



# Article Adsorption and Oxidation of CO on Ceria Nanoparticles Exposing Single-Atom Pd and Ag: A DFT Modelling

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Abstract: Various CO<sub>x</sub> species formed upon the adsorption and oxidation of CO on palladium and silver single atoms supported on a model ceria nanoparticle (NP) have been studied using density functional calculations. For both metals M, the ceria-supported MCO<sub>x</sub> moieties are found to be stabilised in the order MCO < MCO<sub>2</sub> < MCO<sub>3</sub>, similar to the trend for CO<sub>x</sub> species adsorbed on M-free ceria NP. Nevertheless, the characteristics of the palladium and silver intermediates are different. Very weak CO adsorption and the small exothermicity of the CO to CO<sub>2</sub> transformation are found for O<sub>4</sub>Pd site of the Pd/Ce<sub>21</sub>O<sub>42</sub> model featuring a square-planar coordination of the  $Pd^{2+}$  cation. The removal of one O atom and formation of the O<sub>3</sub>Pd site resulted in a notable strengthening of CO adsorption and increased the exothermicity of the CO to CO<sub>2</sub> reaction. For the analogous ceria models with atomic Ag instead of atomic Pd, these two energies became twice as small in magnitude and basically independent of the presence of an O vacancy near the Ag atom. CO<sub>2</sub>-species are strongly bound in palladium carboxylate complexes, whereas the CO<sub>2</sub> molecule easily desorbs from oxide-supported AgCO<sub>2</sub> moieties. Opposite to metal-free ceria particle, the formation of neither PdCO<sub>3</sub> nor AgCO<sub>3</sub> carbonate intermediates before CO<sub>2</sub> desorption is predicted. Overall, CO oxidation is concluded to be more favourable at Ag centres atomically dispersed on ceria nanostructures than at the corresponding Pd centres. Calculated vibrational fingerprints of surface CO<sub>x</sub> moieties allow us to distinguish between CO adsorption on bare ceria NP (blue frequency shifts) and ceria-supported metal atoms (red frequency shifts). However, discrimination between the CO2 and  $CO_3^{2-}$  species anchored to M-containing and bare ceria particles based solely on vibrational spectroscopy seems problematic. This computational modelling study provides guidance for the knowledge-driven design of more efficient ceria-based single-atom catalysts for the environmentally important CO oxidation reaction.

**Keywords:** CeO<sub>2</sub>-based nanomaterials; density functional calculations; single-atom catalysts; structure; reactivity

## 1. Introduction

Ceria, as a component of catalysts containing transition metals (M) Pd or Ag, is used in numerous applications ranging from the abatement of soot, volatile organic compounds, and CO [1–10] to the production of syngas [11] and CO<sub>2</sub> activation [12,13]. As an active reducible support, ceria facilitates the dispersion of metals and MO<sub>x</sub> phases on the surface [7,14–17] and provides lattice O atoms to oxidise reactants [2,9,17–20]. For instance, the interactions within Pd–ceria interfaces allow the synergistic oxidation/reduction of both subsystems [21], promote the oxidation of CO by lattice O atoms, and the oxidation of the reduced ceria by O atoms of CO<sub>2</sub> [22]. Supported transition metals can also enhance



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**Copyright:** © 2021 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). the redox performance and oxygen storage capacity of ceria [23]. Often, high catalytic efficiency is achieved using a nanostructured ceria support via enhanced metal–support interaction, which improves the dispersion of metal particles and suppresses their sintering at elevated temperatures [3,4,7–9,11,19,24–27].

The M-containing surface phases of the aforementioned systems are represented by M<sub>m</sub> [6,7,16,24,28,29] and MO<sub>x</sub> [6,25,26,30] nanoparticles (NPs), charged metal clusters [7,24,31], solid  $M_xCe_{1-x}O_{2-\delta}$  solutions [25–27,32–35], and dispersed  $M_1$  or  $O_xM_1$  adspecies [3,6-9,16,36,37]. Analysis of the crystalline environment of the Pd<sub>1</sub> ad-species revealed that each Pd<sup>2+</sup> ion in Pd/CeO<sub>2</sub> catalysts prepared by the solution combustion method is coordinated, on average, by three O atoms [34]. This coordination mode of  $Pd_1$ is a feature of adsorption complexes with CO such as  $O_2Pd_1$ -CO/CeO<sub>2</sub>(11) [8],  $O_1Pd_1$ -CO/CeO<sub>2</sub>(111) [8], and O<sub>1</sub>Pd<sub>1</sub>-CO/CeO<sub>2</sub>(100) [9], while a Pd<sub>1</sub>-CO/CeO<sub>2</sub>(100) complex exhibits an O-Pd-O bridge [9]. In many cases,  $Pd_1$  is in a square-planar environment.  $Pd_1$ centres in Ce<sub>1-x</sub>Pd<sub>x</sub>O<sub>2- $\delta$ </sub> crystals (x  $\leq$  0.15) reside on O<sub>4</sub> units adjacent to Ce centres [32]. The doping of ceria with Pd results in a structure with the dopant ion displaced from the initial cationic position to the centre of the  $O_4$  unit [23]. Furthermore,  $Pd_1$  species are attached to the  $O_4$  unit of the  $Pd_xCe_{1-x}O_{2-x-\delta}$  lattice incorporating products of water dissociation [27]. Replacing every second upper-layer Ce<sup>4+</sup> cation and one adjacent to it  $O^{2-}$  anion on the CeO<sub>2</sub>(110) surface with Pd<sup>2+</sup> leads to a complex reconstruction and a low-energy surface geometry, with the dopant ion residing close to the centre of the square-planar O<sub>4</sub> site [6]. Stable structures with square-planar O<sub>4</sub>Pd are also communicated on Pd-doped  $CeO_2(111)$  [33] and edges of ceria NPs at intersecting  $\{111\}$  and  $\{100\}$ nanofacets [35]. Four-fold coordinated Pd adatoms are identified in the most stable O4Pd structures on the  $CeO_2(110)$  surface [20] and {100} nanofacets [37]. The surface  $O_4$  sites are also capable of suppressing the sintering of Ag<sub>1</sub> species, despite the fact that the Ag atom binds to the {100}-O<sub>4</sub> pocket more weakly than other Group VIII–XI metal atoms [37,38]. The aforementioned thermally stable structures are relevant to the development of the single-atom catalysts [36,37,39].

To understand how the role of the ceria support varies in specific catalytic processes, it is crucial to examine the interactions of the involved reactants with various active centres of CeO<sub>2</sub>. For the ceria-supported metal catalysts of CO oxidation or CO<sub>2</sub> activation, of primary interest are the interactions of  $O_2$ , CO, and CO<sub>2</sub> molecules with the metal-support interfaces [2,8,9,18,20,33,40–43]. Thus, modelling based on density functional theory (DFT) has recently addressed a variety of sites with M1-O2, M1-CO, and M1-CO2 entities on ceria [8,9,18,20,40]. The adsorption of CO on a single Pd atom embedded in the defectfree CeO<sub>2</sub>(111) surface and that containing O vacancies followed by NO reduction with CO was explored [18,40]. Surface complexes of CO and  $CO_2$  taking part in the catalytic cycle of CO oxidation on Pd1/CeO2(110), including Pd1-CO, Pd1-CO2, O1Pd1-CO species on the stoichiometric CeO<sub>2</sub>(110) surface and Pd<sub>1</sub>-CO<sub>2</sub>, Pd<sub>1</sub>-O<sub>2</sub>, O<sub>2</sub>-Pd<sub>1</sub>-CO ones on the O-deficient  $CeO_2(110)$  surface, were calculated [20]. The CO oxidation routes passing via O<sub>1</sub>Pd<sub>1</sub>-CO, Pd<sub>1</sub>-CO, Pd<sub>1</sub>-CO<sub>2</sub>, O<sub>2</sub>Pd<sub>1</sub>-CO, and O<sub>1</sub>Pd<sub>1</sub>-CO<sub>2</sub> moieties on the defectfree CeO<sub>2</sub>(111) surface [8] as well as on regular and O-deficient CeO<sub>2</sub>(100) surfaces [9] were also quantified. The catalytic CO oxidation according to the Mars-van-Krevelen mechanism combines the elementary steps of oxygen donation from a surface active centre to adsorbed CO and the subsequent replenishment of the support by stream oxygen; the much slower conversions of the first step are found to be rate-determining [8,9,19,20]. In addition to the M<sub>1</sub>-CO<sub>2</sub> structures, surface carbonate complexes can be formed on M<sub>1</sub>-ceria interfaces before the desorption of CO<sub>2</sub>. To this end, the formation of tridentate Pd<sub>1</sub>-CO<sub>3</sub> carbonates upon  $CO_2$  adsorption at the interface of  $Pd_1$  and O-deficient  $CeO_2(111)$  surface was simulated [13] and the CO vibration frequencies of various ceria-supported Pd-CO species were calculated [8,9]. Unlike the quite extensive computational studies of  $Pd_1CO_x$ ceria systems outlined above, no simulations of analogous Ag<sub>1</sub>CO<sub>x</sub>-ceria systems have been communicated so far to the best of our knowledge.

Previous studies have developed structural models of low-energy CeO<sub>2</sub> NPs [44–47] and established that their {100} nanofacets notably stabilise single d-metal atoms [38,41,44,48–51]. In this work, we consider monoatomic Pd and Ag species located on a Ce<sub>21</sub>O<sub>42</sub> NP [49] as models appropriately describing surface composites formed by single-atom Pd and Ag with nanostructured ceria. These two metals, which are neighbouring in the Periodic Table, interact very differently with ceria and behave as M-based species involved in CO oxidation. The quantification and in-depth understanding of such differences are still missing in the literature.

This study aims to (i) determine the structures of the lowest-energy complexes with  $CO, CO_2$ , and  $CO_3^{2-}$  moieties resulting from the interaction of CO with Pd and Ag single atoms anchored to the O<sub>4</sub>-pocket sites of the stoichiometric and O-deficient ceria NPs, (ii) analyse the structure and properties of these nanostructured adsorption systems versus earlier investigated analogues formed on extended ceria surface containing M<sub>1</sub> centres, (iii) evaluate and rationalise the reactivity differences of Pd<sub>1</sub>/NP{100} and Ag<sub>1</sub>/NP{100} sites as active centres for CO oxidation (including the effect of M-atom on the formation of  $CO_3^{2-}$  prior to  $CO_2$  desorption), and (iv) examine the vibrational fingerprints of the  $CO_x$  units accompanying the formation of various surface species. Obtained results related to all these aspects are summarised in the Conclusions section.

