds release toxic gases during combustion that are potentially harmful to human health if absorbed. p-Cresol can enter human bodies in different ways as it is used as a flavoring agent in foods and in some traditional medicines and is always present in tea, oil, and tap water [170]. Recently, an optical fiber-based lossy mode resonance sensor built by using NPs of ZnO/MoS_2 and the MIP technique was used to detect *p*-cresol [171]. To achieve the specificity of the sensor, tyrosinase enzyme is commonly used for *p*-cresol detection [172]. Tyrosinase (polyphenol oxidase) is an enzyme that catalyzes phenol o-hydroxylation yielding o-diphenol (monophenolase activity) that subsequently oxidized to o-quinone [173]. Very recently, Wang et al. [174] reported that a localized plasmonbased sensor was developed for p-cresol detection, consisting of a nonadiabatic 40 µm of tapered optical fiber experimentally fabricated and computationally analyzed using the beam propagation method. For performance optimization of the sensor, two probes were proposed, where probe 1 was immobilized with Au NPs and probe 2 was immobilized with the Au NPs along with ZnO NPs. To increase the specificity of the sensor, the probes were functionalized with tyrosinase enzyme. Different solutions of p-cresol in the concentration range of 0-1000 µM were prepared in an artificial urine solution for sensing purposes. Different analytes were prepared for selectivity measurement. The linearity range, sensitivity, and LOD of the probe using ZnO NPs were 0-1000 μM, 5.6 nm/mM (accuracy 0.981), and 57.43 μM, respectively, making the overall performance of such a probe much better, due to the inclusion of ZnO NPs, which increases the biocompatibility of the sensor probe.

5.3.2. Phthalate Esters

Phthalate esters, also known as PAEs, are a type of synthetic chemical substance that has been allocated to the list of priority pollutants because of their endocrine-disrupting and toxic effects on the human body. In 2019, Cennamo et al. [175] reported a D-shaped

plasmonic optical fiber biosensor to detect the presence of naphthalene in sea water. The D-shaped configuration was obtained by removing the cladding of the POF by a polishing process. In this study, an antibody specific to the naphthalene molecule was designed and produced and, for that, a retro-synthetic chemical strategy was applied to modify the NAPHTA structure in a derivative structure to obtain the anti-NAPHTA antibodies. This modified NAPHTA structure was coupled to a protein carrier and was used for immunization. The capability of the antibody to bind to naphthalene was assessed by enzyme-linked immunosorbent assay (ELISA) tests. Through N-ethyl-N'-(3dimethylaminopropyl)carbodiimide (EDC)/N-hydroxysuccinimide (NHS) chemistry, the gold surface was derivatized and functionalized with the produced antibody. Subsequently, tests using real matrices of sea water were performed using the produced biosensor. Through the obtained results, it was possible to observe that the POF biosensor was able to sense the presence of naphthalene in a sea water solution with a LOD of 0.76 ng/mL (0.76 ppb), which is lower than the limit value of naphthalene (0.13 µg/mL). In a 2020 study, Lamarca et al. [176] prepared a label-free U-shaped immunosensor for the detection of ciprofloxacin (CIP) in wastewater samples, as CIP is a broadly utilized antibiotic to treat infections and is a common contaminant of wastewater treatment plants. The glass optical fiber surface was functionalized with PANI, followed by immobilization of anti-CIP antibody. CIP could be detected with a LOD of 3.30×10^{-3} ng/L, in a linear range between 0.01 ng/L and 10,000 ng/L, and with a quantification limit of 0.01 ng/L. In addition, the immunosensor offered a high average recovery of 91%. Benzo(a)pyrene (B(a)P) is one of the most toxic polycyclic aromatic hydrocarbons and a carcinogen, making monitoring its concentration levels essential for human health and environmental contamination avoidance. To monitor B(a)P contamination levels in the water, Gao et al. [177] proposed an in-line fiber optofluidic immunosensor using a hollow-core fiber with its surface immobilized with antibodies. As such, the immunoreaction between the antibody and the B(a)P molecule induced a significant change in the RI inside the in-line optofluidic channel. The attained results presented a LOD of 1.65 pM and sensitivities of up to 23 pm/pM.

5.3.3. Gases and Volatile Compounds

There are myriad gases and volatile materials that researchers and engineers are interested in studying for a number of reasons (ranging from safety issues to general analytical analysis) in the oil/gas industry, from which three will be focused on: formaldehyde (CH_2O) , hydrogen sulfide (H_2S) , and carbon dioxide (CO_2) , where the identification of specific molecules in diverse media and under extreme conditions has been carried out.

