



# Article Molybdenum Oxide Thin Films Grown on Flexible ITO-Coated PET Substrates

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**Abstract:** Molybdenum oxide thin films were deposited on stiff and flexible substrates by reactive DC magnetron sputtering. Two sets of samples were prepared. The first with different  $O_2/Ar$  flow rate ratios and the second, fixing the oxygen content, with different time of deposition. As the  $O_2/Ar$  flow rate ratio varies from 0 up to 0.56, a threshold was found, ranging from crystalline to amorphous nature, and from a nontransparent appearance with metallic-like electrical conductivity to transparent and dielectric behaviour. From the second set, all transparent, the MoO<sub>x</sub> films present a compact/dense and featureless morphology with thickness from 190 up to 910 nm, depending on the time of deposition. Their structure was corroborated by XPS and Rutherford Backscattering Spectrometry (RBS) and density measurements were performed by RBS and X-ray reflectivity (XRR), revealing a value of 2.4 g/cm<sup>3</sup>. The surface roughness is in the order of a few nanometers and the maxima optical transmission, in the visible range, is around 89%. Electrochemical cyclic voltammograms showed noticeable color reversibility and reproducibility on the flexible substrates opening new framework possibilities for new electrochomic devices.

**Keywords:** MoO<sub>x</sub>; flexible substrates; electrochromic; DC magnetron sputtering

## 1. Introduction

Thin films based on oxide materials have gained momentum in recent years due to their properties that include high optical transparency, good electrical carrier mobility, and mechanical stress tolerance. The possibility of tailoring their physical properties, resulting in different electronic and geometric structures that can exhibit metallic, semiconducting, or insulating characters, makes this class of materials unique. This increases their importance as potential materials in a wide range of applications that include transparent electronics, optoelectronics, magneto electronics, photonics, spintronics, thermoelectrics, piezoelectrics, power harvesting, hydrogen storage and environmental waste management [1,2].

Beyond all amazing properties related to oxide materials, there is a well-known group of metal oxides (from Ti, V, Cr, Mn, Fe, Co, Ni, Nb, Mo, Rh, Ta, W and Ir) that can exhibit electrochromic behaviour, reversely changing their optical characteristics in a persistent mode under the action of an electrical field. The optical changes quantified by transmittance



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**Copyright:** © 2021 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). and/or reflectance are associated with an electrochemically induced oxidation-reduction reaction [3–6].

The electrochromic behaviour of molybdenum oxide (MoO<sub>3</sub>) was discovered around 1980 [7]. This oxide, in parallel with tungsten oxide (WO<sub>3</sub>), is one of the most studied electrochromic materials [3,5]. Several techniques of deposition have been used to grow MoO<sub>3</sub> films such as sputtering [8–10], thermal evaporation [11], sol-gel [12,13], electron beam evaporation [13,14], chemical vapour deposition [15,16], pulsed laser deposition [17,18], spray pyrolysis [19,20] and atomic layer deposition [21,22]. However, the electrochromic performance of a device based on these films is still a big scientific and technological challenge. The electrochromic performance depends on multiple factors, for example: deposition technology used to produce thin films with controlled porosity; type of microstructure; stoichiometry; combination of high optical transparency with high electrical conductivity of the transparent conducting oxide (TCO) used; adhesion of the thin film layers in the sandwich approach; high chemical stability; management of the process of charge insertion/extraction and charge balancing; and electrolytes quality when in contact with the electrolyte environment [23–26].

The improvement of the electrochromic performance and properties is being investigated by mixing some of the simple metal oxides, as  $MoO_3$ ,  $WO_3$  or  $V_2O_5$ , at the nanometric scale, forming new materials with coloring tuning, low response time, higher contrast ratio, etc. [4,27–30]. It is known that many cycles between coloring and bleaching tend to degrade the electrochromic performance and some efforts to study the rejuvenation of  $MoO_3$  thin films are being done [7,31]. Despite the electrochromic characteristic,  $MoO_3$ films are also interesting materials for energy applications [3].

In this work,  $MoO_x$  thin films were grown by reactive DC magnetron sputtering, in function of  $O_2/Ar$  flow rate ratio and deposition time, onto different substrates that include polyethylene terephthalate (PET) and glass, both with indium tin oxide (ITO) conductive layer, and glass plates without ITO and silicon. The structural, chemical, optical, electrical, and electrochemical properties were studied.

#### 2. Materials and Methods

#### 2.1. Thin Films Preparation

 $MoO_x$  thin films were grown by reactive DC magnetron sputtering, using a 60 L custom-made laboratory sized deposition system [32].  $MoO_x$  thin films were obtained using a  $200 \times 100 \times 6 \text{ mm}^3$  metallic Mo target with 99.95% of purity, sputtered by applying a current density of 100 A/m<sup>2</sup> in a plasma composed of argon (Ar, working gas), and a reactive gas that was oxygen (O<sub>2</sub>). The target was placed at 7 cm from the rotating substrate holder. Before the film deposition, the sputtering system was evacuated to  $1 \times 10^{-3}$  Pa. Argon and oxygen, both of 99.997% of purity, were introduced through mass flow-controlled gas inlets.

For the first series of samples, different  $O_2/Ar$  flow rate ratios were used to prepare a set of films with different Mo oxidation states. The flow of Ar was kept constant in all depositions (25 sccm, corresponding to a partial pressure of  $3.9 \times 10^{-1}$  Pa) and only the  $O_2$ flow was changed. The  $O_2/Ar$  flow ratios used were 0, 0.16, 0.32, 0.48 and 0.56. The time of deposition was kept for 60 min in all of these samples. In the second series, the  $O_2/Ar$ gas flow ratio was fixed at 0.56 and the deposition time was decreased from 60 min (used in the first series) to 30 min and to 15 min, to prepare  $MoO_x$  with different thicknesses. No external heating was used during the deposition and the substrates were grounded.

