

Article



# **Effect of Ti Atoms on Néel Relaxation Mechanism at Magnetic Heating Performance of Iron Oxide Nanoparticles**

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**Abstract:** The study was based on understanding the relationship between titanium (Ti) doping amount and magnetic heating performance of magnetite (Fe<sub>3</sub>O<sub>4</sub>). Superparamagnetic nanosized Ti-doped magnetite ((Fe<sub>1-x</sub>,Ti<sub>x</sub>)<sub>3</sub>O<sub>4</sub>; x = 0.02, 0.03 and 0.05) particles were synthesized by sol-gel technique. In addition to (Fe<sub>1-x</sub>,Ti<sub>x</sub>)<sub>3</sub>O<sub>4</sub> nanoparticles, SiO<sub>2</sub> coated (Fe<sub>1-x</sub>,Ti<sub>x</sub>)<sub>3</sub>O<sub>4</sub> nanoparticles were produced as core-shell structures to understand the effects of silica coating on the magnetic properties of nanoparticles. Moreover, the magnetic properties were associated with the Néel relaxation mechanism due to the magnetic heating ability of single-domain state nanoparticles. In terms of results, it was observed that the induced RF magnetic field for SiO<sub>2</sub> coated (Fe<sub>0.97</sub>,Ti<sub>0.03</sub>)<sub>3</sub>O<sub>4</sub> nanoparticles caused an increase in temperature difference ( $\Delta$ T), which reached up to 22 °C in 10 min. The  $\Delta$ T values of SiO<sub>2</sub> coated (Fe<sub>0.97</sub>,Ti<sub>0.03</sub>)<sub>3</sub>O<sub>4</sub> nanoparticles were very close to the values of uncoated Fe<sub>3</sub>O<sub>4</sub> nanoparticles.

**Keywords:** oxide semiconductor; point defects; Néel relaxation; magnetic hyperthermia; superparamagnetic nanoparticles

# 1. Introduction

The crystal structure of magnetite (Fe<sub>3</sub>O<sub>4</sub>) has a spinel cubic structure with Fd-3m space group [1,2]. Fe<sub>3</sub>O<sub>4</sub> crystal structures obtain tetrahedral and octahedral sublattices, occupied by Fe<sup>2+</sup> and Fe<sup>3+</sup> cations coordinated with 32 oxygen atoms [1,2]. Ferrimagnetic properties are governed by the coupling of cation spins in octahedral and tetrahedral sites [1,2]. Properties, such as low toxicity, suitable magnetic properties, and easy fabrication, make the ferrite particles suitable for hyperthermia usages [3]. Each magnetic domain in a magnetic nanoparticle is oriented along the direction of the externally applied magnetic field. Also, especially for single-domain magnetic nanoparticles, all particles can be oriented in the direction of the magnetic field (Brownian motion), and the magnetic domains inside a particle can rotate without rotating the magnetic particle (Neel motion) [4]. In both cases, due to the desire to be in line with the external magnetic field, alternating magnetic fields initiate heating in and around single-domain magnetic particles.