González-Vila et al. [178] developed an MIP coating synthesized around a metal-coated optical fiber sensor by an electropolymerization process, working as an electrode, where the oxidation takes place at the surface of the metal. The electrodeposition was figured by the TFBG–SPR sensor and, as a result of the MIP coating, the TFBG-based sensor acquired sensitivity in gaseous atmospheres. A sensitivity of 2.10 pm/ppm when detecting tiny formaldehyde concentrations in the gaseous state was achieved. In addition, the sensor exhibits a selective behavior to this molecule such that the presence of other volatile compounds did not produce a substantial change in the sensor's response.

H₂S is usually found in nature by the decomposition of organic materials, especially in oil and gas production chains. It can also be found in mineral environments such as coal and salt deposits, as well as other mineral extractions containing sulfur. Physical, chemical, and biological agents are the main agents responsible for the formation of this compound. Various systems have been developed using optical fiber sensors due to their large advantages. Sultangazin et al. [101] proposed a low-cost H₂S sensor based on plastic optical fiber functionalized with silver deposition on the fiber's outer surface. The sensor is integrated with a smartphone used as an interrogator unit, and the response time is just over a few minutes. Ke et al. [179] presented an optical fiber evanescent-wave sensor. The sensing probe is fabricated by etching a standard single-mode fiber, where the transmitted optical power centered at 1631.9 nm is monitored by a power meter. As the H₂S gas

increases, more energy of the evanescent wave is absorbed, leading to a reduction in the optical power. Prado et al. [180] very recently presented a new study that considers the advantages of optical fiber sensors in conjunction with the colorimetric detection abilities of surface plasmon resonance in Au NPs, reporting the development of a H_2S detector based on optical fiber coated with Au NPs. The proposed configuration and method used make it possible to detect the presence of H_2S in gaseous systems while operating at room temperature and with important advantages regarding easy production and short response time. The developed detector has the ability to sense H_2S levels in the range of 0.4 to 2.0 ppm at room temperature.

CO₂ is another critical parameter in the oil/gas industry. There are many materials to use with an optical fiber that have a specific reaction with CO₂. A range of different materials has been reported in the literature over recent years, such as xerogels doped with 1-hydroxy-3,6,8-pyrenetrisulfonic acid trisodium salt (HPTS) [181], also known as pyranine. This CO₂ sensor is based on pyranine, a pH-sensitive fluorescent indicator dye. In the presence of CO₂, the dye has an ion transfer that alters the absorption features of pyranine and xerogel at 396 nm and 460 nm, effectively reducing absorption at these wavelengths, causing reactions at room temperature, and having rational response times. Another approach uses oxidation or reduction reactions in which an electron transfer process occurs between a gas and a material. An example of this is the hybrid nickel oxide/reduced graphene oxide (NiO/rGO), which is a structured coating material. The reactionary wavelengths are 670 nm and 771 nm, and it is the combination of the nanostructured material and its chemical composition in the presence of CO2 that triggers the reaction, altering the distribution of the radicals within the materials, changing the electron density and, hence, the permittivity. The result is a small but significant change in the emission wavelengths [182]. A redox reaction using single-wall carbon nanotubes in localized surface plasmon structures has been used for CO₂ detection [183], where the chemical selectivity is conferred in terms of the activation energy, allowing for room temperature operation.

5.3.4. Aquaculture Monitoring

In agricultural and fisheries engineering, many studies have been explored in terms of soil pollution and critical parameters in water fish tanks. When stress is persistent and uncontrollable, it is considered pathological, which can trigger depression and cardiovascular diseases [184], for example, and therefore the development of technology capable of monitoring stress is essential. Stress involves a large number of neuronal circuits and, once it is promoted, leads to the release of glucocorticoids, in particular cortisol [184]. The substantial variation in this hormone occurs due to exposure to psychological, environmental, or emotional stress [185]. As a result, cortisol is one of the most important stress biomarkers.

One area in which stress also presents significant influence is aquaculture; especially, the impact of such a stress hormone is a challenge to be overcome in recirculating aquaculture systems. When small variations in the water composition or quality occur in these systems, stress induction can arise as well as reduced food intake, which leads to reduced fish growth and, consequently, leads to possible mortality when acute or chronic stress is high. For these reasons and others, it is then essential to monitor cortisol in water [186]. In 2020, Sharma et al. [187] simulated an SPR fiber optic immunosensor for cortisol detection at the wavelength of 830 nm. The sensor consisted of an Ag layer with 2D materials, conventional (graphene, tungsten disulfide (WS₂), and MoS₂) and transition metal carbides (MXenes: Ti_3C_2 , $Ti_3C_2O_2$, $Ti_3C_2F_2$, and $Ti_3C_2(OH)_2$), considered one at a time. The sensor that showed a superior balanced set of performance parameters under both modes was the $Ti_3C_2O_2$ -based probe. Through simulation, this probe was able to achieve a LOD of 15.7 fg/mL. The cortisol sensor developed by Leitão et al. [21], described in the last section, also can be applied to monitor cortisol levels in aquaculture water.