#### 2. Models and Details of Calculations

Surface sites of the CeO<sub>2</sub> substrate were represented by a putative global-minimum structure of stoichiometric NP Ce<sub>21</sub>O<sub>42</sub> [46,47] exposing four O atoms on its top {100} nanofacet (a so-called O<sub>4</sub>-pocket [49]); see Figure 1a. M<sub>1</sub>/Ce<sub>21</sub>O<sub>42</sub> models were created via anchoring a single M<sub>1</sub> atom (M = Pd or Ag) to the O<sub>4</sub>-site of the NP; see Figure 1b,c. The removal of one O atom from the O<sub>4</sub>-pocket results in an O-deficient M<sub>1</sub>/Ce<sub>21</sub>O<sub>41</sub> model with one O vacancy in ceria (not shown in Figures; see Supplementary Material for xyz-structures). CO adsorption does not change the number of O vacancies, whereas the oxidation of CO to CO<sub>2</sub> and further transformation to CO<sub>3</sub><sup>2-</sup> require the expulsion of one or two O atoms from ceria, generating O vacancies. In the following, the number and origin of O vacancies in the NP Ce<sub>21</sub>O<sub>42</sub> are labelled as NP[*n*/*l*], where *n* is a number of O vacancies created by the transfer of O atoms from ceria to the adsorbed CO to form CO<sub>2</sub> or CO<sub>3</sub><sup>2-</sup> moieties (*l* = 0, 1, 2). For instance, the stoichiometric and O-deficient NP models M<sub>1</sub>/Ce<sub>21</sub>O<sub>42</sub> and M<sub>1</sub>/Ce<sub>21</sub>O<sub>41</sub> are referred to as M/NP[0/0] and M/NP[1/0], respectively.



**Figure 1.** Models of  $Ce_{21}O_{42}$  NP: (a) without metal atom; (b) with Pd atom anchored at {100} facet; (c) with Ag atom anchored at {100} facet. The Ce<sup>4+</sup>, Ce<sup>3+</sup>, Pd, Ag, and O atoms are shown in yellow, light blue, dark cyan, grey, and red, respectively. The atoms under the transparent plane are omitted for clarity in the Figures 2–4.

The Vienna ab initio simulation package (VASP) [52,53] was employed to determine equilibrium structures of various isomers of the  $MCO_x/NP[n/x - 1]$  complexes formed by the interaction of a CO molecule with Pd/NP[n/0] and Ag/NP[n/0] sites and transition state structures connecting selected equilibrium structures. The plane-wave basis with a 415 eV cutoff for the kinetic energy was used along with the projector-augmented wave

description of the interactions of valence electrons  $(2s^22p^4 \text{ for O}, 4s^14d^9 \text{ for Pd}, 5s^14d^{10} \text{ for Ag and } 5s^25p^66s^25d^14f^1 \text{ for Ce})$  with the atomic cores [54,55]. The NP models were separated by a vacuum space of ~1 nm in the three Cartesian directions (typical cell dimensions  $2 \times 2 \times 2 \text{ nm}^3$ ) sufficient to eliminate the interaction between periodically repeated NP images [56,57]. All calculations were performed at the  $\Gamma$ -point of the reciprocal space. A generalised-gradient corrected (GGA) exchange-correlated functional PW91 [58] was utilised with the Hubbard-type on-site corrections U [59,60] for Ce4f states providing an improved description of Ce<sup>3+</sup> ion formation in redox transitions [61,62]. The value of U = 4 eV (PW91 + U = 4 setup) were used in line with previous studies [45,48,61,63,64], though the usage of even such small U values may overestimate the formation energy of carbonates [62,65]. The minimum energy reaction paths MCO/NP[n/0]  $\rightarrow$  MCO<sub>2</sub>/NP[n/1]  $\rightarrow$  MCO<sub>3</sub>/NP[n/2] were represented with the points of the string method, and the transition states were approximated with polynomial splines [66,67].

Stabilities of the studied systems were quantified based on their formation energies E':

$$E^{f} = E(MCO_{x}/NP[0/x-1]) - E(CO) - E(M) - E(NP[0/0])$$
(1)

for the MCO<sub>x</sub>/NP models obtained from defect-free MCO/Ce<sub>21</sub>O<sub>42</sub> structure and

$$E^{f} = E(MCO_{x}/NP[1/x-1]) - E(CO) - E(M) - E(NP[0/0]) + 0.5 \times E(O_{2})$$
(2)

for the  $MCO_x/NP$  models obtained from O-deficient  $MCO/Ce_{21}O_{41}$  structure, where  $E(O_2)$  is total energy of a free  $O_2$  molecule. The binding energies  $E_b$  of CO and CO<sub>2</sub> molecules in the models under scrutiny were calculated as follows:

$$E_b(CO) = E(MCO_x/NP[n/x-1]) - E(CO) - E(M/NP[n/0]), x = 1-3; n = 0, 1$$
(3)

$$E_b(CO_2) = E(MCO_x/NP[n/x-1]) - E(CO_2) - E(M/NP[n/1]), x = 2, 3; n = 0, 1$$
(4)

The energies (3) and (4) of the same  $MCO_2/NP$  complex were used to estimate the energy of the overall oxidation process  $CO(gas) + M/NP[n/0] \rightarrow CO_2(gas) + M/NP[n/1]$ :

$$E_{ox} = E_b(CO)(MCO_2/NP[n/1]) - E_b(CO_2)(MCO_2/NP[n/1])$$
(5)

or

$$E_{ox} = E_b(CO)(MCO/NP[n/0]) + E_{CO2}^* - E_b(CO_2)(MCO_2/NP[n/1]), n = 0, 1$$
(6)

where  $E^*_{CO2}$  is the energy of CO to  $CO_2$  oxidation at a metal site; i.e., the heat of MCO/NP[n/0]  $\rightarrow$  MCO<sub>2</sub>/NP[n/1] transformation:

$$E_{CO2}^{*} = E(MCO_2/NP[n/1]) - E(MCO/NP[n/0])$$
(7)

where, the asterisk indicates that CO and CO<sub>2</sub> molecules are adsorbed.

Harmonic vibrational frequencies of  $CO_x$  groups were calculated by diagonalising the mass-weighted Hessian matrix constructed of differences of the first derivatives of total energy, obtained by displacements by  $\pm 0.015$  Å in all Cartesian directions of the M, C, and  $O_x$  atoms as well as neighbouring ceria atoms within 3.6 Å around them.

#### 3. Results and Discussion

Lowest-energy geometries and formation energies  $E^t$  of metal-free CO<sub>x</sub>/NP[n/x - 1] and metal-containing MCO<sub>x</sub>/NP[n/x - 1] complexes are shown in Figures 2–4 (for xyz-structures see Supplementary Material). Table 1 displays the parameters used to specify attachment modes (coordination) of the CO<sub>x</sub> groups. Along with the notations MCO<sub>x</sub>/NP[n/x - 1], shorter ones **MxL** and **MxLV** were used for complexes with n = 0 and n = 1, respectively, where **L** is a sequential identifier of the relative energy of a given isomer among isomers with the same **M**, **x**, and **n** (**a**—the most stable, **b**—the second most

stable, **c**—the third most stable). Ce ions with magnetic moments close to 1, at a variance to 0 for most of the cations, were qualified as  $Ce^{3+}$  ions resulting from the reduction of  $Ce^{4+}$  by electrons of MCO<sub>x</sub> moieties or O vacancies.

**Table 1.** Calculated parameters of the surface complexes shown in Figures 2–4, created by the interaction of CO with the pristine ceria  $Ce_{21}O_{42-n}$  and metal–ceria  $M/Ce_{21}O_{42-n}$  sites (M = Pd, Ag; n = 0, 1): interatomic distances—r, number of the Ce<sup>3+</sup> ions—N, difference between the numbers of spin-up and spin-down electrons—m, binding energies of CO— $E_b(CO)$  (Equation (3)) and  $CO_2-E_b(CO_2)$  (Equation (4)). Negative energy values correspond to exothermic processes.

System <sup>a</sup>	r(M-C)	r(M-O)	r(Ce-O) <sup>b</sup>	r(C-O) <sup>c</sup>	Ν	m	$E_b(CO)$	$E_b(CO_2)$
	pm	pm	pm	pm			eV	eV
Complexes w	vith CO							
Pd1a	296 <sup>d</sup>	-	-	114	2	2	-0.26	-
Pd1b	241	-	-	115	2	2	-0.13	-
Pd1c	188	-	-	116	1	0	0.13	-
Pd1aV	187	-	-	116	4	4	-1.74	-
Ag1a	201	-	-	115	1	1	-0.80	-
Ag1aV	198	-	-	115	3	3	-0.78	-
1a	293 <sup>d</sup>	-	-	114	0	0	-0.26	-
1aV	297 <sup>d</sup>	-	-	114	2	2	-0.23	-
Complexes w								
Pd2a	206	-	247	$2 \times 127$	2	0	-1.43	-1.20
Pd2b	-	-	314	$2 \times 118$	2	0	-0.41	-0.18
Pd2aV	193	237	268	129; 124	4	4	-2.33	-0.72
Pd2bV	330	-	321	$2 \times 118$	4	4	-1.77	-0.16
Ag2a	356	-	316	$2 \times 118$	3	3	-1.37	-0.16
Ag2b	234	-	268	122; 121	3	1	-1.22	-0.01
Ag2aV	207	-	249	$2 \times 132$	3	1	-1.13	-0.10
Ag2bV	342	-	348	2  imes 118	5	1	-1.10	-0.07
2a	-	-	300	2  imes 118	2	0	-1.69	-0.25
2aV	-	-	310	2  imes 118	4	2	-1.16	-0.15
2bV	-	-	254	$2 \times 125$	3	2	-0.42	0.59
Complexes w	vith $CO_3^{2-}$							
Pd3a	249	207	260	2  imes 133; 125	4	2	-1.68	-1.44
Pd3b	265	209	248	129; 135; 129	3	2	-0.99	-0.75
Pd3c	-	212	245	133; 143; 121	2	0	-0.74	-0.50
Pd3aV	275	211	258	132; 132; 128	4	2	-2.67	-1.06
Pd3bV	288	211	247	131; 143; 122	4	0	-2.28	-0.67
Ag3a	277	221	245	130; 135; 127	3	1	-2.72	-1.49
Ag3b	289	233	247	132; 141; 122	3	1	-2.09	-0.86
Ag3aV	273	216	265	130; 132; 130	5	1	-2.36	-1.40
Ag3bV	291	211	255	129; 146; 122	5	1	-2.09	-1.13
3a	-	-	237	133; 138; 122	2	0	-2.59	-1.15
3b	-	-	250	129; 132; 129	2	2	-2.28	-0.84
3aV	-	-	262	130; 131; 130	4	2	-3.23	-2.22
3bV	-	-	229	133; 137; 122	4	2	-2.53	-1.52

<sup>*a*</sup> For notations, see Figures 2–4; <sup>*b*</sup> Average distances of the O of CO<sub>x</sub> moiety and nearest neighbour Ce atom; <sup>*c*</sup> Bond lengths within CO<sub>x</sub> moiety; <sup>*d*</sup> Ce-C contact.