All substrates were previously cleaned by plasma treatment using a Zepto plasma system (Diener Electronic, Ebhausen, Germany), first with oxygen for 5 min and then with argon for 15 min. It must be noted that the polyethylene terephthalate (PET) and glass substrates, both with an indium tin oxide (ITO) conductive layer, present a sheet resistance of 5  $\Omega/\Box$  and 60  $\Omega/\Box$ , respectively. The p-type silicon substrate (Boron doped) with (100) orientation, shows a resistivity from 5–10  $\Omega \times cm$ .

#### 2.2. Characterization

#### 2.2.1. Microstructural and Morphologic Characterization

The as-deposited films were undergone for characterization. Their structure was analyzed by an X-ray diffractometer (PANalytical XPert-Pro, Almelo, The Netherlands) with CuK $\alpha$  radiation ( $\lambda = 1.54060$  nm), operating at a scan rate of  $0.01^{\circ}$ /s over a 20 range of 10-80°. The crystallite size was calculated using the Scherrer equation. X-ray reflectivity (XRR) measurements were performed on a Bruker D8 AXS diffractometer equipped with a Cu source and a soller slit in the secondary beam path to decrease the horizontal divergence. A 0.2 mm width slit and a Ni filter were placed at the primary beam side to reduce the horizontal beam divergence and to eliminate the K-beta line, respectively. Atomic compositions and the depth concentration profile (in the growth direction) were measured by Rutherford Backscattering Spectrometry (RBS), with a 2 MeV  $\alpha$  particle beam of 1 mm diameter obtained from a Van de Graaff accelerator. The random spectra were obtained by tilting the sample by 5 degrees and rotating the sample during the measurement to avoid channeling in the substrate. The backscattered particles were detected using two PIN diode detectors mounted at different backscattering angles of  $165^{\circ}$ and  $140^{\circ}$ . The compositions were extracted as a function of depth by a fitting procedure, using the nuclear data furnace code (NDF) [33]. XPS spectra were performed with a Kratos Axis Ultra HAS spectrometer using the Al K $\alpha$  excitation line at 1487 eV. XPS signals were calibrated using the C1s peak at 285 eV resulting from the accidental hydrocarbons layer present on the sample surface. For these measurements, an energy step size of 0.1 eV and and a pass energy of 40 eV were used, while the analyzed area was 300  $\mu$ m  $\times$  700  $\mu$ m. Raman spectra were obtained with a Horiba Jobin Yvon HR800 spectrometer using the 532 nm line, of a 50 mW He-Cd laser. Surface, cross-section morphology and thin-film thickness estimation were carried out by scanning electron microscopy (SEM) using a Hitachi SU-70 and surface roughness was evaluated using an S-NEOX 3D Optical Profiler system (Sensofar Metrology, Terrassa, Spain).

## 2.2.2. Optical Characterization

For the optical characterization, the transmittance spectra under normal incidence in the visible range were recorded using a Shimazu UV-2100 spectrometer (Shimadzu Corporation, Kyoto, Japan).

#### 2.2.3. Electrical Characterization

DC and AC electrical measurements were performed at room temperature. The resistivity of the bulk  $MoO_x$  thin films grown on a silicon substrate at 0, 0.16, 0.32 and 0.48  $O_2$ /Ar flow rate ratio resulted from the measurement of I-V characteristics, and the apparatus used was a multimeter 34401A (HP) and a laboratory power supply ISO-Tech Dual Tracking with 5 V fixed model IPS 2303 DD (Isotech, Conchester, VT, USA). The thin film grown on PET-ITO substrate at 0.56  $O_2$ /Ar flow rate ratio was also analyzed by impedance spectroscopy (Agilent 4294A Precision Analyser, Santa Clara, CA, USA). For that, a matrix of gold electrode pads of 2 mm in diameter, separated by 2 mm, were deposited on the thin film surface using a sputter coater SC7620 (Quorum, East Sussex, UK). The connection between the sample and the instrument cell was made using micropositioners with gold probes.

#### 2.2.4. Electrochromic Characterization

Cyclic voltammetry experiments were performed inside a glove box with an argoncontrolled atmosphere, using a Methrom Autolab PGstat 302N, in a 3-electrode arrangement, with the samples as "working electrodes", a platinum wire as counter-electrode and a low leakage Ag/AgCl/KCl (3 M) as reference electrode. The voltammograms were swept from -2.5 to +2.0 V<sub>Ag/AgCl/KCl (3 M)</sub>. In situ and simultaneously with the cyclic voltammetry, optical transmittance was recorded using an Ocean Optics USB4000 spectrometer (Ocean Optics Inc., Dunedin, FL, USA).

## 3. Results and Discussion

## 3.1. Microstructural and Morphologic Characterization

## 3.1.1. Structural Analysis

Figure 1a presents the XRD patterns of the first series of  $MoO_x$  films deposited onto silicon substrates. All films deposited at a gas flow rate ratio of  $O_2/Ar$  less or equal to 0.32 are polycrystalline. The films deposited at flow rate ratios less or equal to 0.16  $O_2/Ar$  have body-centered cubic (bcc) molybdenum structure matching with JCPDS Card No. 04-002-0890 [34–36]. Films deposited at a flow rate ratio of 0.32  $O_2/Ar$  show face-centered cubic (fcc) molybdenum structure with correspondence with JCPDS Card No. 01-0882331. This transition, from bcc to fcc, can be related to strains induced by the oxygen accommodation on the crystal lattice [37] and it was responsible for the diffraction peak shift to lower angles. The calculated crystallite sizes (Table 1) showed a decrease from 28 nm to 13 nm with the increase of  $O_2/Ar$  flow ratio.



**Figure 1.** X-ray diffraction patterns of: (a) the first series of  $MoO_x$  thin films prepared with different gas flow rate ratio of  $O_2/Ar$  (all samples of the first series have metallic-like appearance); (b) the second series of  $MoO_x$  thin films prepared with the same gas flow rate ratio of  $O_2/Ar$  and different times of deposition (all samples have transparent appearance).