One of the key limiting factors in aquaculture is the presence of ammonia; therefore, its early detection in small concentrations prevents fish mortality and improves the production

quality [188]. Yi Zhu et al. [189] reported a combination of optical fibers and tapered optical fibers (SMF-MMF-taper-MMF-taper-SMF) as miniature interferometry-based optical fiber ammonia gas sensors. A range from 0 to 5460 µg/L of ammonia in the gas chamber was performed. The sensing material involved ZnO nanoflowers deposited on the sensing area (middle MMF and tapers), by a drop of a ZnO solution and 6 h drying in a vacuum oven at 60 °C. Ammonia sensitivity performance is compared between ZnO nanoflowers and ZnO nanospheres of around the same size (1 μ m), showing sensitivities of 5.75 pm/(μ g/L) versus ~2.25 pm/(μg/L), respectively. Shrivastay et al. [190] presented a contemporary approach to design and developed a hypersensitive ammonia gas sensor producing a Mach–Zehnder interferometer (MZI) by an SMF–PCF–SMF fiber substrate to perform the interference by immobilizing PANI@SnO₂ nanocomposite to achieve sensing. In this case, excitation of core and cladding modes of PCF is achieved using collapse region that is formed at the junction of SMF and PCF specialty fiber, achieving very fast response and recovery times of 7 and 2 s, respectively, which can detect as low as 8.09 ppt (47.59 fM). The reusable probe showed the potential for rapid detection of ultra-trace ammonia with high selectivity and reproducible features. In 2020, Leal-Junior et al. [191] presented a low-cost fiber-optic probe for the early detection of ammonia. The sensor was based on the chemical interaction between the Oxazine 170 perchlorate layer, deposited in an uncladded polymer optical fiber, and the ammonia dissolved in water. In addition, a thin metallic layer (composed of gold and palladium) was deposited in the fiber end facet and acted as a reflector for the optical signals, enabling the use of the proposed sensor in reflection mode. Different configurations of the sensor were tested, where the effects of the PDMS protective layer, thermal treatments, and the use of reflection or transmission modes were compared in the assessment of ammonia concentrations in the range of 100 ppb to 900 ppb. Results showed better performance (as a function of the sensor sensitivity and linearity) of the sensor with the annealing thermal treatment and without the PDMS layer. Then, the proposed fiber-optic probe was applied to the ammonia detection in high-salinity water, and ammonia concentrations as low as 100 ppb were detected.

Another problematic factor is the usage of myriad pesticides and insecticides in agriculture that are dangerously applied to the soil, such as organophosphorus pesticides that are used worldwide for agricultural purposes. Miliutina et al. [192] reported a functional plasmonic sensor aimed at the monitoring of pesticide spreading and determination of their concentration. It utilized the functionalization of a plasmon-supported fiber surface decorated with a metal-organic framework compound with the formula Zn₄O(BDC)₃, also known as the MOF-5 layer. The MOF-5 layer provides the extraction of pesticides from the surrounding medium, which causes the shift in the plasmon resonance absorption band position. The created system demonstrated high selectivity and sensitivity towards organophosphorus pesticides. Particularly, the sensors were successfully applied for the detection of fenitrothion and paraoxon. The lower detectable concentration was found to be as small as 1 pM, which makes the proposed sensors comparable to common analytical approaches to pesticides detection. The proposed cost-effective technique allows for simple and straightforward pesticide detection, even in the complex samples, without any sample's pretreatment step, and it can be easily and scalable transferred to outdoor conditions. In 2019, Kant et al. [193] presented an SPR based fiber-optic sensor for the pesticide fenitrothion utilizing Ta_2O_5 nanostructures sequestered onto a rGO matrix. A thin layer of silver was deposited on the unclad core of silica optical fiber for SPR generation and was followed by the deposition of a sensing surface comprising a layer of tantalum oxide NPs sequestered in a nano-scaled matrix of rGO. The sensing mechanism is based on the interaction of fenitrothion with the silver film, which leads to a change in the RI. Characterized by a wavelength interrogation scheme, the fiber-optic sensor exhibited a redshift equaling 56 nm corresponding to fenitrothion concentration in the range from 0.25 to 4 μ M, including the blank solution. The spectral sensitivity is 24 nm/ μ M, the limit of detection is 38 nM, and the response time is as short as 23 s. The sensor is selective, repeatable, and works at ambient temperature.