To specify the coordination modes of CO<sub>2</sub> in MCO<sub>2</sub>/NP[n/1] and of CO<sub>3</sub><sup>2-</sup> in MCO<sub>3</sub>/NP[n/2] complexes, additional three-digit indices were used [68] resulting in the notation that being invoked only when discussing the coordination modes of CO<sub>x</sub> species. For carbonate complexes, a notation abc (integer a, b and c range from 0 to 3) determines numbers of substrate atoms, to which each O atom in CO<sub>3</sub><sup>2-</sup> is coordinated. The middle digit b corresponds to the O atom with the highest coordination. A dot in indices a.bc or ab.c specifies that two O atoms of CO<sub>3</sub><sup>2-</sup> (corresponding to a and b or b and c, respectively) form a bidentate bond with one atom of the substrate. Dotless abc identifiers designate CO<sub>3</sub><sup>2-</sup> groups with each O atom coordinated to a different substrate atom. When the three-digits notation is used to specify coordination of CO<sub>2</sub> in MCO<sub>2</sub>/NP[n/1]

complexes, the first digit a gives the number of substrate atoms coordinated to C atom. For instance, in Figure 3, each of the atoms of  $CO_2$  in structure 111-**Pd2a** contacts different ions of the Pd/Ce<sub>21</sub>O<sub>41</sub> subsystem, and in complex 1.21-**Pd2aV**, one of the O atoms of CO<sub>2</sub> attaches to Pd and Ce atoms of Pd/Ce<sub>21</sub>O<sub>40</sub>, and C atom also to Pd.

# 3.1. Structure, Charge State, and Relative Energies of Surface Complexes with CO<sub>x</sub>3.1.1. Adsorption Sites

The particle Ce<sub>21</sub>O<sub>42</sub> exposes a {100} nanofacet (Figure 1a), which can bind metal atoms much stronger than its {111} nanofacets do [38]. We anchored single Pd and Ag atoms to the {100} nanofacet composed of four nearly coplanar two-coordinated oxygen centres forming an O<sub>4</sub>-pocket with diagonals of 459 and 443 pm. This arrangement is appropriate for accommodating transition metal cations with typical M-O bond lengths of 185–210 pm [37]. The distances between the O<sub>4</sub>-pocket atoms and the neighboring Ce atoms, 214–215 pm, are shortened versus the Ce-O distances of three-coordinated surface O atoms and four-coordinated inner O atoms, 230–250 pm.

The Ag atom binds by 2.27 eV to the  $O_4$ -site (Figure 1b), making the two diagonals of the latter almost equal. The two types of Ag-O bond lengths, 236 and 242 pm, agree with Ag-O distances of 239 and 241 pm calculated at the same theory level on a larger  $Ce_{40}O_{80}$  NP [37]. Furthermore, the present location of the Ag atom 96 pm above the  $O_4$ -plane is close to the elevation by 90 pm on  $Ce_{40}O_{80}$  [37]. A slightly larger above-plane elevation, by 102 pm, was calculated for a two-fold coordinated Ag single atom adsorbed on the Fe<sub>3</sub>O<sub>4</sub>(001) surface with energy 2.75 eV [69]. Ag atom adsorption moves O<sub>4</sub> centres upwards, elongating the involved Ce-O distances to 218–223 pm. One Ce ion reduced to +3 state, pointing to the Ag<sup>+</sup> oxidation state.

The Pd adatom is located nearly in the O<sub>4</sub> plane (Figure 1c), with an out-of-plane displacement of 18 pm, adopting a favourable planar coordination. All four Pd-O bonds are equal to 205 pm—exactly the same as for the Pd/Ce<sub>40</sub>O<sub>80</sub> model [37]. The formation of PdO<sub>4</sub> species moves O atoms closer to Pd, contracting diagonals of the O<sub>4</sub> square to 408 pm and elongating the corresponding Ce-O distances to 230 pm—the value typical for three-coordinated O centres. The high adsorption energy of the Pd atom, 4.24 eV, further increases upon interactions with the O<sub>4</sub>-site of larger ceria NPs [37]. Two Ce<sup>3+</sup> ions appeared upon Pd adsorption, indicating the oxidation state Pd<sup>2+</sup>.

An O-deficient site is created by the removal of the weakest bonded O atom (with  $E(O_V) = 1.87 \text{ eV}$ , Table 2) from the O<sub>4</sub>-pocket of the Ce<sub>21</sub>O<sub>42</sub> particle (Figure 1). Two Ce<sup>3+</sup> centres resulting from the O removal in the second Ce layer cause the elongation of the Ce<sup>3+</sup>-O bonds by ~0.1 Å. O vacancy formation in the Ag/Ce<sub>21</sub>O<sub>42</sub> model requires 2.08 eV (Table 2), nearly the same amount as for the metal-free O<sub>4</sub>-site. Thus, an extra energy cost due to breaking the Ag-O bond (in addition to two Ce-O bonds of Ce<sub>21</sub>O<sub>42</sub>) is estimated at ~0.2 eV. Two Ag-O bonds are contracted to 2.15 Å upon O removal, whereas the third bond is elongated from 2.36 to 2.49 Å. A similar bonding situation has been found in our previous work, where an Ag single atom was anchored to the bottom {100} nanofacet of Ce<sub>21</sub>O<sub>42</sub> NP [38]. Alternatively, the creation of an O-deficient Pd/Ce<sub>21</sub>O<sub>41</sub> site requires substantial energy costs of 3.08 eV (Table 2).

Model <sup><i>a</i></sup>		$CO  ightarrow CO_2$			$CO_2 \to CO_3$	
	$E(O_V)$	$E_{ox}^{b}$	$E^*_{CO2}$ <sup>c</sup>	E≠	$E^*_{CO3} d$	E≠
Ce <sub>21</sub> O <sub>42</sub>	1.87	-1.44	-1.43	1.02	-0.90	0.27
Ce <sub>21</sub> O <sub>41</sub>	2.31	-1.01	-0.92	0.82	-2.07	0.82
$Ag/Ce_{21}O_{42}$	2.08	-1.21	-0.52	0.88	-1.50	0.78
$Ag/Ce_{21}O_{41}$	2.28	-1.03	-0.35	0.53	-1.23	1.73
$Pd/Ce_{21}O_{42}$	3.08	-0.23	-1.30	0.51	-0.25	1.83
$Pd/Ce_{21}O_{41}$	1.70	-1.61	-0.59	1.06	-0.34	1.80

**Table 2.** Reaction and activation energies  $(E^{\neq})$  on the pristine ceria Ce<sub>21</sub>O<sub>42-n</sub> and metal–ceria M/Ce<sub>21</sub>O<sub>42-n</sub> models (M = Pd, Ag; n = 0, 1) depicted in Figures 2–4. Oxygen vacancy formation energies  $E(O_V)$  are also shown. Negative energy values correspond to exothermic processes. All energies are in eV.

<sup>*a*</sup> System on which initial CO adsorption takes place; <sup>*b*</sup> reaction energy of the CO oxidation  $CO(gas) + NP[n/0] \rightarrow CO_2(gas) + NP[n/1]$  or  $CO(gas) + M/NP[n/0] \rightarrow CO_2(gas) + M/NP[n/1]$ ; <sup>*c*</sup> reaction energy of the CO<sub>2</sub> formation  $CO/NP[n/0] \rightarrow CO_2/NP[n/1]$  or  $MCO/NP[n/0] \rightarrow MCO_2/NP[n/1]$ ; <sup>*d*</sup> reaction energy of  $CO_2/NP[n/1] \rightarrow CO_3/NP[n/2]$  or  $MCO_2/NP[n/1] \rightarrow MCO_3/NP[n/2]$  conversion.



**Figure 2.** Structures and formation energies  $E^{f}$  (in eV) of the complexes created by the interaction of CO with the Ce<sub>21</sub>O<sub>42-n</sub> sites (n = 0, 1): CO/Ce<sub>21</sub>O<sub>42</sub> (**1a**), CO/Ce<sub>21</sub>O<sub>41</sub> (**1aV**), CO<sub>2</sub>/Ce<sub>21</sub>O<sub>41</sub> (**2a**), CO<sub>2</sub>/Ce<sub>21</sub>O<sub>40</sub> (**2aV, 2bV**), CO<sub>3</sub>/Ce<sub>21</sub>O<sub>40</sub> (**3a, 3b**) and CO<sub>3</sub>/Ce<sub>21</sub>O<sub>39</sub> (**3aV, 3bV**). Activation energies ( $E^{\neq}$ ) of selected transition states are shown near the dashed arrows connecting corresponding initial reagent and product. The C atoms are shown in dark grey. O atoms numbered with "1", "2", and "3" enter CO, CO<sub>2</sub>, and CO<sub>3</sub> moieties, respectively. For further explanations, refer to the main text and the Figure 1 caption.

Evidently, an extra energy of ~1.2 eV is needed to cleave one Pd-O bond. Note that the Pd atom remains after O removal in a virtual square-planar environment with one coordination site empty and the lengths of the remaining three Pd-O bonds unchanged, 206–209 pm, vs. the PdO<sub>4</sub> site of Pd/Ce<sub>21</sub>O<sub>42</sub>. No reduction of Ag<sup>+</sup> and Pd<sup>2+</sup> ions occurs upon O vacancy creation, since the leaving O atom donates two electrons to two Ce<sup>4+</sup> ions, increasing the number of Ce<sup>3+</sup> centres by two (to three for Ag/NP and to four for Pd/NP).