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During the last few decades, magnetic hyperthermia has become a widely researched topic (especially for cancer treatments) [5-11]. In recent technological applications, ferrite nanoparticles have emerged as the most popular candidate for magnetic hyperthermia applications. The heating performance of magnetic nanoparticles is investigated mainly depending on size and magnetic domain number in the particle [12]. The heating performance of multi-domain magnetic nanoparticles is related to energy losses from magnetic hysteresis due to the domain wall motions and eddy currents induced in magnetic grains [12,13]. Unlike multidomain nanosized particles, the magneto-heating performance of single magnetic domain nanosized particles is mainly correlated with magnetic anisotropy, inter/intra particles interactions and homogeny distributed magnetic nanoparticles [14]. Magnetic relaxation is generally the dominant mechanism in magnetic anisotropy and thus in magneto heating performance of nanoparticles. The magneto-heating performance of single magnetic domain nanosized particles is associated with two mechanisms, Néel and Brownian relaxations [15–17]. Brownian motion is created by mechanical fluctuation that performing by an entire nanoparticle rotating its own axis. On the other hand, the Néel relaxation mechanism is independent from the rotational motion of particles. The internal flip of spins with respect to the crystalline lattice creates magnetic heating associated with Néel relaxation [18–20]. Both relaxation mechanisms cause an increase in the temperature of magnetic nanoparticles. Even though magneto-heating can be observed with temperature increase, the magneto-heating performance of magnetic nanoparticles can be measured by specific absorption rate (SAR) values. SAR values can be defined by absorbed/converted magnetic energy into thermal energy and in technological applications (especially in cancer treatments) mainly modified properties of magnetic nanoparticles [3]. For magnetic hyperthermia applications, Néel relaxation has certain advantages due to its high SAR values [3]. On the other hand, Brownian relaxation is largely dependent on the viscosity of the medium surrounded by particles, and in cancer treatment SAR value is not enough for magnetoheating processes [3]. Liquid environmental conditions, such as different viscosities of the medium, agglomeration of nanoparticles within different cells, or fixation of nanoparticles in cell membranes (or extracellular tissue), weaken or inhibit the mechanical rotation of magnetic nanoparticles [21]. During the cancer treatments, due to the lack of movement or fluctuation of magnetic nanoparticles inside cancer cells Néel relaxation became the dominant factor for magnetic heating in compare to Brownian movement [3,20,21].

The study is based on the evaluation of the magneto heating performance of individual, non-interacting, and monodisperse particles in highly viscous environments confirming the Néel relaxation mechanism. The magneto heating mechanism is investigated for superparamagnetic  $Fe_3O_4$  nanoparticles doped with Ti atoms. The magneto heating performance of magnetic nanoparticles is not sufficient for cancer treatments, and therefore the surface of magnetic nanoparticles must be functionalized by ligands [22,23]. The general way to improve the surface modification and functionalization ability of magnetic nanoparticles is to coat each nanoparticle with a  $SiO_2$  shell [23]. Optical transparency, highly biocompatibility, biodegradability, and manufacturability with porous surfaces make SiO<sub>2</sub> very useful as a biomaterial [24]. The study also includes understanding the effect of  $SiO_2$  on the heating performance of  $(Fe_{1-x}, Ti_x)_3O_4$  nanoparticles. Although coating with SiO<sub>2</sub> has no effect on crystal structures, the SiO<sub>2</sub> thickness has caused a remarkable decrease in internal magnetization value of  $Fe_3O_4$  [25]. One of the consequences of the reduction of internal magnetization can be observed as a decrease in the magneto heating performance of magnetic nanoparticles [26,27]. On the other hand, SiO<sub>2</sub> coating reduces interparticle interactions and inhibits aggregation, thereby increasing the heating capabilities of magnetic nanoparticles in AC fields [28–30]. Therefore, the magneto heating performance of magnetic nanoparticles is correlated with interacting or non-interacting particles [26-30].

The magnetic nanoparticles such as  $Fe_3O_4$  seem as efficient nano-heaters in biomedical applications [31–33]. Surface modified ( $Fe_{1-x}$ ,  $Ti_x$ )<sub>3</sub> $O_4$  nanoparticles are expected to be potential bio-materials which suitable for loading anticancer drugs and heating under both

excitations, RF magnetic field, and UV radiation, which utilize the magnetic nanoparticles useful at clinical hyperthermia applications.

#### 2. Experimental

Magnetite nanoparticles were prepared via the co-precipitation method. Ferrous chloride tetrahydrate (FeCl<sub>2</sub>·4H<sub>2</sub>O) and ferric chloride (FeCl<sub>3</sub>) were used as iron precursors containing different valance states. On the other hand, Tetraisopropil Ortotinatate (TIPO) was used as a Titanium source. Natrium hydroxide (NaOH: 25% by weight) in H<sub>2</sub>O and hydrochloric acid in water (HCl) were used as the precipitating agent. To prevent the agglomeration of magnetic nanoparticles, nanoparticles were coated with oleic acid at the end of the process. Ethanol and DI water were used to remove excessive coating agent. The schematic diagram of the procedure was demonstrated in Figure 1. The synthesis was performed on magnetic stirrer at 90 °C. Firstly, ferrous and ferric chloride iron salts were dissolved in water with HCl under Argon gas flow. Secondly, after half an hour, TIPO and NaOH were dropped into the solution. Oleic acid was added to the solution as the last step.