6. Conclusions

This paper presents a review of optical fiber-based cost-effective biosensing platforms, covering different geometries, interrogation techniques, encapsulation methods, and, finally, presenting applications in three crucial fields of sensing for wellbeing: cancer and cardiovascular biomarkers detection and environmental monitoring.

The development and application of different types of optical fiber biosensing has increasing over the years at an impressive rate, with the demystification of some associated prejudices. In the past, optical fiber biosensing was labeled as high cost, and for this reason was only considered for highly specific applications. However, as fiber optic biosensor technologies have continued to show excellent performance, the instrumentation associated with their interrogation has dropped in cost, and essentially new methods of designing interrogating sensors have been developed. Fiber-optic sensors working in the visible range are naturally low-cost. Furthermore, the possibility to use the instrumentation of a smartphone to perform the readout of these types of sensors, in addition to considerably lowering the system cost, can be the basis for more democratic and accessible self-monitoring.

Fragility was also one of the points associated with fiber optic sensors. However, various encapsulation techniques are currently used to make these sensors highly robust while maintaining or even amplifying their sensitivity. Currently, we are able to obtain sensors with very high performance, which are highly robust and with practical and low-cost reading techniques.

The potential of all these cost-effective optical fiber sensing techniques reviewed in this article is supported by the numerous applications and excellent results shown in the three monitoring areas addressed. Nonetheless, as in any other biosensing technique, several aspects still need to be improved and worked on, namely, in the enhancement of their repeatability and reproducibility and by developing techniques to enable their reuse. In the future, it is expected that further research will be devoted to make progress on these topics.

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Abbreviations

The following abbreviations are used in this manuscript:

AMI Acute myocardial infarction APTMS (3-Aminopropyl)trimethoxysilane

ARS Alizarin Red S
B[a]P Benzo(a)pyrene
BSA Bovine serum albumin
CCDs Charge-coupled devices
CD44 Cluster of differentiation 44

CIP Ciprofloxacin CK7 Cytokeratin 7 CK17 Cytokeratin 17

CLIA Chemiluminescence assay
CVDs Cardiovascular diseases
DNA Deoxyribonucleic acid
DTC Differentiated thyroid cancer
EBL Electron beam lithography
ECG Electrocardiography

EDC N-ethyl-N'-(3-dimethylaminopropyl)carbodiimide

eFBG Etched Fiber Bragg grating

ELISA Enzyme-linked immunosorbent assay

ELP Electroless deposition

eTFBG Etched tilted Fiber Bragg grating FWHM Full width at half maximum

GNG Gold nanograting
GNS Gold nanostars
GNR Gold nanorods
GO Graphene oxide
GOx Glucose oxidase

HER2 Human epidermal growth factor receptor 2

HPTS1 Hydroxy-3,6,8-pyrene trisulfonic acid trisodium salt

HTR Human transferrin

IEMA Immunoenzymometric assay IRMA Immunoradiometric assay

ITO Indium tin oxide
LDF Light diffusing fibers
LEDs Light-emitting diodes
LOD Limit-of-detection

LRSPP Long-range surface plasmon–polariton LSPR Localized surface plasmon resonance MIPs Molecularly imprinted polymers

MMFs Multimode fibers

MZI Mach-Zehnder interferometer

NA Numerical aperture NHS *N*-hydroxysuccinimide

NPs Nanoparticles

OFB Optical fiber biosensor OFS Optical fiber sensor

OF-SPR Optical fiber-based surface plasmon resonance

OTA Ochratoxin A PANI Polyaniline

PBS Phosphate buffered saline

PCS Plastic-clad silica
PDs Photodetectors
PDMS Polydimethylsiloxane
PMMA Poly(methyl methacrylate)

PM-SMF Polarization-maintaining single-mode fiber

POC Point-of-care **POFs** Plastic optical fibers Refractive index RI RIU Refractive index units rGO Reduced graphene oxide **SMFs** Single-mode fibers **SNR** Signal-to-noise ratio **SPR** Surface plasmon resonance

Tn Troponin
Tg Thyroglobulin
TNT Trinitrotoluene

VEGF Vascular Endothelial Growth Factor

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