## 3.1.2. Carbonyl Species

Let us first consider the adsorption of a CO molecule on a Ce ion of bare ceria particles (Figure 2). For both the pristine and O-deficient structures, CO adsorption energy is very small, at -0.26 and -0.23 eV, respectively (Table 1). No geometry changes are calculated upon bringing together the CO molecule and ceria species. In particular, the C-O bond is retained at 114 pm as in the free CO molecule. The C end of the CO adsorbate is 293 (1a) and 297 pm (1aV) away from the nearest Ce<sup>4+</sup> ion of NP (Table 1). CO forms angles at 176° and 172° with the Ce site. Thus, CO is rather physisorbed than chemisorbed on the Ce site of the pristine and reduced forms of the metal-free NP. The same geometry and adsorption energy of -0.26 eV (Table 1) was calculated for the Pd-containing Pd1a structure with the CO molecule bound to a Ce<sup>4+</sup> ion (Figure 3), indicating that the filling of the O<sub>4</sub>-pocket by the Pd atom mainly affects the local structure of the PdO<sub>4</sub> site.

An even smaller CO adsorption energy, -0.13 eV, was calculated for the coordinatively saturated Pd centre of the PdO<sub>4</sub> site of the unreduced model (**Pd1b**, Figure 3). Here, CO binds the Pd atom at an angle of 131° and rather long Pd–C contact of 241 pm.

CO adsorption causes minor distortions in the Pd/Ce<sub>21</sub>O<sub>42</sub> structure: Pd-O bonds extend from 205 to 207 pm, and Pd moves by 27 pm above the O<sub>4</sub> plane. The C-O bond elongates by just 1 pm vs. gas-phase CO. No CO adsorption was reported on Pd atoms in the PdO<sub>4</sub> environment saturated by two O centres of the CeO<sub>2</sub>{100} surface and two O adatoms [9]. Furthermore,  $E_b(CO) < -0.2$  eV was calculated for coordinatively saturated Pd atoms in PdO(100). Thus, a very low  $E_b(CO)$  for **Pd1b** is related to the saturation of the coordination sphere of the Pd atom in the O<sub>4</sub>-pocket of the {100} nanofacet by O atoms; the formation of an additional Pd-C bond competes with quite strong Pd-O bonds. The exceptional stability of Pd<sup>2+</sup> ions in square-planar oxygen environment in CeO<sub>2</sub> materials is also claimed in other experiments [6,27,32]. Note that CO adsorption on all Pd-containing models does not change the oxidation state Pd<sup>2+</sup>, except for **Pd1c** with an endothermic mode by 0.13 eV CO adsorption (Table 1), where a Ce<sup>3+</sup>  $\rightarrow$  Ce<sup>4+</sup> transition indicated a reduction to Pd<sup>+</sup>.

In an O-deficient structure, Pd1aV (Figure 3), the adsorbed CO occupies one of four places around Pd. This results in Pd-C bond shortening to 187 pm (Table 1), as in Pd<sub>1</sub>- $CO/CeO_2(110)$  complexes bonded with two (Pd-C = 184 pm) and three O surface atoms (Pd-C = 188 pm) [20]. The Pd-C-O angle in Pd1aV is close to  $180^{\circ}$ . The elongation of the C-O distance from 114 to 116 pm indicates noticeable d  $\rightarrow 2\pi^*$  back-donation. In **Pd1aV**, CO binds at a vacant coordination site around Pd, no bonds are broken, and adsorption induced geometry changes are minor. As a result,  $PdCO/Ce_{21}O_{41}$  is stabilised by 1.7 eV with respect to separated  $Pd/Ce_{21}O_{41}$  and CO fragments (Table 1). Interestingly, similarly strong CO binding as that for Pd1aV, 1.77 eV, and a Pd-C distance of 186 pm were calculated for the two-fold coordinated single Pd atom in  $Pd_1/Fe_3O_4(001)$  [69]. The CO bond to the isolated Pd atom is also similarly strong, at 1.8 eV [70]. The values of 1.6 eV [8] and 1.9 eV [9] were calculated for Pd<sub>1</sub>-CO/CeO<sub>2</sub>(111) and O<sub>1</sub>Pd<sub>1</sub>-CO/CeO<sub>2</sub>(100) complexes, respectively, with Pd<sub>1</sub> coordinated to three O atoms including that of isolated O<sub>1</sub>Pd<sub>1</sub> species. A CO adsorption energy of -1.5 eV was calculated for the Pd<sub>1</sub>-CO/CeO<sub>2</sub>(110) complex with the Pd<sup>+</sup> ion between two three-fold O atoms [20]. In the series of complexes Pd<sub>1</sub>- $CO/CeO_2(111)$ ,  $O_1Pd_1$ - $CO/CeO_2(111)$ , and  $O_2Pd_1$ - $CO/CeO_2(111)$ ,  $E_b(CO)$  decreases (by module) with the growth of the  $Pd_1$  coordination number from 1.6 to 0.9 and to 0.6 eV [8].



**Figure 3.** Structures and formation energies  $E^{f}$  (in eV) of the complexes created by the interaction of CO with the PdCe<sub>21</sub>O<sub>42-n</sub> sites (n = 0, 1): PdCO/Ce<sub>21</sub>O<sub>42</sub> (**Pd1a, Pd1b, Pd1c**), PdCO/Ce<sub>21</sub>O<sub>41</sub> (**Pd1aV**), PdCO<sub>2</sub>/Ce<sub>21</sub>O<sub>41</sub> (**Pd2a, Pd2b**), PdCO<sub>2</sub>/Ce<sub>21</sub>O<sub>40</sub> (**Pd2aV, Pd2bV**), PdCO<sub>3</sub>/Ce<sub>21</sub>O<sub>40</sub> (**Pd3a, Pd3b, Pd3c**) and PdCO<sub>3</sub>/Ce<sub>21</sub>O<sub>39</sub> (**Pd3aV, Pd3bV**). Activation energies ( $E^{\neq}$ ) of selected transition states are shown near the dashed arrows connecting corresponding initial reagent and product. For further explanations, refer to the main text and the captions to Figures 1 and 2.

The binding of CO to O-defect-free Ag1a and O-deficient Ag1aV complexes is moderately strong, at about 0.8 eV (Table 1), and essentially independent of the coordination- $AgO_4$  or  $AgO_3$ —of the Ag atom. This adsorption energy value fits the calculated values of -0.85 eV for CO adsorption at the two-fold coordinated Ag single atom in the  $Ag_1/Fe_3O_4(001)$  site and -0.94 eV for the two-fold coordinated Ag atom at the AgO<sub>2</sub>(111) surface well [69]. Structures of the AgCO fragments in Ag1a and Ag1aV are very similar (Figure 4). The Ag-C distance of ~200 pm (Table 1) coincides with the value computed for the  $Ag_1CO$  moiety at  $Fe_3O_4(001)$ . The C-O bond, 115 pm, is 1 pm longer than that in the free CO molecule. The Ag-C-O angle of ~170° points to slight deviation from the typically favoured linear bonding geometry. In both structures, CO adsorption triggers a further displacement of Ag atom out-of-plane of neighbouring O atoms, reaching values of 132 and 152 pm for Ag1a and Ag1aV, respectively (vs. 96 and 60 pm for CO-free structures). Ag-O bond lengths vary substantially, ranging from 237 to 258 pm for the  $AgO_4$  unit and from 224 to 264 pm for the AgO<sub>3</sub> unit. The notable distortions at the Ag/CeO<sub>2</sub> interface are not reflected in CO adsorption energies. This finding is in line with the quite weak Ag-O bonding estimated at 0.2 eV (see Section 3.1.1).

In summary, the modification of the ceria NP with single Pd and Ag atoms strongly affects its affinity to CO. Effects of Pd and Ag atoms are different. Due to strong Pd-O

bonds, the saturated  $PdO_4$  site is almost inactive towards CO adsorption, whereas the  $PdO_3$  unit with a vacant coordination place readily traps CO with a substantial energy gain. In contrast,  $AgO_4$  and  $AgO_3$  centres with weak Ag-O bonds and a more flexible geometry are more prone to adsorb CO molecules with equally moderate energies.

#### 3.1.3. Carbon Dioxide Species

For the metal-free ceria, only weakly adsorbed  $CO_2$  species were calculated: binding energies are -0.25 and -0.15 eV for **2a** and **2aV** models, respectively (Table 1). These energies are comparable with the values calculated for linearly-adsorbed  $CO_2$  at extended  $CeO_2(100)$ , (110), and (111) surfaces [71–76]. In both complexes, the linear geometry of  $CO_2$ as in the gas-phase molecule is preserved: C-O bond lengths are 118 pm and the O-C-O angle is 178°. The C atoms of  $CO_2$  molecules are located above the O vacancy parallel to the surface of NP (Figure 2); in case of **2a**, the O atom from the low-lying layer moves up and forms an O-C contact of 288 pm (vs. 390 pm in **2aV**).



**Figure 4.** Structures and formation energies  $E^{f}$  (in eV) of the complexes created by the interaction of CO with the Ag/Ce<sub>21</sub>O<sub>42-n</sub> sites (n = 0, 1): AgCO/Ce<sub>21</sub>O<sub>42</sub> (**Ag1a**), AgCO/Ce<sub>21</sub>O<sub>41</sub> (**Ag1aV**), AgCO<sub>2</sub>/Ce<sub>21</sub>O<sub>41</sub> (**Ag2a, Ag2b**), AgCO<sub>2</sub>/Ce<sub>21</sub>O<sub>40</sub> (**Ag2aV, 2bV**), AgCO<sub>3</sub>/Ce<sub>21</sub>O<sub>40</sub> (**Ag3a, Ag3b**) and AgCO<sub>3</sub>/Ce<sub>21</sub>O<sub>39</sub> (**Ag3aV, Ag3bV**). Activation energies ( $E^{\neq}$ ) of selected transition states are shown near the dashed arrows connecting corresponding initial reagent and product. For further explanations, refer to the main text and the captions to Figures 1 and 2.

