



Review

Majority and Minority Charge Carrier Traps in n-Type 4H-SiC Studied by Junction Spectroscopy Techniques

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Abstract: In this review, we provide an overview of the most common majority and minority charge carrier traps in n-type 4H-SiC materials. We focus on the results obtained by different applications of junction spectroscopy techniques. The basic principles behind the most common junction spectroscopy techniques are given. These techniques, namely, deep-level transient spectroscopy (DLTS), Laplace DLTS (L-DLTS), and minority carrier transient spectroscopy (MCTS), have led to recent progress in identifying and better understanding the charge carrier traps in n-type 4H-SiC materials.

Keywords: 4H-SiC; Schottky barrier diodes; defects; DLTS

1. Introduction

Today, 4H-SiC materials are regarded as one of the most promising types of materials for electronic devices. It has the largest bandgap among all SiC polytypes, and due to the high and isotropic mobility of charge carriers, it is preferred as a material for power electronics [1], bipolar devices [2], and quantum sensing [3]. Moreover, it has found applications in radiation detection too [1,4,5].

Even wider applicability of 4H-SiC has mostly been hindered by electrically active deep level defects present in the material. Electrically active defects introduce energy levels into the bandgap that act as traps for charge carriers. The majority charge carrier traps in n-type 4H-SiC material have been systematically investigated for many years. On the other hand, minority charge carrier traps in n-type 4H-SiC material have been much less investigated [6–9].

Junction spectroscopy technique is a term describing measurements performed on a semiconductor junction using electrical or electro-optical techniques [10]. The role of the junction is to create a depletion region, as its usage brings an important advantage over other bulk techniques. The advantage is that it is much easier to manipulate the occupancy of defects producing energy levels in the bandgap within the depletion region than in bulk region [10]. The most common junction spectroscopy technique is deep-level transient spectroscopy (DLTS). DLTS can detect defects in concentrations around 10^{10} cm^{-3} [10]. It provides information regarding the activation energy for electron and hole emission, capture cross-section, and concentration of defects. However, the main problem associated with DLTS is the lack of energy resolution (i.e., it is almost impossible to resolve two closely spaced deep energy levels). An improvement came in another junction spectroscopy technique called Laplace DLTS (L-DLTS), which brought an order of magnitude better energy resolution (meV) [11].

While DLTS is mostly used for studying the electrically active defects associated with majority charge carrier traps, minority charge carrier traps have been much less investigated. The basic principles of minority transient spectroscopy techniques were described by Hamilton et al. [12], and later by Brunwin et al. [13].

In this review, we provide an overview of the main majority and minority charge carrier traps in n-type 4H-SiC materials. We focus this overview on the recent advances



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achieved by means of junction spectroscopy techniques. Moreover, the basic principles behind the junction spectroscopy techniques responsible for these advances are provided. The main aim of this mini-review is to provide a basic description of junction spectroscopy techniques supported by examples from different studies on n-type 4H-SiC materials for graduate students and early-stage researchers.

2. Materials and Methods

2.1. Schottky Barrier Diode

As mentioned above, we need a junction to apply junction spectroscopy techniques for studying the electrically active defects in semiconductors. One of the basic junctions is the Schottky barrier diode (SBD). Figure 1 shows a schematic cross-section of a typical n-type 4H-SiC SBD.

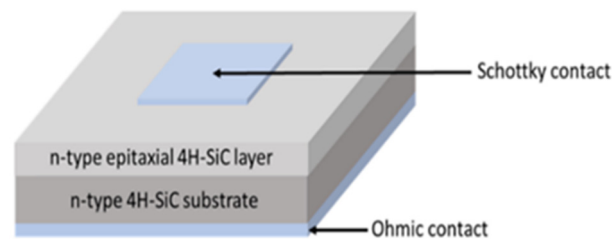


Figure 1. A schematic cross-section of a typical n-type 4H-SiC Schottky barrier diode (SBD).