Sample –	Intense Peaks		Crystallite Size
	D-Spacing (Å)	(hkl)	(nm)
	2.24323	(110)	
0 O <sub>2</sub> /Ar-60 min	1.57868	(200)	$28\pm3$
_	1.29034	(211)	
	2.23039	(110)	
0.16 O <sub>2</sub> /Ar-60 min	1.58409	(200)	$16\pm3$
_	1.28323	(211)	
	2.42993	(111)	
0.32 O <sub>2</sub> /Ar-60 min –	2.10633	(200)	$- 13 \pm 2$
	1.48369	(220)	13 ± 2
	1.26528	(311)	

Table 1. Structural parameters of crystalline MoO<sub>x</sub> films and crystallite size (nm).

The increase of the oxygen content caused a broadening of the diffraction peaks. The sample deposited with the flow rate ratio of 0.48  $O_2/Ar$  showed a broadened band, centered at  $2\theta \approx 35^\circ$ , which suggested a coexistence between an amorphous structure and some few nanocrystallites dispensed in the matrix. At the flow rate ratio of 0.56  $O_2/Ar$ ,

corresponding to the second series of depositions, the diffractogram pattern showed in Figure 1b, with broadband centered at  $2\theta \approx 25^{\circ}$ , confirmed the formation of an amorphous structure based on the MoOx units [38–40]. In the samples grown in silicon, the intensity of the peak corresponding to silicon was observed only on samples with sputtering time below 60 min, which inferred a lower sample thickness.

## 3.1.2. Composition of the Films

Figure 2a-d show the RBS spectra of the MoO<sub>x</sub> thin films, deposited onto the glass-ITO substrate in the transition region (from metallic to transparent like appearance) and in transparent condition for different times of deposition. The RBS spectrum of the sample grown at the flow rate ratio of  $0.48 O_2/Ar$  was characterized by two barriers assigned to the Mo and O and no barriers assigned to the substrate were present (Figure 2a). The absence of RBS signal from the substrate is due to the high thickness of the film, which should be higher than 1250 nm, a value estimated using the NDF code (considering 6.5 g/cm<sup>3</sup> the volumetric density of MoO<sub>2</sub> [41]). A good fit was obtained assuming a homogenous composition profile in depth with a Mo/O ratio of  $\sim 0.5$ . However, it should be noticed that the sensitivity of RBS to O was strongly reduced due to its low mass and to the overlap of the O-signal with the Mo signal. The spectra obtained from the samples grown at 0.56 flow rate ratio of  $O_2/Ar$ , showed the end of the  $MoO_x$  layer meaning that these films were thinner, which is in agreement with XRD results that showed the appearence of the substrate (Si) peak. Furthermore, three extra barriers assigned to indium (In), calcium (Ca) and silicon (Si) from the substrate are marked in Figure 2b–d. To estimate the composition and the depth concentration profiles of Mo, these spectra were fitted using the NDF code (red lines in Figure 2). Regarding the film grown at the flow rate ratio of  $0.56 O_2/Ar$  for 60 min, the fit suggested a compositional gradient with an increase of the molybdenum concentration with depth. The Mo/O ratio as a function of the depth, estimated by the fit, changed from  $\sim 0.44$  at the surface to  $\sim 0.48$  at the interface with the substrate. The depth profile analysis, which was also a film thickness measurement, indicated that for an equal deposition time, the increase of  $O_2/Ar$  ratio promoted a thinner film. The sample with a 0.48 O<sub>2</sub>/Ar ratio had a thickness higher than 1250 nm. This thickness decreased to approximately 450 nm for the sample with 0.56  $O_2/Ar$  ratio. For an equal  $O_2/Ar$  ratio (0.56) and decreasing the sputtering time, the thickness decreased until around 200 nm for the sample growth with 30 min and 100 nm for the sample growth with 15 min. Another observation, for the samples grown at 30 and 15 min, is that the molybdenum profile was not homogeneous across the thickness. Both samples showed a maximum of molybdenum around half of the film thickness.

#### 3.1.3. Surface Chemistry Analysis

The thin film's surface chemistry was analyzed by X-ray photoelectron spectroscopy. Exhibited in Figure 3 is the high-resolution XPS spectrum of the Mo 3d state for samples grown at different  $O_2$  / Ar ratios. The different spectra from (a) to (e) showed the spin-orbit splitting characteristic of the Mo 3d core level. This spectrum was deconvoluted into two doublets, each of them related to a specific oxidation state. Fitting with Lorentzian-Gaussian line shapes and assuming a spin-orbit splitting of 3.13 eV and an intensity ratio between  $3d_{5/2}$  and  $3d_{3/2}$  of 1.5, revealed the presence of Mo<sup>6+</sup>, Mo<sup>5+</sup> and Mo<sup>0</sup> oxidation states in the sample prepared without the introduction of oxygen (Figure 3a). The presence of  $Mo^{6+}$  and  $Mo^{5+}$  oxidation states might come from surface oxidation phenomena once the as-deposited films were exposed to atmospheric conditions. Figure 3b shows the core level XPS spectrum for the film deposited at a flow rate ratio of  $O_2/Ar$  of 0.16, revealing that the surface composition of was similar to the previous. For the Mo<sup>0</sup> state, the Mo  $3d_{5/2}$ and Mo 3d<sub>3/2</sub> positions were found at 227.37 eV and 230.50 eV (Figure 3a) and at 227.83 eV and 230.96 eV (Figure 3b), respectively for 0 and 0.16  $O_2$  / Ar ratios. The Mo<sup>5+</sup> and Mo<sup>6+</sup> oxidation states had Mo 3d<sub>5/2</sub> and Mo 3d<sub>3/2</sub> peak positions at 231.17 eV and 234.30 eV, and 232.17 eV and 235.30 eV (Figure 3a) and at 231.53 eV and 234.66 eV and 232.60 eV

and 235.73 eV (Figure 3b), correspondingly. All these values were in agreement with the literature [42–46].