The contacts of O atoms of the CO<sub>2</sub> group with Ce atoms are ~300 pm in **2a** and 304–331 pm in **2aV** (Table 1). All attempts to locate more stable structures with a bent CO<sub>2</sub> moiety have led to carbonate moieties (described in detail in Section 3.1.4). No minima corresponding to bent CO<sub>2</sub> structures of carboxylate type were found. The only located structure with distorted CO<sub>2</sub> (**2bV** at Figure 2) was found to be 0.75 eV less stable relative to the linear **2aV** model. It features C-O bond lengths of 125 pm and an O-C-O angle of 131°. The O atoms of CO<sub>2</sub> are at 252 and 257 pm from Ce atoms. The positive CO<sub>2</sub> binding energy in **2bV** of 0.6 eV (Table 1) indicates exothermic CO<sub>2</sub> release. Since only one Ce<sup>3+</sup> ion is present after CO to CO<sub>2</sub> transformation, CO<sub>2</sub> in **2bV** is negatively charged, forming a CO<sub>2</sub><sup>-</sup> anion. Thus, CO<sub>2</sub> binds weakly and easily desorbs from bare Ce<sub>21</sub>O<sub>42-x</sub> NP.

The metal-containing Pd/NP and Ag/NP models also feature structures with "linear" and "bent"  $CO_2$  (Figures 3 and 4). Similar to the bare ceria, linear  $CO_2$  is weakly bound in Pd2b, Pd2bV, Ag2a, and Ag2bV, by less than 0.2 eV (Table 1). All distances between atoms of CO<sub>2</sub> molecule and ceria are longer than 300 pm. In contrast, in "bent" CO<sub>2</sub> structures Pd2a, Pd2aV, Ag2b, and Ag2aV, CO<sub>2</sub> approaches more closely to the NP surface and forms C-M and O-Ce bonds at 193-234 and 250-270 pm, respectively. Interestingly, despite the sizable changes at the metal-adsorbate interface when going from "linear" to "bent" structures, the pairs Ag2aV/Ag2bV and Ag2a/Ag2b are isoenergetic within 0.15 eV. This observation supports our earlier finding that the formation of a new Ag-ligand bond (e.g., Ag-C with  $CO_2$ ) does not require substantial energy (limited to 0.2 eV). The bonds of  $CO_2$  in **Ag2b** are symmetrically stretched to 121 pm and the O-C-O angle is reduced to 145°, thus indicating the formation of a weakly-bound carboxylate-like complex. CO<sub>2</sub> is coordinated to Ag via the C atom at a distance of 234 pm and forms an (O)C-Ag-O angle of 175° with the Ag-O bond in trans-position. The Ag atom in  $AgCO_2/Ce_{21}O_{41}$  rises above the O<sub>3</sub>-plane by 52 pm (vs. 132 pm in AgCO/Ce<sub>21</sub>O<sub>42</sub>), similar to O-deficient adsorbate-free Ag/Ce<sub>21</sub>O<sub>41</sub> complex (60 pm). Thus, the surrounding of the Ag center in Ag2b is nearly square-planar. The formation of the  $CO_2$  moiety in Ag2aV does not lead to the entire detachment of the surface O\* atom by adsorbed CO (Figure 4). Rather, the Ag2aV model can be viewed as a complex with CO inserted into the O\*-Ag bond; as a result, the O\*-Ag distance is stretched from 227 to 287 pm. Upon this insertion, Ce-O\* bonds of the Ce-O\*-Ce bridge are elongated from 218–233 pm in Ag1aV to 249–257 pm in Ag2aV (Table 1). The shortest Ce-O\* distance in **Ag2b** is 268 pm. The CO<sub>2</sub> fragment in **Ag2aV** is more strongly distorted than in **Ag2b**: C-O bonds are 130 pm and the O-C-O angle is quite small, at 114°. The Ag-CO<sub>2</sub> distance (207 pm) is typical for metal–CO<sub>2</sub> complexes with an  $\eta^1$ -C type of CO<sub>2</sub> coordination [77]. Energies of Ag2b and Ag2aV models differ by 1.9 eV, the value close to the O vacancy formation energy at the  $O_4$ -pocket (Table 2).

The PdCO<sub>2</sub>-containing complexes, **Pd2a** and **Pd2aV**, with a "bent" CO<sub>2</sub> moiety (Figure 3) are characterised by substantial CO<sub>2</sub> binding energies of -1.2 and -0.7 eV, respectively (Table 1). Both models exhibit a carboxylate-like structure of CO<sub>2</sub>: C-O bonds are 124–130 pm long and the O-C-O angle is  $130 \pm 3^{\circ}$ . In the **Pd2a** structure, CO<sub>2</sub> is  $\eta^1$ -coordinated to the Pd atom with the Pd-C distance, 206 pm, comparable to three Pd-O bonds around the metal center, 208–220 pm. In the **Pd2aV** complex, the CO<sub>2</sub> molecule is  $\eta^2$ -coordinated to Pd via the short Pd-C bond, 193 pm, and long Pd-O contact of 237 pm; two other Pd-O bond lengths are 205–220 pm. The Pd atom in both **Pd2a** and **Pd2aV** complexes is nearly in a square-planar environment with Pd shifting from the ligand plane by 25–30 pm. Interestingly, no ceria-supported PdO<sub>3</sub> structures with calculated CO<sub>2</sub> binding energies more than -0.4 eV (by module) were found in the literature [8,9,20]. The larger values -0.82 eV [20] and -0.96 eV [8] were calculated for PdO<sub>2</sub> units at CeO<sub>2</sub>(110) and CeO<sub>2</sub>(111) surfaces. In part, this can be explained by a favourable square-planar geometry within the O<sub>3</sub>Pd-CO<sub>2</sub> fragment formed at the {100} nanofacet of Ce<sub>21</sub>O<sub>42</sub> NP. Indeed, such structures are hindered at extended CeO<sub>2</sub> surfaces by geometrical constrains.

Thus, similar to the CO case, the strongest bonding is calculated for the  $PdO_3$  site with one vacant coordination place. This is followed by the  $PdO_2$  unit with two vacant valences, which binds  $CO_2$  in a side-on fashion. Both structures are characterised by  $CO_2$  bending.

The adsorption of  $CO_2$  in linear mode results in an energy gain less than 0.2 eV for both Pd and Ag derivatives. In contrast to the Pd derivatives, both "linear" and "bent" modes have similar small adsorption energies for the Ag systems. Thus, AgO<sub>3</sub> and AgO<sub>2</sub> sites again do not show differences in adsorption properties, whereas their Pd analogs do.

#### 3.1.4. Carbonate Species

Carbonate species are formed upon the coordination of the  $CO_2$  molecule via the C atom to an O center of ceria NP. Thus,  $CO_3^{2-}$  formation can be considered as a form of  $CO_2$  adsorption, and  $CO_2$  binding energy can be applied to the estimation of the stability of these carbonate species.

The carbonate-like  $CO_2$  adsorption at bare ceria NPs with binding energies of 0.84–2.22 eV (Table 1) is the most exothermic of all types of  $CO_2$  coordination discussed above. The strongest binding of 2.22 eV is calculated for the tridentate 2.2.2-structure (3aV, Figure 2), with each O atom of the  $CO_3$  moiety coordinated to two Ce ions. The flat-lying carbonate fragment of **3aV** fills two O vacancies with its oxygen atoms; i.e., all O atoms of the CO<sub>3</sub> moiety are in neighbouring O-positions of the  $O_4$ -pocket. The  $E_b(CO_2)$  energy for **3aV** is between -1.9 and -2.3 eV, respectively, calculated for its formal 2.2.2-CO<sub>3</sub> analog on regular  $CeO_2(100)$  [78] and bidentate 220-CO<sub>3</sub> on O-deficient CeO<sub>2</sub>(110) [74]. The 2.2.2model features three C-O bonds equally stretched to 130 pm and O-C-O angles of  $118 \pm 2^{\circ}$ . Three of four  $Ce^{3+}$  ions are located in the  $Ce_4$  square just below the  $O_4$ -pocket of the {100} nanofacet. The other tridentate 1.21-structure 3b is stabilised by only 0.84 eV relative to the sum  $CO_2$  + NP (Table 1). Here, each of the two O atoms of  $CO_2$  has only one contact with Ce ions of 238 and 247 pm. As in the 3aV model, the CO<sub>3</sub> fragment in 3b has a slightly distorted  $C_{3V}$  symmetry, with C-O bonds of 130  $\pm$  2 pm and O-C-O angles of  $117-122^{\circ}$ . Its CO<sub>3</sub> plane is tilted by  $40^{\circ}$  relative to the Ce<sub>4</sub> layer. The CO<sub>2</sub> binding energy in the 1.21-carbonate is comparable to value -0.72 eV calculated for the flat-lying 2.21-CO<sub>3</sub> structure [71] formed upon CO adsorption at the stoichiometric  $CeO_2(100)$  surface with every second atom of the O layer removed. The deletion of an O atom from 1.21-3b to give 2.2.2-3aV costs only 0.67 eV. Thus, energy for the removal of a lattice O atom (~2.3 eV for formation of a 2nd O vacancy) is compensated by the formation of two new O-Ce contacts in 3aV with ~1.6 eV energy gain. At variance, the energy difference between two 1.20structures  $3a \rightarrow 3bV$ , with a bidentate coordination, 1.68 eV, does not decrease with respect to  $E(O_v)$  for  $Ce_{21}O_{42} \rightarrow Ce_{21}O_{41}$  transition, since both models have a similar bonding pattern of carbonate species. Thus, not surprisingly, the  $CO_2$  binding energies for **3a** and **3bV** models are not very different, at -1.15 and -1.52 eV, respectively (Table 1). In both **3a** and **3bV** models, the CO<sub>3</sub> moiety is tied by three O-Ce bonds at 227, 238, and 245 pm. The three O-C bond lengths are different and increase from 122 to 133 and to 137 pm with the growing coordination number of O atom from 0 to 1 and to 2, respectively. Notably, at the O-deficient ceria, the flat-lying structure 3aV is preferred over the standing one 3bV by 0.7 eV. A similar difference of 0.84 eV was calculated for CO<sub>3</sub> moieties oriented parallel and perpendicular to the O-defective  $CeO_2$  (100) facet [71]. The trend of destabilising carbonate species at ceria substrates upon surface enrichment by O atoms [71] is also supported by the present data.