SBD is formed when a metal and an n-type semiconductor are joined. Due to band alignment, a depletion space charge region (a volume depleted of charge carriers) is created at the semiconductor side of the junction [5]. As we apply the external bias, we enable the transport of electrons, and a net current will flow through the SBD. The SBD is reversely biased when a negative voltage (V_R) is applied to the metal with respect to the semiconductor. The SBD is kept under reverse bias when junction transient spectroscopies are applied.

Typical n-type 4H-SiC SBDs are produced on nitrogen-doped 4H-SiC epitaxial layers. The epitaxial layers are grown on thick silicon carbide substrate. The Schottky barriers are formed by thermal evaporation of nickel through a metal mask with openings, while ohmic contacts are formed on the backside of the silicon carbide substrate by nickel sintering.

2.2. Measurement Techniques

Electrically active defects introduce energy levels into the bandgap of a semiconductor, which act as traps for charge carriers. We can distinguish between majority and minority charge carrier traps. Majority carrier traps in semiconductors have been successfully and extensively studied by DLTS for decades [10,11], while minority charge carrier traps have been much less studied. In principle, it is possible to investigate these traps using DLTS by applying forward bias [14], but more reliable results can be obtained by using minority carrier transient spectroscopy (MCTS). The main principles of DLTS and MCTS will be described next. More detailed information about these techniques is given elsewhere [10,11].

2.2.1. Deep-Level Transient Spectroscopy (DLTS)

Basic DLTS measurement consists of the repetitive filling and emptying of deep levels (E_T) in the depletion region of the SBD by a bias pulse, as shown in Figure 2. The n-type SBD is operated under reverse bias V_R (Figure 2a), which is reduced to V_P during a bias pulse (Figure 2b). Empty traps, residing in the former depletion region, will be able to capture free carriers and become occupied (Figure 2b). After restoring the original bias V_R , the charge in the depletion region will be lower than before, due to the trapped charge carriers (Figure 2c). These carriers will be released again through thermal emission, which proceeds exponentially in time. This thermal discharging of the occupied traps is monitored

by measuring the capacitance of the reverse biased diode as a function of time after the filling pulse (Figure 2d).