Figure 3c shows the core level XPS spectrum for representative thin films deposited at a flow rate ratio of  $O_2/Ar$  of 0.32. For the  $Mo^0$  state the Mo  $3d_{5/2}$  and Mo  $3d_{3/2}$  positions were found at 228.71 eV and 230.96 eV and the  $Mo^{5+}$  and  $Mo^{6+}$  oxidation states had Mo  $3d_{5/2}$  and Mo  $3d_{3/2}$  peak positions at 231.40 eV and 234.53 eV and 232.37 eV and 235.52 eV, respectively. On Figure 3d, corresponding to the thin films deposited at a flow ratio of  $O_2/Ar$  of 0.48, the oxidation state of the  $Mo^0$  was not detected but an extra  $Mo^{4+}$  oxidation state with Mo  $3d_{5/2}$  and Mo  $3d_{3/2}$  peak positions at 229.46 eV and 232.59 eV was observed [43,46]. The oxidation states  $Mo^{5+}$  and  $Mo^{6+}$  had Mo  $3d_{5/2}$  and Mo  $3d_{3/2}$  peak positions at 231.50 eV and 234.63 eV and at 232.44 eV and 235.57 eV, correspondingly. The presence of  $Mo^{5+}$  and  $Mo^{4+}$  oxidation states suggested the formation of  $MoO_2$  and  $Mo_2O_5$ . Increasing the flow rate ratio to 0.56  $O_2/Ar$  (Figure 3e), a single oxidation state was presently associated with the  $Mo^{6+}$  oxidation state with  $Mo 3d_{3/2}$  peak positions at 232.55 eV and 235.68 eV [45,46]. This spectrum was compatible with the presence of  $MoO_3$ . All binding energies values for the molybdenum oxides species are displayed in Table 2.



**Figure 2.** Rutherford Backscattering Spectrometry (RBS) spectrum obtained for measurements of the last sample from the first series: (a) deposited with a flow rate ratio of 0.48  $O_2$ /Ar for 60 min; and of the samples of the second series, prepared with the same flow rate ratio of 0.56  $O_2$ /Ar and different deposition times, (b) 60 min, (c) 30 min and (d) 15 min. The vertical arrows correspond to the surface/interface scattering energies of the different elements assigned to the film/substrate (sub).



**Figure 3.** High-resolution XPS spectra of the Mo 3d sate showing the spin-orbit splitting. The deconvoluted peaks associated with the  $3d_{3/2}$  and  $3d_{5/2}$  orbitals are also indicated for sample (**a**)  $0 O_2/Ar-60 \min$ , (**b**) sample  $0.16 O_2/Ar-60 \min$ , (**c**) sample  $0.32 O_2/Ar-60 \min$ , (**d**) sample  $0.48 O_2/Ar-60 \min$ , (**e**) sample  $0.56 O_2/Ar-60 \min$ .

Sample	Mo 3d Orbitals	Mo (0)	MoO <sub>2</sub> (IV)	Mo <sub>2</sub> O <sub>5</sub> (V)	MoO <sub>3</sub> (VI)
0 O <sub>2</sub> /Ar-	Mo 3d <sub>3/2</sub>	230.50	-	234.30	235.30
60 min	Mo 3d <sub>5/2</sub>	227.37	-	231.17	232.17
0.16 O <sub>2</sub> /Ar-	Mo 3d <sub>3/2</sub>	230.96	-	234.66	235.73
60 min	Mo 3d <sub>5/2</sub>	227.83	-	231.53	232.60
0.32 O <sub>2</sub> /Ar-	Mo 3d <sub>3/2</sub>	230.96	-	234.53	235.52
60 min	Mo 3d <sub>5/2</sub>	228.71	-	231.40	232.37
0.48 O <sub>2</sub> /Ar- 60 min	Mo 3d <sub>3/2</sub>	-	232.59	234.63	235.57
	Mo 3d <sub>5/2</sub>	-	229.46	231.50	232.44
(0	Mo 3d <sub>3/2</sub>	-	-	-	235.67
	Mo 3d <sub>5/2</sub>	-	-	-	232.56
From literature	Mo 3d <sub>3/2</sub>	230.85	232.6	234.62	235.85
[41,42,44]	Mo 3d <sub>5/2</sub>	227.70	229.4	231.63	232.65

Table 2. Binding energies (eV) of the Mo 3d orbitals obtained by XPS and from literature.

#### 3.1.4. Raman Analyses

The Raman spectra of the first and second series of thin films are presented in Figure 4a,b. In general, the Raman spectra of the deposited thin films are in agreement with the principle that the structure is formed by  $MoO_6$  octahedra, where stretching and bending vibrational modes occur in the regions 200–400 cm<sup>-1</sup> and 500–1000 cm<sup>-1</sup>, respectively, and the narrow Raman bands detected between 920–1000 cm<sup>-1</sup> are due to the vibration of terminal Mo=O bonds [47].



**Figure 4.** Raman spectra of the: (a) the first series of  $MoO_x$  films grown at a different flow rate ratio of  $O_2$ /Ar; (b) the second series of  $MoO_x$  films deposited with a 0.56 flow rate ratio of  $O_2$ /Ar and different deposition times.

The Raman spectra of the samples deposited at the flow rate ratios up to 0.48  $O_2/Ar$  were very similar, presenting one broadband and three broad peaks. The broad featureless band observed between 200 cm<sup>-1</sup> and 400 cm<sup>-1</sup> could be assigned to amorphous structural nature, the first broad peak located at 520 cm<sup>-1</sup> was related to the optical phonon vibrations of the silicon substrate and the bands located at 817 cm<sup>-1</sup> and 940 cm<sup>-1</sup>, more intense and defined, could be related to the stretching vibration of doubly-connected bridge oxygen (Mo-O-Mo) groups and Mo = O vibration bond [47,48], respectively. It is important to mention that the as-deposited samples were exposed to atmospheric conditions before Raman analyses, corroborating the presence of MoO<sub>3</sub> oxides as also observed by XPS.

Increasing the flow rate ratio to 0.56  $O_2/Ar$  (Figure 4b), a shift of the 817 cm<sup>-1</sup> and 940 cm<sup>-1</sup> bands to higher wavenumbers was observed, indicating a predominance of the MoO<sub>3</sub> base structure [49,50].

## 3.1.5. Morphological and Thin Film Growth Features

The two-dimensional (2D) surface and cross-section morphologies of  $MoO_x$  thin films deposited on silicon and PET-ITO substrates have been investigated using SEM images as shown in Figures 5–7.