Figure 1. The schematic diagram of the synthesis procedure of pure and Ti-doped magnetite.

The coated nanoparticles were washed with DI water and ethanol to remove chloride ions and excessive coating materials. Furthermore, the same procedure was performed for the synthesis of the pure magnetite nanoparticles.

The nanoparticles were also coated with  $SiO_2$  by base-catalyzed silica formation from tetraethylorthosilicate (TEOS) in a water-in-oil microemulsion technique, which was mentioned in previous study [34]. The resulting mixture was vigorously stirred for 24 h.

The crystal structures of samples were investigated by x-ray powder diffractometer (XRD), employing a MiniFlex model XRD, produced by Rigaku Corporation (Tokyo, Japan). The XRD patterns were taken under Cu K<sub> $\alpha$ </sub> radiation (1.5406 Å) in 20 range from 10°

to 90°. A JEM-2010F high-resolution transmission electron microscope (HR-TEM) (JEOL Ltd., Tokyo, Japan) was used to picture the structural morphologies of nanoparticles. Dc magnetization ( $\sigma$  (H)) measurements were performed at 300 K temperatures in a magnetic field range of  $\pm$  2 T.

Magneto-thermal characterization was taken by a homemade setup constructed using the equipment with a frequency of 150 kHz (power generator, thermometer, etc.). Experiments were performed in a custom-made setup with an alcohol thermometer, a covered glass tube, a water-cooled magnetic coil (diameter 50 mm and four turns for 160 Oe), and an AC power generator (Istanbul, Turkey) with a constant frequency of 150 kHz. The SAR measurements were conducted in non-adiabatic conditions as in many publications. The colloidal solution was put into a glass tube and the tube was placed in the coil. The thermometer was directly inserted into the solution. The temperature was measured using the thermometer as a function of time for a duration of 15 min. The filling level of the solution in the tube was adapted to the half-length of the coil to minimize the effects of magnetic field inhomogenity. Between the glass tube and the coils, we used styrofoam as an insulating material.

## 3. Result and Discussions

The structural analyses were performed employing XRD patterns for SiO<sub>2</sub> coated and uncoated particles as shown in Figure 2a,b respectively. The patterns were in an agreement with Fe<sub>3</sub>O<sub>4</sub> diffraction pattern shown in ICDD card (PDF# 74-0748). No contamination or unexpected phase such as TiO<sub>2</sub> based structures, was detected on the XRD patterns. As seen in Figure 2b, even though having high background intensity, originating from the amorphous phase of SiO<sub>2</sub>, the Fe<sub>3</sub>O<sub>4</sub> patterns were distinctly distinguished at each XRD pattern.



Figure 2. The XRD patterns of (a) Ti-doped magnetite and (b) SiO<sub>2</sub> coated Ti-doped magnetite.

Not observing Ti elements or compounds in the xrd patterns indicated the displacement of Ti atoms inside of  $Fe_3O_4$  lattice. The ionic radius of  $Ti^{4+}$  is approximately 0.61Å, which is close to the ionic radius of  $Fe^{3+}$  (0.64 Å). Thus,  $Ti^{4+}$  ions are expected to settle instead of by  $Fe^{3+}$  ions in the octahedral lattice sites. Due to charge neutrality, a  $Ti^{4+}$  ion replacement in an octahedral site gives rise to change the valence state of  $Fe^{3+}$  ion to  $Fe^{2+}$  ion as shown in the chemical equation of (1) [35]. Due to its charge neutrality, the substitution of  $Ti^{4+}$ , a tetravalent positive ion, causes an  $Fe^{3+}$  ion to change its valency into  $Fe^{2+}$ . Chemical Equations (1) and (2) assign the charge neutrality occurring in  $(Fe_{1-x},Ti_x)_3O_4$  lattice.

$$Fe^{2+} \leftrightarrow Fe^{3+} + e^{-} \tag{1}$$

$$2Fe^{3+} \rightarrow Ti^{4+} + Fe^{2+}$$
 (2)

Increase in Ti<sup>4+</sup> substitution amount (0.01  $\leq x \leq$  0.025) in (Fe<sub>1-x</sub>,Ti<sub>x</sub>)<sub>3</sub>O<sub>4</sub> lattice cause to formation of Fe<sup>2+</sup><sub>A</sub>(Fe<sup>2+</sup>, Ti<sup>4+</sup>)<sub>B</sub>O<sub>4</sub> (A, tetrahedral side; B, octahedral side), which inhibits the hopping mechanism between iron ionic states. And thus, the new configuration cause to increase in magnetic anisotropy as mentioned in literature [35,36]. Excessive Ti<sup>4+</sup> substitution (x  $\geq$  0.025) easily bypasses Fe<sup>2+</sup> ions, resulting in the formation of differential vacancies [35].