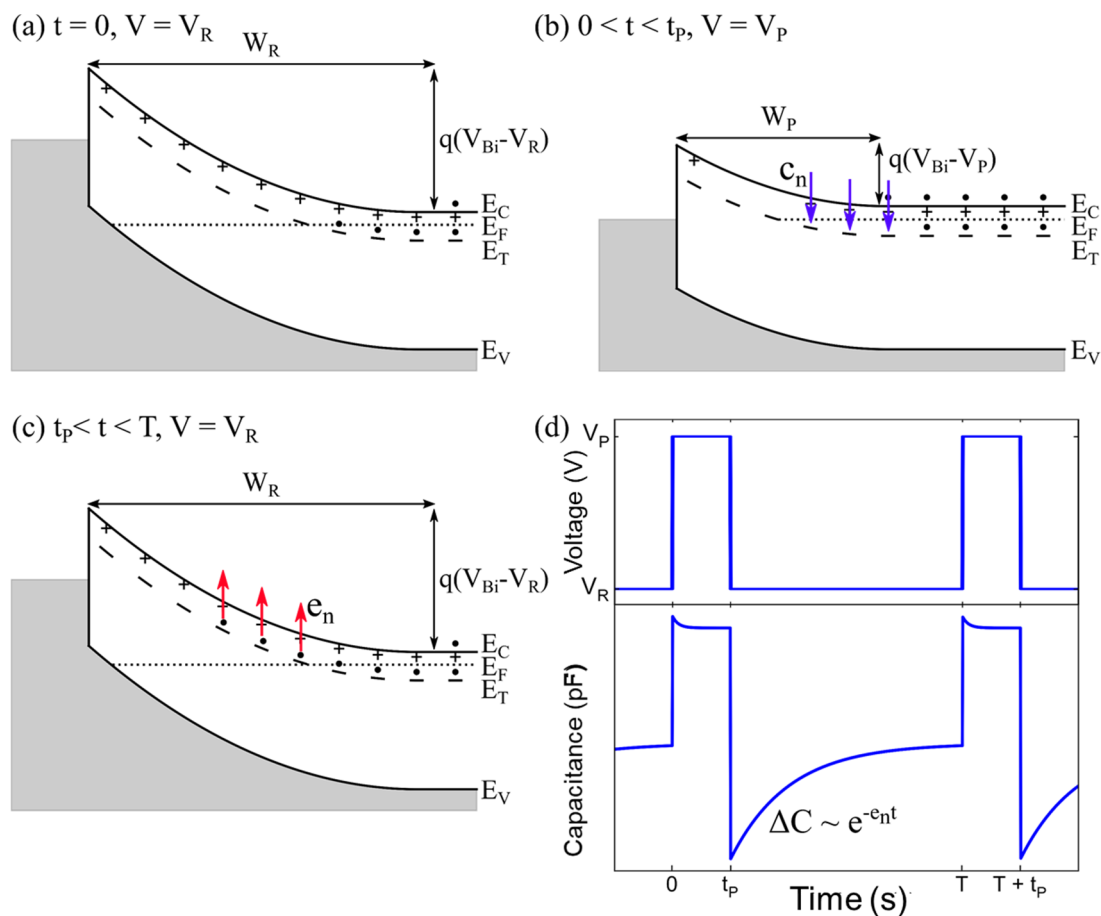


Figure 2. The basic principle of deep-level transient spectroscopy (DLTS) measurement: (a) equilibrium state with the applied reverse bias; (b) capture of the majority charge carriers in an n-type material while the pulse is applied V_P ; (c) emission of the trapped charge carriers; (d) the measured capacitance transient as a function of time. Here, t_p is the filling pulse duration, W_R and W_P are the depletion region widths for the applied V_R and V_P , respectively, e_n is the electron emission rate, and c_n is the electron capture rate. Figure adapted from Ref. [15].

The transient analysis and the problems with the signal–noise ratio not only for DLTS, but for MCTS and L-DLTS as well, will not be addressed here, as that goes beyond this mini-review paper. This information can be found elsewhere [10–13].

2.2.2. Minority Carrier Transient Spectroscopy (MCTS)

The basic principle of MCTS measurement is slightly different, as minority charge carriers are optically generated by use of above-bandgap light [10]. The MCTS measurement consists of the repetitive filling and emptying of deep levels (E_T) by optical pulses with an energy ($h\nu$) just above the bandgap energy (E_g), as shown in Figure 3. The minority charge carriers are captured during the optical pulse (Figure 3a), followed by the emission after the optical pulse (Figure 3b). Optical excitation can be applied through a semi-transparent Schottky contact, or from the back. If the sample thickness is greater than the minority carrier diffusion length, then the sample should be thinned [10].