The SEM micrographs show that Mo films grown in pure Ar atmosphere, exhibited a well-defined structure with flat elongated plate-like shaped grains in a top view and columnar grain growth in cross-section [37]. The samples deposited with a ratio of  $O_2/Ar$  of 0.16, 0.32 and 0.48 present islands and clusters in top view that became more expressive with the increase of oxygen content. The cross-section micrographs revealed a columnar grain growth that became dense and featureless with the increase of the  $O_2$  flow rate.



**Figure 5.** Surface SEM images of the first series of  $MoO_x$  films deposited on a silicon substrate with (**a**)  $OO_2/Ar-60$  min; (**b**)  $0.16 O_2/Ar-60$  min; (**c**)  $0.32 O_2/Ar-60$  min; (**d**)  $0.48 O_2/Ar-60$  min. Cross-section SEM images of the first series of  $MoO_x$  films deposited on a silicon substrate with (**e**)  $0 O_2/Ar-60$  min; (**f**)  $0.16 O_2/Ar-60$  min; (**g**)  $0.32 O_2/Ar-60$  min; (**h**)  $0.48 O_2/Ar-60$  min; (**f**)  $0.16 O_2/Ar-60$  min; (**g**)  $0.32 O_2/Ar-60$  min; (**h**)  $0.48 O_2/Ar-60$  min; (**f**)  $0.16 O_2/Ar-60$  min; (**g**)  $0.32 O_2/Ar-60$  min; (**h**)  $0.48 O_2/Ar-60$  min; (**f**)  $0.16 O_2/Ar-60$  min; (**g**)  $0.32 O_2/Ar-60$  min; (**h**)  $0.48 O_2/Ar-60$  min; (**f**)  $0.16 O_2/Ar-60$  min; (**g**)  $0.32 O_2/Ar-60$  min; (**h**)  $0.48 O_2/Ar-60$  min; (**f**)  $0.16 O_2/Ar-60$  min; (**g**)  $0.32 O_2/Ar-60$  min; (**h**)  $0.48 O_2/Ar-60$  min.

The sample deposited at 0.56 ratio of  $O_2/Ar$ , Figure 6, revealed a compact/dense and inexpressive morphology on both top view and cross-sections, without evidence of the presence of cracks or porosity.

In Figure 7 it is possible to observe the morphology of the thin films grown on PET-ITO substrates, which was similar to the one reported for the films deposited on silicon. Moreover, it was visible that the thickness decreased from around 910 nm (sample  $0.56 \text{ O}_2/\text{Ar-60}$  min) to 400 nm (sample  $0.56 \text{ O}_2/\text{Ar-30}$  min). Sample  $0.56 \text{ O}_2/\text{Ar-15}$  min showed a thickness of around 190 nm.

These film thicknesses were slightly different from the ones calculated through RBS measurements (Section 3.1.2) since those calculations depended on the material's density in the bulk form and not of the thin film. Taking this feature into account, Table 3 shows the results of a deeper study of RBS spectra, calculating the film thickness considering different stoichiometries and calculating the film density considering the thickness estimated from the SEM measurements. The calculated density results were lower than each possible stoichiometry of molybdenum oxide, suggesting the existence of voids in the film structure. Moreover, and considering a possible statistical error of about 10 to 20% in the calculated values, the density was approximately the same for all samples indicating a similar stoichiometry. From this, it was also suggested that the calculated values were an average of a mixture of several oxidation states of molybdenum.



**Figure 6.** Surface SEM images of the second series of  $MoO_x$  films grown on a silicon substrate with (**a**) 0.56  $O_2$ /Ar-60 min; (**b**) 0.56  $O_2$ /Ar-30 min; (**c**) 0.56  $O_2$ /Ar-15 min. Cross-section SEM images of the second series of  $MoO_x$  films grown on a silicon substrate with (**d**) 0.56  $O_2$ /Ar-60 min; (**e**) 0.56  $O_2$ /Ar-30 min; (**f**) 0.56  $O_2$ /Ar-15 min.



**Figure 7.** Surface SEM images of the second series of  $MoO_x$  films grown on polyethylene terephthalate with indium tin oxide (PET-ITO) substrate with (**a**) 0.56  $O_2$ /Ar-30 min; (**b**) 0.56  $O_2$ /Ar-15 min. Cross-section SEM images of the second series of  $MoO_x$  films grown on polyethylene terephthalate with indium tin oxide (PET-ITO) substrate with (**c**) 0.56  $O_2$ /Ar-30 min; (**d**) 0.56  $O_2$ /Ar-15 min.

**Table 3.** Samples thickness calculated using RBS fit results and considering different molybdenum oxides and density estimation considering the thickness measured by SEM.

Sample –	Areal Density	Thickness (nm)		+ Density
	(10 <sup>15</sup> atm/cm <sup>2</sup> )	*	**	(g/cm <sup>3</sup> )
0.48 O <sub>2</sub> /Ar-60 min	429,949	-	-	-
0.56 O <sub>2</sub> /Ar-60 min	4900	385	531	2.74
0.56 O <sub>2</sub> /Ar-30 min	1658	122	166	1.97
0.56 O <sub>2</sub> /Ar-15 min	862	61.4	84.7	2.09

\* considering the MoO<sub>2</sub> density 6.47 g/cm<sup>3</sup> [51]; \*\*considering the MoO<sub>3</sub> density 4.69 g/cm<sup>3</sup> [16]; \* Density estimated considering the stoichiometry obtained from fit of RBS data and the thickness estimated from the SEM measurements.