After understanding the crystal structure and possible defects in a lattice, the particle size distributions were investigated employing TEM micrographs. TEM Figures assigned homogeny distributed nanoparticles. In addition, due to covering with oleic acid, no agglomeration across the entire particle distribution was realized as seen in Figure 3.



Figure 3. TEM micrographs of (a) uncoated and (b) coated (Fe<sub>0.97</sub>,Ti<sub>0.03</sub>)<sub>3</sub>O<sub>4</sub> nanoparticles.

The size frequencies of particles were calculated by the subprogram of Image-J2. Particles with high differences in particle size were selected in the calculations. As shown on Figure 4, the particles size distributions which were confirmed by ImageJ, were found as  $10.3 \pm 0.6$  nm and  $18.7 \pm 0.5$  nm for uncoated and coated (Fe<sub>0.97</sub>, Ti<sub>0.03</sub>)<sub>3</sub>O<sub>4</sub>, respectively. The particles' size distribution indicated that the particles were in superparamagnetic regions.

The magnetization ( $\sigma$  (H)) measurements at the room temperature were illustrated in the Figure 5 for both samples uncoated and coated (Fe<sub>1-x</sub>,Ti<sub>x</sub>)<sub>3</sub>O<sub>4</sub>. As seen from the Figure, the zero remanent magnetization assign the overcoming thermal energy to the magnetic anisotropy energy barrier at all samples. For both coated and uncoated (Fe<sub>1-x</sub>,Ti<sub>x</sub>)<sub>3</sub>O<sub>4</sub> particles, only difference in magnetization curves was the decrease in magnetization value by Ti amount in lattice. In order to understand the magnetic domain states of particles, room temperature magnetic hysteresis curves were obtained as shown in Figure 5.



**Figure 4.** Calculated size distributions of (**a**) uncoated and (**b**) coated (Fe<sub>0.97</sub>, Ti<sub>0.03</sub>)<sub>3</sub>O<sub>4</sub> nanoparticles via ImageJ subprogram.



Figure 5. The magnetization measurements of (a) uncoated and (b)  $SiO_2$  coated pure and Tidoped magnetite.

Not observing coercivity and remanence values on hysteresis curves at room temperature proved that the particles were in superparamagnetic states [37,38]. Heating performance of an individual, non-interacting, and monodisperse particles have high SAR value due to dominating of Néel relaxation in a highly viscous environment. In the superparamagnetic state at room temperature, the magnetic interaction decreased to the lowest values with the effect of thermal energy. Nanoparticle size, anisotropy and interparticle interaction are the principal factors that influence heat generation [39,40]. Moreover, agglomerations between the magnetic nanoparticles occur due to strong magnetic dipoledipole interactions between particles. As seen in Figure 5, coating with SiO<sub>2</sub> and doping with Ti atoms reduced the magnetization of the particles. A possible mechanism reducing the magnetization of particles is the number of vacancies in the crystal structure. The vacancies may form non-magnetic regions on the particle surface [41]. In addition, the magnetic anisotropy of the easy-axis is another parameter affecting the magnetization values that can be different either parallel or perpendicular to the easy-axis orientation of the domains [42]. These parameters manage the SAR values of particles.

Being predominant of thermal energy to magnetic energy at room temperature indicated that the nanoparticles were in superparamagnetic region. Since the particles were in superparamagnetic state size, the Néel magnetic relaxation was expected to dominate the magneto heating performance of the particles.