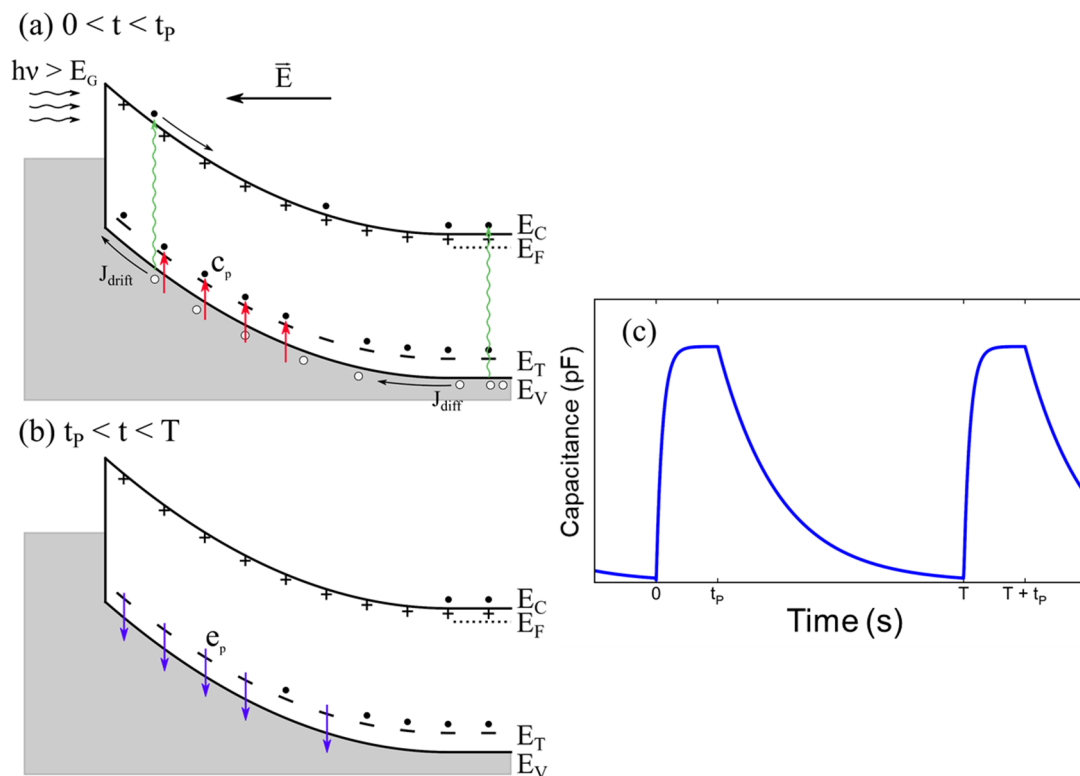


Figure 3. The basic principle of minority carrier transient spectroscopy (MCTS) measurement: (a) capture of the minority charge carriers from the valence band to the deep level (E_T) while optical pulses are applied; (b) emission of the minority charge carriers from the deep level after the optical pulse; (c) measured capacitance transient as a function of time. Figure adapted from Ref. [15]. Here, $h\nu$ is the energy of the optical excitation, J_{drift} and J_{diff} are drift and diffusion current densities, respectively, t_p is the optical pulse duration, e_p is the hole emission rate, and c_p is the hole capture rate.

3. Electrically Active Defects in n-Type 4H-SiC

3.1. Majority Charge Carrier Traps

Figure 4 shows a typical DLTS spectrum for the as-grown n-type 4H-SiC SBD. One peak, labelled as $Z_{1/2}$, is observed. The estimated activation energy for electron emission is $E_C - 0.67$ eV. It should be noted that since 4H-SiC is a wide bandgap material, DLTS measurements are usually performed in the wider temperature range, 80–700 K.

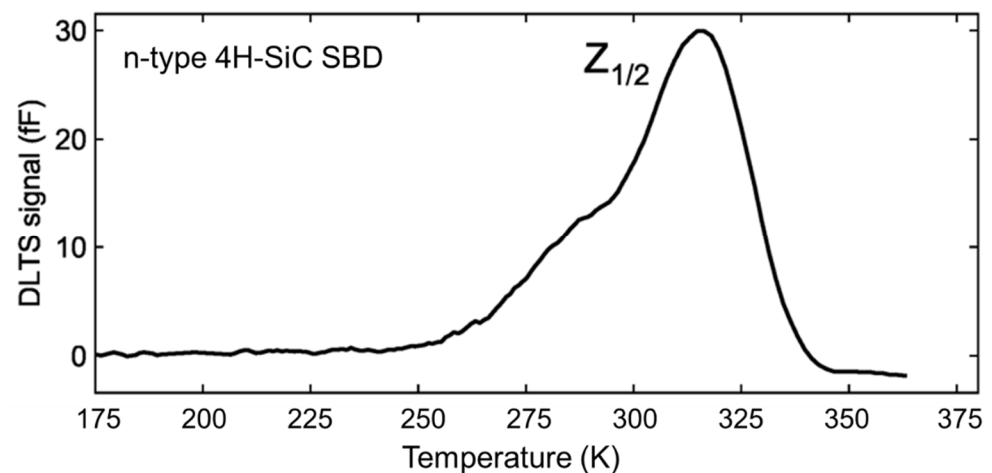


Figure 4. Typical DLTS spectrum for the as-grown n-type 4H-SiC SBD. Data adapted from Ref. [16].