Table 4 shows the roughness of the samples grown at the 0.56 O<sub>2</sub>/Ar flow rate ratio with different deposition times by optical profilometry. Samples deposited on the PET-ITO substrate presented the lowest roughness, between 6 and 8 nm. To confirm the values of density estimated by RBS, XRR measurements were carried out for these samples without the need of using a monochromator despite the very thick layers involved. The XRR curves were characterized by a region of low incident angles for which the X-ray beam intensity was totally reflected and a region above a critical angle ( $\theta_c$ ) for which the intensity fell rapidly with the incident angle with a strong dependence on the roughness. By simulating the curves, it was possible to estimate the complex refractive indexes of the layered structure, *n*. The *n* value is given by:

$$n = (1 - \delta) - i\beta \tag{1}$$

where  $\delta$  parameter ranges from  $10^{-5}$  to  $10^{-6}$  for X-ray wavelength of approximately 1 Å, and  $\beta$  is related to the X-ray absorption. The critical angle  $\theta_c$  is given by:

$$\theta_c = \sqrt{2\delta} \tag{2}$$

 $\delta$  in turn, is proportional to  $\rho$  through (Equation (3)):

$$\delta = \frac{\frac{r_e \lambda^2}{2\pi} N_0 \rho \sum_i x_i (Z_i + f'_i)}{\sum_i (x_i M_i)}$$
(3)

where  $r_e$  is the classical electron radius,  $N_o$  is the Avogrado number,  $\lambda$  is the X-ray wavelength,  $\rho$  is the density,  $x_i$  is the atomic ratio of atom i,  $Z_i$  is the atomic number of atom i,  $M_i$  is the atomic weight of the atom i and f' is the real part anomalous dispersion term of the atomic scattering factors for atom i. By replacing above the f' with f'' (imaginary part of the atomic scattering factors) the third term in the refrective index n,  $\beta$ , is derived (Equation (4)) [52].

$$\beta = \frac{\frac{r_e \lambda^2}{2\pi} N_0 \rho \sum_i x_i (Z_i + f_i'')}{\sum_i (x_i M_i)}$$
(4)

The simulations were performed using the MROX code following the dynamical theory recursive method developed by Parratt [53]. Furthermore, due to the impossibility of obtaining a reasonable fit considering homogeneous MoO<sub>3</sub> layers, a function to estimate the contribution of pores using the kinematical model developed by Maaze et al. [54] was included. The contribution of the pores in the X-ray reflectivity scans was visible around 1.42° where a small shoulder suggested the small-angle X-ray scattering due to the randomly dispersed pores in the matrix. For the simulations, and to estimate the density, pores with aspherical shape and an averaged inter-distances between 10 nm and 500 nm were assumed.

The fits were performed using a standard genetic algorithm and are presented in Figure 8. A density of around 2.4 g/cm<sup>3</sup> was obtained for the surface layer, which was in the same level of the ones obtained by RBS (Table 3) and lower than the bulk density of  $MoO_3$  (4.69 g/cm<sup>3</sup>). The density derived was also an indication of an oxygen overstoichiometry at the surface, as RBS also suggested. Although a clear improvement of the fitting was possible, considering the contribution of pores, the derived thickness in some samples did not match with the ones observed by SEM. This can be due to the limitation of this technique for the estimation of thickness when interference oscillations were not observed or it could suggest that other reasons were affecting the angular dependency of the XRR intensity at angles above the critical angle that were not considered in the fit. Small soller-slit mis-alignments or crystal heterogeneities, such as in the thickness along the high-area probed sample due to the small X-ray angle scattering events are possible



reasons. However, it is out of the scope of this manuscript to deeper the outputs of the X-ray reflectivity simulation.

**Figure 8.** X-ray reflectivity (XRR) spectra and simulations of the film's growth at the  $0.56 \text{ O}_2/\text{Ar}$  flow rate during (**a**) 60, (**b**) 30 and (**c**) 15 min.

Table 4. Surface roughness (root mean square neight, sq) measured by optical promotierry of me
samples grown at flow rate of $0.56 \text{ O}_2/\text{Ar}$ with different deposition times.

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Sample	Substrate	Sq (nm)
0.56 O <sub>2</sub> /Ar-15 min	PET-ITO	8
0.56 O <sub>2</sub> /Ar-30 min	rei-iio	6
0.56 O <sub>2</sub> /Ar-15 min	Glass-ITO	11
0.56 O <sub>2</sub> /Ar-30 min		13
0.56 O <sub>2</sub> /Ar-60 min		42

## 3.2. Optical Response of the Thin Films

Figure 9a,b present the transmittance spectra of MoO<sub>x</sub> films grown on PET-ITO and glass-ITO substrate deposited at 0.56 O<sub>2</sub>/Ar flow rate ratio with different deposition times. It was found that the MoO<sub>x</sub> films deposited onto PET-ITO substrates exhibited higher transmittance, in the visible range, than the ones deposited onto glass-ITO substrates (maxima of 89% and 87%, respectively). The phenomenon observed in Figure 9a related to higher transmittance of MoO<sub>x</sub> films deposited on PET-ITO substrate was already observed by [55,56] and it was justified by an improvement of the anti-reflective property of the ITO layer after the MoO<sub>x</sub> film deposition, increasing the global transmittance. The energy band gaps for direct and indirect transitions of MoO<sub>x</sub> films were obtained by the Tauc method [40], by extrapolation of the linear portion of a plot of  $(\alpha h\nu)^2$  versus  $h\nu$  for the direct

transition, and  $(\alpha h\nu)^{\frac{1}{2}}$  versus  $h\nu$  for indirect transition, where  $\alpha$  is the optical absorbance coeficiente and  $h\nu$  the photon energy. The maximum energy bandgap value obtained for MoO<sub>x</sub> films grown on PET-ITO substrate was 3.92 eV and for MoO<sub>x</sub> films grown on glass-ITO substrate was 3.93 eV. A blue shift of the bandgap value was observed with the increase of the thickness, which can be explained in terms of the extended disorder and presence of defects that could be related to the existence of substoichiometric MoO<sub>3</sub> phases [40]. This direct bandgap energy behaviour could be related to a gradual increase of oxygen ion vacancies in the film with the deposition time [57].



**Figure 9.** Transmission spectra and plots of  $(\alpha h\nu)^2$  versus  $h\nu$  for direct transition in MoO<sub>x</sub> films (**a**) grown on PET-ITO substrate and (**b**) grown on glass-ITO substrate at 0.56 O<sub>2</sub>/Ar flow rate ratio at different deposition time.