Three types of carbonate structures—standing (perpendicular), tilted and flat-lying (parallel)—were also located for metal-containing ceria NPs (Figures 3 and 4). The "standing"-type carbonate species is coordinated in a bidentate way in 1.30-models Pd3c, Pd3bV, Ag3b, and Ag3bV. These structures are very similar to 1.20-carbonates 3a and 3bV at bare ceria NPs, with the difference that the CO<sub>3</sub> moiety is additionally bound with the M center via one short M-O bond of 211 pm (233 pm in Ag3b). This extra metal–oxygen bond induces the elongation of other three O-Ce contacts by 5–15 pm as well as elongates  $O \cdots CO_2$  contact by 4–9 pm, which probably weakens  $CO_2$  binding from -1.15-1.52 eV in metal-free models to -0.86-1.13 and -0.39-0.50 eV in Ag- and Pd-carbonates, respectively (Table 1). Thus, for Pd-systems, carbonate-like CO<sub>2</sub> adsorption in a "standing" mode is

weaker than adsorption in carboxylate PdCO<sub>2</sub> form. Notably, Pd is in a zero-oxidation state in both types of complexes.

Stronger CO<sub>2</sub> bindings, of -1.44 and -1.40 eV, were calculated for the "flat-lying" structures 2.2.1-Pd3a and 2.2.2.-Ag3aV, respectively (Table 1). The CO<sub>3</sub> fragment is bidentately coordinated with two M-O contacts: almost equal in Pd3a (207 and 211 pm) and asymmetric in Ag3aV (216 and 265 pm). In the Pd3a structure, all O atoms of the CO<sub>3</sub> moiety leave their lattice positions to form short Pd-O bonds (Figure 3). This leads to the creation of a  $CO_3^{2-}$  unit with a formal charge of -2 and oxidation of Pd to the +2 oxidation state. Even shorter M-O bonds are formed between metal centres and O atoms of the ceria support: 201–202 pm for Pd-O and 206 pm for Ag-O bonds. Thus, the Pd atom is four-coordinated whereas Ag is three-coordinated by O ligands. Importantly, the PdO<sub>4</sub> unit in **Pd3a** is in a slightly distorted stable square-planar configuration. Likely, this contributes greatly to making Pd3a the most energetically favourable among all studied ceria-supported  $PdCO_x$  species. The **Pd3a** complex is even 0.25 eV more stable than another square-planar structure: carboxylate Pd2a complex with O<sub>3</sub>PdCO<sub>2</sub> unit. The plane of the CO<sub>3</sub> subsystem forms a small angle of about 20° with the O<sub>3</sub>M plane. The "lying" Ag3aV structure is only 0.27 eV stabilised with respect to "standing" Ag3bV, whereas the analogous stabilisation for  $Pd3a \rightarrow Pd3c$  transition reaches 0.87 eV.

The Ag3a 1.21-complex with the "tilted" coordination mode of  $CO_3$  is the most favourable Ag-carbonate structure with a  $CO_2$  adsorption energy of -1.49 eV (Table 1). Similar to the 1.30-model **Ag3b**, the CO<sub>2</sub> moiety in the **Ag3a** structure is tied in an  $\eta$ -C,O fashion with ceria support by means of two bonds. While the distance O-C-O $\cdots$ Ce is the same in Ag3a and Ag3b, the  $O \cdots CO_2$  contact in Ag3a is shortened to 135 pm from 141–146 pm in Ag3b. The shorter Ag-O contact of 221 pm (vs. 233 pm in Ag3b) is formed with the O atom, which has no bonds with CeO<sub>2</sub>; this is comparable to the Ag-O bond length of 218 pm with ceria. Remarkably, "tilted" 1.21-complex Ag3a is by 0.63 eV stabilised with respect to "standing" 1.30-model **Ag3b** with an  $E(CO_2)$  of -0.86 eV and reaches a CO<sub>2</sub> adsorption energy of -1.52 eV of the metal-free "standing" 3bV model. 1.21-Ag3a complex has a nearly identical CO<sub>2</sub> adsorption energy (within 0.1 eV) to the "lying" 2.2.2-Ag3aV model. In both complexes, the cationic Ag<sup>+</sup> center is three-coordinated. Despite the closer contact of the CO<sub>3</sub> unit with ceria support in the 2.2.2-model, its formation from the Ag3a structure by the removal of O bound to the Ag atom requires a substantial energy of 2.18 eV. In contrast, O deletion from 1.21-Pd3b to give 2.2.1-Pd3aV is endothermic by only 1.55 eV. Both "tilted" complexes, Pd3b and Pd3aV, have  $CO_2$  adsorption energies of -0.75 and -1.06 eV, respectively, bracketing the value for the "tilted" metal-free **3b** complex. Note that the  $Pd3b \rightarrow Pd3aV$  transition is associated with the reduction of the Pd centre from  $Pd^+$  to  $Pd^0$ .

In summary, carbonate  $CO_3^{2-}$  species show coordination patterns different from CO and  $CO_2$  moieties: they tie to the ceria support or metal centres via O atoms, whereas the C atom does not participate in adsorbate–substrate interaction. From the comparison of  $CO_2$  binding energies at the metal-free and M-containing sites, it follows that  $CO_2$  binds slightly more strongly in  $MCO_3/NP[0/2]$  structures than in  $CO_3/NP[0/2]$ . Conversely, the creation of an extra O vacancy stronger stabilises the formation of the carbonates at metal-free NPs than that at metal-containing NPs.

#### 3.2. Reaction Energies

In this section, we consider the energies of CO to CO<sub>2</sub> oxidation,  $E^*_{ox}$ , and of CO<sub>2</sub> to CO<sub>3</sub> transformation,  $E^*_{CO3}$ , along with the corresponding activation barriers,  $E^{\neq}$ , calculated for Pd/NP and Ag/NP models in comparison with the bare NP model. Data in Table 2 show that i) all these reactions are exothermic and ii) activation barriers for the carbonate formation are often higher than for the oxidation of CO.

Bare NPs adsorb CO very weakly, by less than 0.3 eV (Table 1, Figure 2). The extraction of lattice O and the formation of ceria-supported CO<sub>2</sub> yield energies of 0.9-1.4 eV and require 0.8-1.0 eV of activation (Table 2). Because of the low CO<sub>2</sub> desorption barriers,

0.15–0.25 eV (Table 1,  $E_{des}(CO_2) = -E_b(CO_2)$  for **2a** and **2aV**), the formed CO<sub>2</sub>/Ce<sub>21</sub>O<sub>42</sub> and CO<sub>2</sub>/Ce<sub>21</sub>O<sub>41</sub> complexes can quite easily decompose. Otherwise, they are expected to exothermically transform into carbonates with activation barriers ranging from low— 0.27 eV (for CO<sub>2</sub>/NP[0/1])—to moderate—0.82 eV (for CO<sub>2</sub>/NP[1/1])—values. The thermodynamic stability of such surface carbonate complexes makes them most probable candidates for experimental detection [68].

CO binds very weakly at the defect-free  $PdO_4$  site of the Pd/NP[0/0] complex with an adsorption energy of -0.1 eV (Pd1b in Table 1). CO<sub>2</sub> can form via the interaction of the adsorbed CO with a lattice O atom. The process, which is exothermic by 1.3 eV, requires the overcoming of a barrier of 0.5 eV (Table 2). The formed CO<sub>2</sub> molecule is quite strongly bound, with a desorption energy of 1.2 eV (Table 1,  $E_{des}(CO_2) = -E_b(CO_2)$  for Pd2a). Overcoming an even higher barrier of 1.8 eV (Table 2) is needed to extract one more lattice O centre and activate the transformation of  $CO_2$  to  $CO_3^{2-}$ , which is exothermic by only 0.25 eV. CO is strongly, by 1.7 eV, adsorbed at the O-deficient  $PdO_3$  site (Pd1aV in Table 1). The formed stable  $O_3PdCO$  species can transform to  $O_2PdCO_2$  with an activation barrier of 1.1 eV and moderate reaction exothermicity of 0.6 eV (Table 2). Slightly exothermic by 0.3 eV, the formation of carbonate is hindered by a high energy barrier of 1.8 eV. A much lower energy barrier of 0.7 eV ( $-E_b(CO_2)$  for **Pd2aV** in Table 1) is required to desorb the CO<sub>2</sub> molecule into the gas phase. Thus, the most likely ceria-supported Pd-intermediates to be detected in reaction medium are saturated square-planar O<sub>3</sub>PdCO/NP (Pd1aV) and O<sub>3</sub>PdCO<sub>2</sub>/NP (Pd2a) complexes whose formation proceeds with a sizable energy release (1.3–1.7 eV) and moderate activation barriers of 0.5 eV and whose decomposition is hindered by substantial barriers of about 1.1-1.2 eV.

The reactivity of Ag-containing systems is different (Table 2). CO adsorption on defectfree and O-deficient Ag/NP systems occurs with a moderate energy gain of 0.8 eV (Table 1). The formed O<sub>4</sub>AgCO and O<sub>3</sub>AgCO species are converted into the corresponding O<sub>3</sub>AgCO<sub>2</sub> and O<sub>2</sub>AgCO<sub>2</sub> species with a similar exothermicity around 0.4 eV, but the activation barrier for the more O-saturated complex O<sub>4</sub>AgCO, 0.9 eV, is 0.4 eV higher than that for O<sub>3</sub>AgCO (Table 2). Despite the high exothermicity of 1.2–1.5 eV, the transformation to carbonates is hindered by barriers of ~0.8–1.7 eV. Alternatively, the decomposition with CO<sub>2</sub> desorption should proceed quite readily (see  $-E_b(CO_2)$  values for **Ag2b** and **Ag2aV** in Table 1). Thus, the carbonyl complexes O<sub>4</sub>AgCO/NP (**Ag1a**) and O<sub>3</sub>AgCO (**Ag1aV**), which are easily formed with notable energy gains, are expected to be detectable in a reaction medium. The detection of carboxylate AgCO<sub>2</sub>/NP (**Ag2b** and **Ag2aV**) complexes seems problematic due to their instability with respect to CO<sub>2</sub> desorption.