Son et al. [17] were the first to ascribe $Z_{1/2}$ to $(= / 0)$ transition from the carbon vacancy (V_C). Carbon vacancy (V_C) is the most studied defect in n-type 4H-SiC. It is introduced during the crystal growth and upon irradiation [18].

As we can see in Figure 4, the DLTS peak is rather asymmetric. Hemmingsson et al. [19] showed that $Z_{1/2}$ is the superposition of two almost identical Z_1 and Z_2 transitions, which cannot be resolved by DLTS. This problem (too closely spaced deep energy levels) is an ideal case for the application of another junction spectroscopy technique, L-DLTS. L-DLTS is an isothermal technique in which the capacitance transient (measurements follow the same principle as described in Figure 2) is averaged at a fixed temperature. It provides a spectral plot of a processed capacitance signal against emission rate rather than against temperature. Moreover, L-DLTS provides an energy resolution in meV. More details on L-DLTS can be found elsewhere [11].

Recently, using L-DLTS measurements, direct evidence that $Z_{1/2}$ consists of two components has been observed, namely Z_1 and Z_2 , with activation energies for electron emission of $E_c - 0.59$ and $E_c - 0.67$ eV, respectively. These were ascribed to $(= / 0)$ transitions of V_C located at hexagonal (-h) and pseudo-cubic (-k) sites of the 4H-SiC crystal, respectively [20,21]. Figure 5 shows an L-DLTS spectrum for the as-grown 4H-SiC SBD. Two emission lines (Z_1 and Z_2) arising from the $Z_{1/2}$ DLTS peak are clearly resolved.

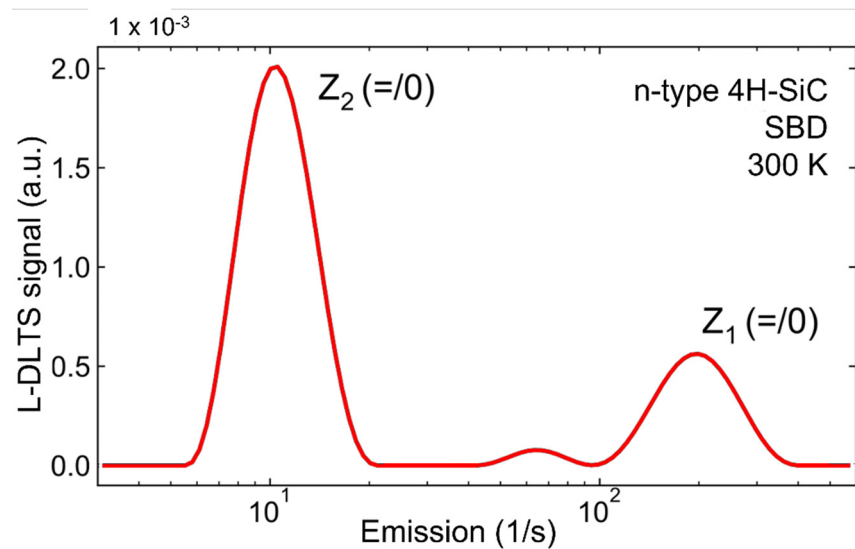


Figure 5. Laplace DLTS (L-DLTS) spectrum for the as grown n-type 4H-SiC SBD measured at room temperature (RT) [15].

Figure 6 shows the DLTS spectrum for the 2 MeV He ion-implanted n-type 4H-SiC SBD. In addition to the $Z_{1/2}$, 2 MeV He ion implantation has introduced two deep level defects, namely S_1 and S_2 . Estimated activation energies for electron emissions for S_1 and S_2 are $E_c - 0.40$, and $E_c - 0.70$, respectively.