## 3.3. Electrical Characterization

The electrical resistivity, at room temperature, of the thin films with  $O_2/Ar$  flow ratio between 0 (oxygen absence) and 0.48 was measured using the four-probe method [58] (Figure 10a). The samples resistivity was calculated using the Equation (5):

$$\rho = \frac{V}{i}wF_3 \tag{5}$$

where  $\rho$  represents the resistivity, *V* the measured voltage, *i* the applied current, *w* the film thickness and *F*<sub>3</sub> a correction factor related to the geometry of the sample [59,60]. Correction factors were obtained through geometric series with further calculations. In the correspondent employed physical modes, the contact area was assumed infinitesimal (true in very good approximation). Geometrically, the four-point probes were collinear and equidistant. As shown in Figure 10, the electrical resistivity of the films depended on the oxygen concentration used in the sputtering atmosphere. The metallic Mo sample (Mo<sup>0</sup>) showed a resistivity of 65.7  $\mu$ Ω.cm which was in agreement with reported studies [61]. Increasing the oxygen content, the resistivity increased to the order of hundreds of  $\mu$ Ω.cm. The sample prepared at 0.48 flow rate ratio of O<sub>2</sub>/Ar presented the highest resistivity value, in the order of thousands  $\mu$ Ω.cm. This sample presented a metallic-like appearance and the value of resistivity could be assigned to the presence of MoO<sub>2</sub> [62]. With further increase of the oxygen concentration in the sputtered atmosphere, optical transparency was achieved, and the electrical properties changed significantly. The samples became highly resistive, a feature that was in agreement with the already results discussed.

Figure 10b presents the dielectric constant ( $\varepsilon'$ ) behavior as a function of the frequency, measured by impedance spectroscopy [63], for the MoO<sub>x</sub> transparent films deposited on PET-ITO substrates, for 30 and 15 min. Both samples showed similar profiles, but the sample 0.56 O<sub>2</sub>/Ar-15 min showed a higher  $\varepsilon'$  in the high-frequency region. In the low-frequency region, the  $\varepsilon'$  was high due to the Maxwell–Wagner polarization effect, i.e., due to polarization phenomenon arising from sample surface-electrode interfaces, although some possible effect of space charge due to injection electrodes, cannot be excluded. The

decrease of  $\varepsilon'$  with the increase of the frequency was expected as part of the relaxation process that should occur at very low frequency. At 1 kHz, and at room temperature, the  $\varepsilon'$  of the sample deposited for 15 min (~5.2) was higher than the one grown for 30 min (~4.4). This difference should be related to the fact that the first one has a MoO<sub>x</sub> stoichiometry with x < 3 (Figure 2). Comparing these results with the ones from the literature [64,65] that shows, in similar experimental conditions, higher dielectric constants, it is possible to claim that the differences observed might be the result of the existence of voids in the films microstructure.



**Figure 10.** (a) Electrical resistivity of  $MoO_x$  thin films at different  $O_2/Ar$  flow rate ratio; (b) Dielectric constant and loss of the samples 0.56  $O_2/Ar$ -30 min and 0.56  $O_2/A$ -15 min.

#### 3.4. Electrochromic Characterization

 $MoO_3$  thin films, with  $Mo^{+6}$  sites, are commonly colorless. The intercalation of charges into the film during the redox reaction, where  $Mo^{+6}$  ions are reduced to  $Mo^{+5}$  causes a blue coloration of the films. Further deintercalation of charges, forces the  $Mo^{+5}$  sites in the film to be reoxidized to  $Mo^{+6}$  and the film becomes to the original colorless state. Coloration and bleaching states are guided by the following electrochemical redox reaction [5,66]:

$$MoO_3 (colourless) + yLi^+ + ye^- \leftrightarrow Li_yMoO_3 (dark)$$
 (6)

On the  $Li_yMoO_3$  chemical formula Equation (1), y is named "insertion coefficient" and represents the fractional number of sites in  $MoO_3$  lattice able to be filled by Li<sup>+</sup> ions. Usually, the final color observed with almost  $MoO_3$  filled is blue (dark). However, in the present work, the lattice has some  $MoO_x$  incorporation (as supported by XPS data) and the final reduction reaction gave a dark brown color.

The cyclic voltammetry experiments were conducted at the sweep rate of 10 mV/s in the 1M electrolyte solution of LiClO<sub>4</sub> dissolved in propylene carbonate. It is important to mention that, for MoO<sub>3</sub> films, cyclic voltammetry experiments should not be performed in aqueous solutions unless the surfaces are chemically modified to create stability [23]. Figure 11 presents typical cyclic voltammetry (CV) curves of MoO<sub>x</sub> films grown on PET-ITO substrate at 0.56 flow rate ratio of O<sub>2</sub>/Ar and different deposition times. It was found that when the MoO<sub>x</sub> films underwent a cathodic sweep (from positive to negative voltage), a continuous enhancement in the dark brown color of the film was observed. On contrary, during an anodic sweep, the film bleached and became almost colorless.

Figure 12a,b presents the evolution of relative transmittance for both samples studied. The optical transmittance decreased over time, indicating that some Li<sup>+</sup> ions intercalated during the coloration cycle could not be deintercalated, resulting in the retention of a small fraction of Li<sup>+</sup> ions. During the cathodic sweep, the current density increased continuously and attained a negative maximum value at the potential of  $-2.5 V_{Ag/AgCl/KCl (3M)}$  and the film exhibited a deep brown coloration when electrons and Li<sup>+</sup> ions intercalated

into the MoO<sub>x</sub> films. Changing the potential from -2.5 to  $2.0 V_{Ag/AgCl/KCl (3M)}$  during the anodic excursion, the current density changed from negative to positive, attaining a maximum (positive) value, decreasing after until the film was bleached. The bleaching process occurred due to the deintercalation of the electrical charges from the film structure. During the anodic scan, the film exhibitd a broad peak, whereas during cathodic scan a characteristic "spike" could be seen at the potential of  $-2.5 V_{Ag/AgCl/KCl (3M)}$  (Figure 11), in agreement with previous results [67,68]. The absence of a well-defined peak in the cathodic region was due to the back electromotive force (back emf) formed within the LiMoO<sub>x</sub> (molybdenum bronze) during ion insertion [69]. In Figure 11 it is evident that the cathodic and anodic peak current densities increased upon deposition time. This fact is supported by the amount of intercalated and deintercalated Li<sup>+</sup> ions determined by integration of the cathodic and anodic current, respectively, as listed in Table 5.