We estimated the propensity of CO to CO<sub>2</sub> transformation by the energy of the CO(gas) +  $M/NP[n/0] \rightarrow CO_2(gas) + M/NP[n/1]$  oxidation reaction,  $E_{ox}$  (Equations (5) and (6)). It is directly connected with the ease of O release from the ceria lattice and O vacancy creation (Table 2).

Let us compare the oxidation reaction energy  $E_{ox}$  calculated for the O<sub>4</sub>-site of CeO<sub>2</sub> NP with the calculated energies for clean ceria surfaces [65,71,74,79,80] and the Pd<sub>1</sub>–ceria interfaces [8,9,20]. This reaction was found to be slightly, by 0.4–0.6 eV, exothermic on bare CeO<sub>2</sub>(111) [74,79] and notably more exothermic on CeO<sub>2</sub>(110) [65,74,78,79], at 1.1–1.8 eV. The calculated reaction exothermicity further drastically increases to 3.1 eV for the CeO<sub>2</sub>(100) surface with the most exposed O atoms [71]. Our calculated  $E_{ox}$  energies for the stoichiometric and O-deficient NP{100} sites, -1.44 and -1.01 eV, respectively (Table 2), are considerably lower than those for the CeO<sub>2</sub>(100) surface, but in a similar range to that for CeO<sub>2</sub>(110) (assuming that the  $E_{ox}$  value should increase by ca. 0.5 eV when the U value is increased from 4 to 5 eV [62]).

At the  $Pd_1/CeO_2(100)$  and  $O_1Pd_1/CeO_2(100)$  sites, the conversion of CO to  $CO_2$  was characterised by energy yields of 0.6 and 1.2 eV [9], respectively—markedly lower than at the pristine  $CeO_2(100)$  surface. The reaction is also moderately exothermic at the  $Pd_1/CeO_2(110)$  interface, by 1.2 eV [20], and highly exothermic at the isolated  $O_1Pd_1$  and  $O_2Pd_1$  species on the  $CeO_2(111)$  surface, by 2.9 and 2.7 eV [8]—much more exothermic than

at the regular CeO<sub>2</sub>(111) surface because of exposing weakly bonded O atoms. Compared to the above-mentioned transformations, that at Pd/NP[0/0] has the lowest  $E_{ox}$  of -0.23 eV. There, the formal migration of an O-atom of the O<sub>4</sub>Pd moiety to become a part of CO<sub>2</sub> molecule is difficult even in comparison with the three-coordinated O centre of the regular CeO<sub>2</sub>(111) surface. For the conversion at the Pd/NP[1/0] site with  $E_{ox} = -1.61$  eV, the relocation of the second O atom of the O<sub>4</sub>Pd moiety is more favourable than of the O atom of the isolated O<sub>1</sub>Pd<sub>1</sub> moiety on CeO<sub>2</sub>(100) or Pd<sub>1</sub> at CeO<sub>2</sub>(110), but less beneficial than that of O<sub>1</sub>Pd<sub>1</sub> or O<sub>2</sub>Pd<sub>1</sub> ad-species on CeO<sub>2</sub>(111). For the reactions on Ag/NP[n/0] sites, the energy  $E_{ox}$ , -1.0-1.2 eV, is slightly lower than on the Pd/NP[1/0] site, thus approaching  $E_{ox}$  values for O<sub>1</sub>Pd<sub>1</sub>/CeO<sub>2</sub>(100) and Pd<sub>1</sub>/CeO<sub>2</sub>(110) systems. Therefore, our model ceria particle with the Pd atom adsorbed on the O-defective {100} nanofacet appears to be more reactive in CO to CO<sub>2</sub> oxidation than its formal analogue of the Pd-doped extended CeO<sub>2</sub>(100) surface.

Thus, the trend for lowering  $E_{ox}$  in bare and M-containing ceria systems is Pd/NP[1/0] > NP[n/0]  $\approx$  Ag/NP[n/0] > Pd/NP[0/0], which correlates with the growth of the O vacancy formation energy in the same row (Table 2). Note that the reactivity in CO oxidation of both Ag/NP complexes—with and without an O vacancy near the Ag atom is similar. Conversely, the presence of an O vacancy in the vicinity of Pd atom is mandatory for CO oxidation at Pd/NP systems to proceed. This makes two consequent steps of CO oxidation at the Pd/Ce<sub>21</sub>O<sub>42</sub> nanoparticle problematic. Such characteristics of the studied models as moderately strong CO adsorption, exothermic overall CO oxidation process, sufficiently low barriers of MCO to MCO<sub>2</sub> transformations, and ease of CO<sub>2</sub> desorption render CO oxidation by lattice ceria oxygen atoms more favourable at the sites with Ag than with Pd. Comparing the reaction and activation energies of CO to  $CO_2$  and  $CO_2$  to  $CO_3$  conversions for M-containing ceria NPs, we conclude that the most probable species to be observed experimentally are AgCO and PdCO carbonyls and carboxylate PdCO<sub>2</sub> species. Unlike purely ceria nanoparticles, the formation of silver and palladium carbonates is prohibited by high activation barriers. We note, however, that for precise information on the species present in the reaction medium at equilibrium, a microkinetic modelling is required, which is out of the scope of the present study.

#### 3.3. CO<sub>x</sub> Vibrational Fingerprints

Several studies reported measured [7,9,16,25,26,28,68,81] and calculated [8,9,68,82–84] vibrational frequencies of CO, CO<sub>2</sub>, and CO<sub>3</sub><sup>2–</sup> groups in the complexes formed by the interaction of CO or CO<sub>2</sub> with cerium-based substrates. The calculated frequencies for selected stretching vibrations of CO<sub>x</sub> subsystems of the structures displayed in Figures 2–4 are collected in Table 3.

Our modelling revealed that the CO stretching frequency, v(CO), for the molecule attached to a cerium ion in M-free systems 1a and 1aV and Pd-containing model Pd1a shifts by 27-34 cm<sup>-1</sup> to the short-wave region, which agrees with the measured blue shifts of 27–32 cm<sup>-1</sup> (vs. 2143 cm<sup>-1</sup> for free molecule [85]) for CO interacting with the Ce<sup>4+</sup> centres of the nanostructured  $CeO_2$  [19,86]. Note that the quantitatively precise reproduction of measured vibrational frequencies of CO on ceria requires going beyond the U-corrected generalised-gradient exchange-correlation functionals to hybrid-type functionals [81]. In contrast, v(CO) for the fragments with M<sub>1</sub>-CO bonding formed on M/NP[n/0] substrates shows redshifts of 50–113 cm<sup>-1</sup>, consistent with the C-O bond elongation by 1–2 pm. Among the models containing Pd<sub>1</sub> species, the Pd1b complex with Pd<sup>2+</sup> cation coordinated by four two-coordinated O anions reveals a medium redshift of 84 cm<sup>-1</sup>, and the maximum redshifts of 109–113 cm<sup>-1</sup> are identified for the other two complexes, Pd1c and Pd1aV, with three-fold coordinated Pd<sup>+</sup> and Pd<sup>2+</sup> cations (Table 3). For the earlier examined systems with Pd-CO bonds, the redshifts  $\Delta v$ (CO) were calculated to increase from 6 to 96 cm<sup>-1</sup> in the order  $O_2Pd_1/CeO_2(111) < O_1Pd_1/CeO_2(100) \approx O_1Pd_1/CeO_2(111) < Pd_1/CeO_2(100)$ < Pd<sub>1</sub>/CeO<sub>2</sub>(111) [8,9]. These values are comparable with those attributed to the CO molecule contacting one Pd atom in experimental studies of Pd/CeO<sub>2</sub>, 13–123 cm<sup>-1</sup> [7,28]

and 10–69 cm<sup>-1</sup> [9,16,25] and PdO/CeO<sub>2</sub>, 58 cm<sup>-1</sup> [26]. Note that the  $\Delta\nu$ (CO) redshifts for CO coordinated to the supported Pd cations are opposite to the blue shift for the free PdCO<sup>+</sup> ion, measured at 63 cm<sup>-1</sup> and calculated at 75 cm<sup>-1</sup> (B3LYP), but approach a redshift for the neutral PdCO molecule, measured at 87 cm<sup>-1</sup> and calculated at 99 cm<sup>-1</sup> [70]. The redshifts of  $\nu$ (CO) for **Ag1a** and **Ag1aV** complexes, at 50 and 61 cm<sup>-1</sup>, respectively, are smaller than those for the Pd<sub>1</sub>-CO molecule, 144 cm<sup>-1</sup> [87]. Similarly large frequency redshifts to those obtained for CO on single Ag and Pd cations on a ceria NP were calculated for CO adsorbed on Pt<sup>+</sup> and Pd<sup>2+</sup> cations anchored to ceria [83].

**Table 3.** Calculated vibrational frequencies  $\nu(CO)$ ,  $\nu(CO_2)$  and  $\nu(CO_3)$  for the systems depicted in Figures 2–4 along with the corresponding frequency shifts  $\Delta\nu(CO)$  and  $\Delta\nu(CO_2)$  with respect to calculated vibrational frequencies of free molecules CO ( $\nu$ (free CO) = 2131 cm<sup>-1</sup>) and CO<sub>2</sub> ( $\nu_{asym}$ (free CO<sub>2</sub>) = 2363 cm<sup>-1</sup>).

S	ν(CO)	Δν(CO)	ν(CO <sub>2</sub> )	$\Delta \nu (CO_2)$	ν(CO <sub>3</sub> )
System <sup>a</sup>	$\mathrm{cm}^{-1}$	$\mathrm{cm}^{-1}$	cm <sup>-1</sup>	cm <sup>-1</sup>	cm <sup>-1</sup>
Complexes with CO					
Pd1a	2158	27	-	-	-
Pd1b	2047	-84	-	-	-
Pd1c	2022	-109	-	-	-
Pd1aV	2018	-113	-	-	-
Ag1a	2081	-50	-	-	-
Ag1aV	2070	-61	-	-	-
1a	2162	31	-	-	-
1aV	2165	34	-	-	-
Complexes with CO <sub>2</sub>					
Linear O-C-O					
Pd2b	-	-	2323	-40	_
Pd2bV	-	-	2332	-31	-
Ag2a	-	-	2348	-15	-
Ag2bV	-	-	2350	-13	-
2a	-	-	2354	-9	_
2aV	-	-	2360	-3	_
Bent O-C-O			2000	0	
Pd2a	-	-	1535	-828	_
Pd2aV	-	-	1628	-735	_
Ag2aV	-	-	1332	-1031	_
Ag2b			1950	-413	
2bV			1657	-706	
Complexes with $CO_3^{2-}$	-	-	1007	-700	-
Tridentate					
Pd3a					1573; 1221
Pd3b	-	-	-	-	1451; 1243
Pd3aV	-	-	-	-	1442; 1309
	-	-	-	-	
Ag3a	-	-	-	-	1511; 1233
Ag3aV	-	-	-	-	1434; 1337
3b	-	-	-	-	1442; 1280
3aV	-	-	-	-	1400; 1353
Bidentate					1744 1110
Pd3c	-	-	-	-	1746; 1118
Pd3bV	-	-	-	-	1729; 1122
Ag3b	-	-	-	-	1689; 1126
Ag3bV	-	-	-	-	1752; 1169
3a	-	-	-	-	1714; 1082
3bV	-	-	-	-	1704; 1101

<sup>a</sup> For system designations see Figures 2–4.