Using the L-DLTS measurements, Bathen et al. [3] have shown that S_1 (in a proton-irradiated n-type 4H-SiC sample) has two emission lines arising from V_{Si} sitting at -k and -h lattice sites. This result was later confirmed by Capan et al. [22] when studying a neutron-irradiated n-type 4H-SiC SBD.

Deep level defect energies $E_c - 0.40$ eV and $E_c - 0.70$ eV have been reported in numerous studies for the proton-irradiated [3], ion-implanted [18,23], and neutron-irradiated [22,24] n-type 4H-SiC samples. All results indicated that these defects are intrinsic defects, either interstitials or vacancies. Recent progress in understanding the S_1 and S_2 defects has been achieved by Bathen et al. [3]. They have identified the S_1 and S_2 as $V_{Si} (-3 / =)$ and $V_{Si} (= / -)$ charge states.

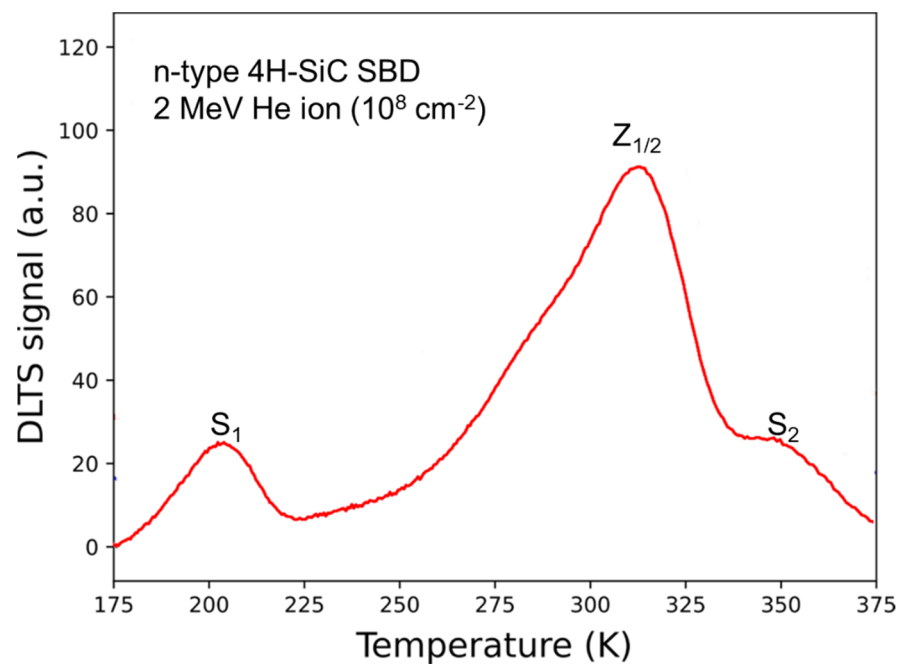


Figure 6. DLTS spectrum for 2 MeV He ion-implanted n-type 4H-SiC SBD. Several deep level defects are observed: S_1 , $Z_{1/2}$, and S_2 . Data adapted from Ref. [16].

Moreover, two deep levels with identical activation energies for electron emission of $E_c - 0.40$ and $E_c - 0.70$ eV have been observed in the low-energy electron irradiated n-type 4H-SiC [25,26]. These defects (not shown here) are labeled as EH_1 and EH_3 and have been identified as carbon interstitial-related (Ci) defects [26].