**Figure 11.** Comparison of typical CV curves of  $MoO_x$  thin films grown on PET-ITO substrate with different deposition times at a scan rate of 10 mV/s. The arrows indicate the scan direction.



**Figure 12.** Evolution of representative transmittance of the sample (**a**)  $0.56 \text{ O}_2/\text{Ar-15}$  min and (**b**)  $0.56 \text{ O}_2/\text{Ar-30}$  min scanned at a rate of 10 mV/s.  $\Delta$ T represents the relative optical transmittance difference between the colored and bleached states.

Sampl	e	0.56 O <sub>2</sub> /Ar-15 min	0.56 O <sub>2</sub> /Ar-30 min	
Scan rate (mV/s)		10	10	
Cathodic current density, j <sub>pc</sub> (A/cm <sup>2</sup> )		$7.57 imes10^{-4}$	$6.83 imes10^{-4}$	
Cathodic peak potential, V <sub>pc</sub> (V <sub>Ag/AgCl/KCl (3M)</sub> )		-2.50	-2.46	
Anodic current density, j <sub>pa</sub> (A/cm <sup>2</sup> )		$5.01 imes10^{-4}$	$6.56 imes10^{-4}$	
Anodic peak potential, V <sub>pa</sub> (V <sub>Ag/AgCl/KCl (3M)</sub> )		-1.00	-0.68	
Diffusion coefficient (cm <sup>2</sup> s <sup>-1</sup> )	D <sub>insertion</sub>	$7.74 imes10^{-10}$	$6.31  imes 10^{-10}$	
	D <sub>extraction</sub>	$3.39 imes10^{-10}$	$5.81  imes 10^{-10}$	
Intercalated charge (C/cm <sup>2</sup> )		0.077	0.109	
Deintercalated charge (C/cm <sup>2</sup> )		0.067	0.099	
Reversibility (%)		87%	91%	
Coloration efficiency at 630 nm ( $cm^2/C$ )		19.6	10.2	

Table 5. Cyclic voltammetry measurements of the  $0.56 O_2/Ar$  films growth at 15 and 30 min.

The film color reversibility, which correlates the intercalated/deintercalated charge, showed a higher value for the sample with a thickness of around 400 nm (deposition time of 30 min). On the other hand, the anodic current peak position in the thinner sample (~190 nm) shifted towards the lower potential inferring that the process involving a change from colored to bleached state was faster for the sample deposited with lower deposition time. The CV curve areas were found to increase with the increase of the deposition time (thinner), suggesting an increase of the intercalation capacity.

The diffusion coefficient of  $Li^+$  ions during the intercalation and deintercalation, into  $MoO_x$  host lattice, could be obtained by the Randles–Sevick relationship [29]:

$$j_p = 2.72 \ x \ 10^5 n^{\frac{3}{2}} D^{\frac{1}{2}} C v^{\frac{1}{2}} \tag{7}$$

where  $j_p$  is the peak current density (A/cm<sup>2</sup>), *n* is the number of electrons involved in the redox process (assumed to be 1), *D* is the diffusion coefficient of Li<sup>+</sup> ions (cm<sup>2</sup>/s), *C* is the concentration of active ions in the electrolyte solution (mol/cm<sup>3</sup>), and *v* is the scan rate (V/s). Table 5 shows these different electrochromic parameters for the MoO<sub>x</sub> films. It is observed that the diffusion coefficient of Li<sup>+</sup> ions for both process intercalation and deintercalation was faster for the thinner sample (190 nm—deposition time of 15 min) and the variation on the mobility of the diffusing ionic species could be related with the texture of the film [68], suggesting that low roughness promoted faster inter-deintercalation processes.

Coloration efficiency  $(\eta)$  is another very important electrochromic parameter that correlates the change in optical density with the charges intercalated per unit of electrode area by the following relationship [70]:

$$\eta(\lambda) = \frac{\Delta OD}{\left(\frac{Q}{A}\right)} = \frac{ln\left(\frac{T_{bleached}}{T_{coloured}}\right)}{\left(\frac{Q}{A}\right)}$$
(8)

where  $\Delta OD$  is the change in optical density,  $\frac{T_{bleached}}{T_{coloured}}$  are the transmittances of the film in colored and bleached states, at a certain wavelength, Q is the intercalated charge, and A is the film electrode area. The maximum coloration efficiency obtained was 19.6 cm<sup>2</sup>/C (Table 5), observed on the MoO<sub>x</sub> film grown for 15 min. Although not very high, such a value was in agreement with some previous work found in the literature [71] considering the device stability. Moreover, in the present work, the material was deposited on flexible substrates with low surface roughness.

## 4. Conclusions

Films of  $MoO_x$ , with thickness from 190 up to 910 nm, were prepared by DC magnetron sputtering in stiff and flexible substrates.

XRD analysis shows a transition from a crystalline Mo structure to a MoOx amorphous state, occurring between 0.48 and 0.56  $O_2/Ar$  flow rate ratio. XPS indicates the chemistry evolution, with the increase of the  $O_2/Ar$  ratio, from Mo<sup>0</sup> up to the Mo<sup>6+</sup> oxidation state. Raman corroborates these results. It was found a density value of around 2.4 g/cm<sup>3</sup>, supported by XRR. RBS revealed a decrease of the Mo/O ratio with the rise of  $O_2/Ar$  flow rate ratio and time of deposition. It showed that the molybdenum profile is not homogeneous across the thickness.

The transparent and flexible films present an optical transmittance of 89%, in the visible range and an energy bandgap value of 3.92 eV. The dielectric response of these films revealed an  $\varepsilon'$ , at 300 K and 1 kHz, between 4.4 and 5.2 depending on the time of growth. The electrochromic response showed that the 0.56 O<sub>2</sub>/Ar-15 min film exhibits the maximum coloration efficiency (19.6 cm<sup>2</sup>/C), but the 0.56 O<sub>2</sub>/Ar-30 min film shows better reversibility.

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