In Figure 6, the heating performance of particles was investigated under an ac magnetic field, approximately 13 kA/m field intensity and frequency of 150 kHz (the biological limits are  $5 \times 10^9$  A/(m.s). The magneto heating measurements were taken immediately after arranging nanoparticles as magnetic fluids in 1 ml ethanol media. As seen on the Figure, the temperature difference ( $\Delta$ T) reach up to 30 °C for pure Fe<sub>3</sub>O<sub>4</sub> and for the same time interval Ti doping lowered the  $\Delta$ T value down to 20 °C ((Fe<sub>0.97</sub>,Ti<sub>0.03</sub>)<sub>3</sub>O<sub>4</sub> nanoparticles). For (Fe<sub>0.97</sub>,Ti<sub>0.03</sub>)<sub>3</sub>O<sub>4</sub> nanoparticles, the  $\Delta$ T value was measured as 22 °C, which was good enough for use in vivo studies.



**Figure 6.** The magnetic heating performance of (**a**) uncoated and (**b**)  $SiO_2$  coated pure and Ti doped Fe<sub>3</sub>O<sub>4</sub>.

Then, for each particle, the SAR value was calculated as shown in Table 1. The physical quantity of SAR value was determined by defined as the heat released from colloidal magnetic nanoparticles in unit time by Equation (3) [43].

$$SAR = \frac{Q}{\Delta t m_{mag}}$$
(3)

where  $Q = mc\Delta T$ ,  $m_{mag}$  is the mass of magnetic nanoparticle, c is the specific heat of the colloid (only ethanol is taken into account, the contribution of magnetic nanoparticles, oleic acid, and SiO<sub>2</sub> to the specific heat are neglected). The calculations were performed for the heat capacity and the density of ethanol 2.57 kJ/(kgK), 0.789 g/mL, respectively.

**Table 1.** SAR values of  $(Fe_{1-x}, Ti_x)_3O_4$  (x = 0.00, 0.02 and 0.03) nanoparticles.

Nanoparticles	SAR (W/g)
Fe <sub>3</sub> O <sub>4</sub>	155
x = 0.02	70
x = 0.03	3
SiO <sub>2</sub> coated Fe <sub>3</sub> O <sub>4</sub>	104
x = 0.02 (SiO <sub>2</sub> coated)	34
x = 0.03 (SiO <sub>2</sub> coated)	116

The SAR values were illustrated in Table 1. As understood from Table 1, coating with  $SiO_2$  lowered the SAR values of  $(Fe_{1-x},Ti_x)_3O_4$  nanoparticles. However, for  $SiO_2$  coated samples increase in  $Ti^{4+}$  ions amount in lattice caused an increase in SAR value, which getting closer to the value of pure  $Fe_3O_4$  nanoparticles.

## 4. Conclusions

In the study, homogeny size distributed  $(Fe_{1-x},Ti_x)_3O_4$  ferrite nanoparticles in oleic acid and at SiO<sub>2</sub> matrix were synthesized via a chemical route. The particles were obtained as superparamagnetic Fe<sub>3</sub>O<sub>4</sub> nanoparticles to dominate the Néel relaxation over Brownian relaxation mechanism. Furthermore, lowering the particle size down to superparamagnetic region, coating with SiO<sub>2</sub> and Ti doping into the lattice was the tuned parameters to produce individual, non-interacting, and monodisperse particles. Then, the heating mechanism of SiO<sub>2</sub> coated Ti doped Fe<sub>3</sub>O<sub>4</sub> nanoparticles were only correlated with Ti atoms amount in the lattice. Due to the expected coupling between Ti<sup>4+</sup>-Fe<sup>2+</sup> ions in the octahedral site, the heating performance by Ti doping was lower than pure Fe<sub>3</sub>O<sub>4</sub>. On the other hand, for SiO<sub>2</sub> coated (Fe<sub>0.97</sub>,Ti<sub>0.03</sub>)<sub>3</sub>O<sub>4</sub> nanoparticles, the increase in the amount of Ti<sup>4+</sup> ions in lattice cause an increase in SAR value ( $\Delta T = 22$  °C in 10 min), while decreasing for uncoated nanoparticles. The heating performance of (Fe<sub>0.97</sub>,Ti<sub>0.03</sub>)<sub>3</sub>O<sub>4</sub> nanoparticles coated with SiO<sub>2</sub> was almost close to the heating performance of pure magnetite.

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