In addition to the above-mentioned majority charge carrier traps ($Z_{1/2}$, $S_{1/2}$, $EH_{1/3}$), we should not forget other majority carrier traps, which are usually present in the as-grown n-type 4H-SiC materials or are introduced by radiation or additional annealing. The most common traps are $EH_{4/5}$ and $EH_{6/7}$. The $EH_{4/5}$ introduce peaks in the DLTS spectrum at around 450–500 K, while the $EH_{6/7}$ introduce a broad DLTS peak at around 650 K (not shown here). They are assigned to carbon antisite-carbon vacancy (CAV) complex [27,28] and $(0/++)$ transition of the Vc [17], respectively. The evidence that $EH_{6/7}$ consists of two components was provided by Alfieri and Kimoto [29]. In their work, they resolved two energy levels at $E_c - 1.30$ and $E_c - 1.49$ eV, for EH_6 and EH_7 , respectively, using L-DLTS measurements.

3.2. Minority Charge Carrier Traps

In the following text, we will focus on the minority charge carrier traps in an n-type 4H-SiC. Most published works are aluminum or boron related. In this paper, we will provide examples which are boron related. Boron is introduced in SiC intentionally for p-type doping, or unintentionally during crystal growth [2,7].

Figure 7 shows a MCTS spectrum for the n-type 4H-SiC SBD. Two MCTS peaks labeled as B and D-center (multiplied $10\times$, for clarity) with activation energies for hole emissions of $E_v + 0.28$ and $E_v + 0.54$ eV are detected.

Two boron-related defects, labelled as B (shallow boron) and D-center (deep boron state) have been extensively studied. They are assigned to substitutional boron atoms occupying the Si and C-site, respectively [30–32].

Similar to $Z_{1/2}$, $S_{1/2}$, and $EH_{6/7}$, further improvements in studying the D-center by L-MCTS measurements have been made. Recent results have shown that the D-center has two emission lines, D1 and D2 (Figure 8), with activation energies for hole emissions of $E_v + 0.49$ eV and $E_v + 0.57$ eV, respectively. These emission lines are assigned to an isolated boron sitting at the C site, -h and -k site, respectively [31].

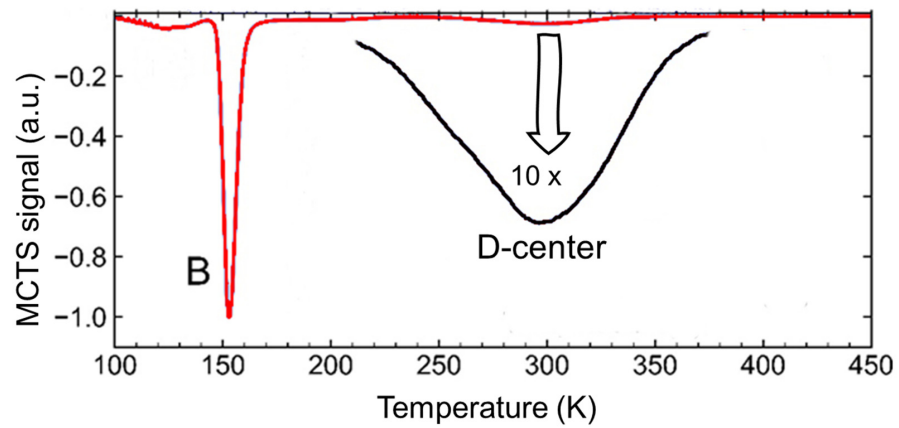


Figure 7. A normalized MCTS spectrum for the n-type 4H-SiC SBD. Data adapted from Ref. [32].

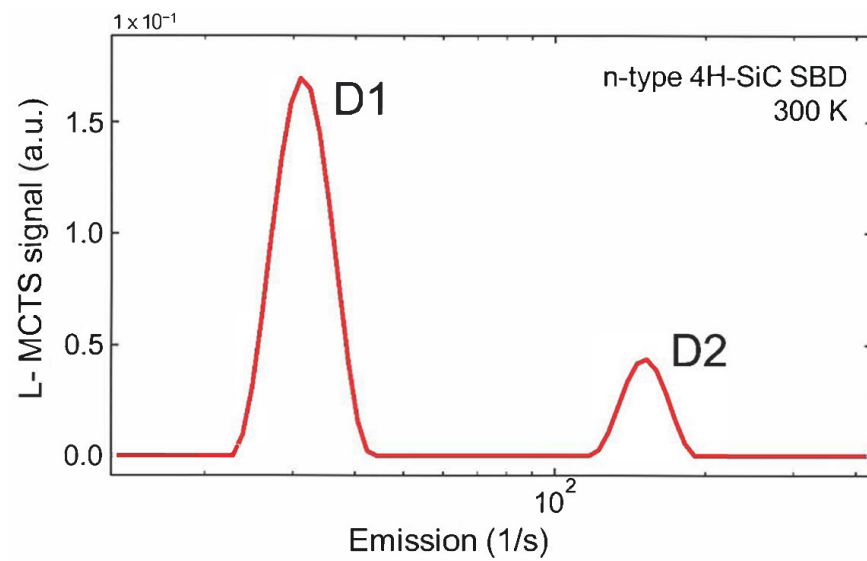


Figure 8. L-MCTS spectrum for D-center in then-type 4H-SiC SBD measured at 300 K. Data adapted from Ref. [15].

At the end of this section, we summarize the relevant information of the above-discussed majority and minority charge carrier traps. Figure 9 shows a graphical representation of the energy levels based on DLTS measurements. More information, including the energy levels observed by L-DLTS and MCTS measurements, is given in Table 1.

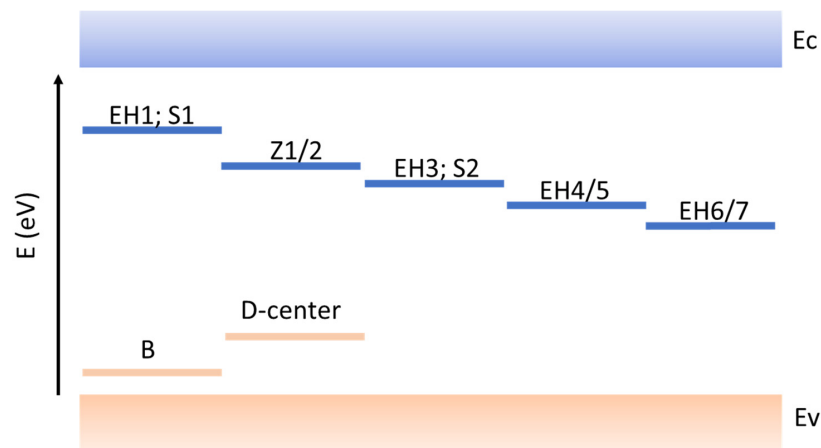


Figure 9. Graphical representation of the energy levels, based on DLTS measurements.

Table 1. Details on the majority and minority charge carrier traps in n-type 4H-SiC materials.

Trap Label	Identification	Activation Energy (eV)	References
EH ₁	C _i	Ec – 0.40	[25,26]
EH ₃	C _i	Ec – 0.70	[25,26]
S ₁	V _{Si} (–3/–)	Ec – 0.40	[3,22]
S ₂	V _{Si} (=/-)	Ec – 0.70	[3,22]
Z ₁	V _c (=0)	Ec – 0.59	[17,19–21]
Z ₂	V _c (=0)	Ec – 0.67	[17,19–21]
EH _{4/5}	C _{Si} -V _c (+/0)	Ec – 1.10	[27,28]
EH ₆	V _c (0/++)	Ec – 1.30	[17,18,29]
EH ₇	V _c (0/++)	Ec – 1.40	[17,18,29]
B	B _{Si}	Ev + 0.28	[2,7,30–32]
D-center	B _C	Ev + 0.54	[2,7,30–32]

4. Conclusions

In this review paper, the basic principles behind the most used junction spectroscopy techniques, DLTS and MCTS, were given. We have provided examples from different studies on majority and minority charge carrier traps. Successful applications of junction spectroscopy techniques, which have led to better understanding of the charge carrier traps in n-type 4H-SiC materials, were highlighted.

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