The highest frequencies of CO<sub>2</sub> stretching vibrations in 2a, 2aV, Ag2a, Ag2bV, Pd2b, and Pd2bV complexes with a linear CO<sub>2</sub> moiety are close to that of the IR active CO<sub>2</sub> asymmetric stretching frequency of the free molecule; the negative shift  $\Delta v(CO_2)$  does not exceed 40 cm<sup>-1</sup> (Table 3). The redshift  $\Delta v$ (CO<sub>2</sub>) for a bent CO<sub>2</sub> fragment in **2bV**, 706  $cm^{-1}$ , is comparable to the experimental value for the free anion  $CO_2^{-1}$ , 691  $cm^{-1}$  [88] (vs. measured  $\nu(CO_2)$  of free molecule 2349 cm<sup>-1</sup> [76]) and CO<sub>2</sub><sup>-</sup> species at TiO<sub>2</sub> surface, 709 cm<sup>-1</sup> [89]. The redshift  $\Delta v(CO_2)$  for the MCO<sub>2</sub>/NP[*n*/1] systems having the M-C bond from 413 to 1031 cm<sup>-1</sup> is associated with the bending of the CO<sub>2</sub> moiety and C-O bond elongation by 3 to 14 pm. The redshifts for  $PdCO_2/NP$  systems, 828 cm<sup>-1</sup> (Pd2a) and 735 cm<sup>-1</sup> (Pd2aV), are in the range of values 610–850 cm<sup>-1</sup> reported for coordination compounds with CO<sub>2</sub> attached to a single d-metal atom [76], while redshift  $\Delta v$ (CO<sub>2</sub>) in Ag2b is smaller, at only 413 cm<sup>-1</sup>.  $\Delta v(CO_2)$  for the CO<sub>2</sub><sup>2-</sup> group in Ag2aV with CO bonds elongated by 14 pm compared to those of free CO<sub>2</sub> matches the measured redshifts of 1020–1079 cm<sup>-1</sup> for carbonite ions at CeO<sub>2</sub> [90] and in Cs<sub>2</sub>CO<sub>2</sub> [91]. Thus, the increased redshift  $\Delta v$ (CO<sub>2</sub>) seems to correlate with C-O bond elongation, **Ag2b** < **Pd2aV** < **Pd2a** < Ag2aV (Table 1).

Differences in the length of the intramolecular C-O bonds are considered among the main factors determining the frequency splitting of  $CO_3^{2-}$  stretching vibrations [68]. In tridentate  $CO_3/NP[n/2]$  and  $MCO_3/NP[n/2]$  complexes, the C-O bond lengths are 125–135 pm for the 1.11- and 1.21-modes, and in less symmetrical bidentate 1.20- and 1.30-structures, the C-O bonds cover a wider interval 121–146 pm.

Overall, the splitting of C-O frequencies is lower for more symmetrical species. In particular, for the two highest frequencies ( $v_1$  and  $v_2$ ), it is only 47 cm<sup>-1</sup> for 2.2.2-structures, increasing to 97–133 cm<sup>-1</sup> for 2.2.1-isomers and 162–252 cm<sup>-1</sup> for 1.2.1- and 1.1.1-isomers (Table 3). For 1.20- and 1.30-carbonates, the splitting becomes as high as 563-632 cm<sup>-1</sup>. According to earlier calculations, carbonate groups attached to nanostructured ceria surface by three oxygen atoms, represented by sets of 1.21-, 1.2.1- and 1.3.1-structures, are characterised by  $v_1$  values from 1590 to 1490 cm<sup>-1</sup> [68]. With the inclusion of  $v_1$  values for tridentate structures **3b** and **3aV**, this range extends to 1400 cm<sup>-1</sup>. The calculated  $v_1$ for tridentate carbonates corresponds to a broad experimental region of  $1620-1450 \text{ cm}^{-1}$ attributed to the high-frequency vibrations of the carbonate groups formed on surfaces with oxygen vacancies, as well as on facet, edge, and corner sites of ceria particles [68]. The v1 frequencies for tridentate carbonate Ag3a, Pd3a, Pd3b, Ag3aV, and Pd3aV complexes, of 1573–1434 cm<sup>-1</sup>, fall between the limits of the clean surface of the nanostructured ceria. The  $v_1$  values for 1.20-carbonate groups in **3a** and **3bV** systems are in a narrow range of 1718–1698 cm<sup>-1</sup> for bidentate carbonates [68] corresponding to the measured frequency interval of 1732–1722 cm<sup>-1</sup> [68]. The  $v_1$  frequencies of the M-containing 1.30-complexes Pd3c, Pd3bV, Ag3b, and Ag3bV are in a broader range of 1752–1689 cm<sup>-1</sup>, which includes the interval for bidentate  $CO_3^{2-}$  on M-free ceria substrates.

The calculated  $v_2$  frequency range of the complexes on the clean ceria surface is 1353–1227 cm<sup>-1</sup>. This corresponds to the experimental frequencies of 1380 and 1280 cm<sup>-1</sup> [68]. The  $v_2$  values for tridentate 1.21- and 2.2.1-carbonates **P3b**, **Pd3aV**, **Ag3a**, and **Ag3aV** for M-containing systems, at 1340–1220 cm<sup>-1</sup>, are between these limits, and that for the 1.11-**Pd3a** isomer is only 6 cm<sup>-1</sup> below the low-end threshold. The  $v_2$  frequencies are distributed between 1194 and 1082 cm<sup>-1</sup> for bidentate carbonates depicted in Figure 2 and those examined in [68]; the related experimental values are 1147–1133 cm<sup>-1</sup> [68].

The range of  $v_2$  for the CO<sub>3</sub> moieties coordinated in a bidentate way at M/NP systems is 1170–1120 cm<sup>-1</sup>. Thus, the frequency ranges are similar for the metal-free CO<sub>3</sub>/NP and MCO<sub>3</sub>/NP sites, making the discrimination between the metal-containing and bare ceria particles solely on the basis of the vibrational spectroscopy data problematic.

In summary, our calculations show that the C-O stretching vibrations of the ceriasupported PdCO and AgCO fragments feature redshifts up to ~110 cm<sup>-1</sup>, which is at variance with the blue shift at metal-free ceria. Redshifts of CO<sub>2</sub> asymmetric stretching frequency of the M-CO<sub>2</sub> fragments are much higher, up to ~830 cm<sup>-1</sup> for carboxylate MCO<sub>2</sub> and further increasing by ~200 cm<sup>-1</sup> for carbonite AgCO<sub>2</sub>. The two highest  $\nu$ (CO<sub>3</sub>) stretching frequencies of M-CO<sub>3</sub> structures lie in intervals 1755–1690 cm<sup>-1</sup> and 1170–1120 cm<sup>-1</sup> for the CO<sub>3</sub> moiety coordinated in a bidentate fashion and 1575–1430 cm<sup>-1</sup> and 1340–1220 cm<sup>-1</sup> for the CO<sub>3</sub> in the tridentate coordination.

#### 4. Conclusions

 $CO_x$  intermediates formed upon CO adsorption and oxidation on single M = Pd and Ag atoms coordinated to the O<sub>4</sub>-site on the {100} facet of a Ce<sub>21</sub>O<sub>42</sub> nanoparticle have been studied computationally. Equilibrium structures, CO<sub>x</sub> vibrational frequencies, and energetic parameters of various MCO<sub>x</sub>-containing complexes have been determined. The influence of the creation of an O vacancy nearby the M atom has been also investigated. The stability of the CO<sub>x</sub> moieties anchored to the ceria-supported M atom is found to increase in the order MCO < MCO<sub>2</sub> < MCO<sub>3</sub>, similar to the trend for CO<sub>x</sub> species adsorbed on M-free ceria NP.

Except for the Pd atom saturated by four O atoms of the ceria surface O<sub>4</sub>-site, which is unable to properly adsorb CO, the doping of the ceria nanoparticle with Pd and Ag atom increases its propensity to bind the CO molecule with respect to bare ceria material. In particular, the CO adsorption energy value reaches -1.7 eV for a PdCO unit on a ceria nanoparticle with a nearby O vacancy. CO binding in AgCO complexes, regardless of the presence or absence of a nearby O vacancy, is moderately strong, at -0.8 eV. All these species are the most probable candidates to be detected experimentally, also due to the presence of moderate barriers for CO oxidation (0.5–1.0 eV). In contrast to the blue shift for CO adsorbed on pristine ceria, red shifts of the C-O stretching (vs. free CO) have been calculated for MCO species anchored to ceria. The red shifts of the CO stretching frequency are higher for complexes of Pd and increase with the decreasing coordination of M from MO<sub>4</sub> to MO<sub>3</sub> for a particular metal: Ag/Ce<sub>21</sub>O<sub>42</sub> (50 cm<sup>-1</sup>) < Ag/Ce<sub>21</sub>O<sub>41</sub> (61 cm<sup>-1</sup>) < Pd/Ce<sub>21</sub>O<sub>42</sub> (84 cm<sup>-1</sup>) < Pd/Ce<sub>21</sub>O<sub>41</sub> (113 cm<sup>-1</sup>).

Carboxylate  $\text{CO}_2^-$  and carbonite  $\text{CO}_2^{2-}$  (for Ag-doped NP with an O vacancy) complexes featuring a bent  $\text{CO}_2$  moiety are formed upon CO oxidation at the M/ceria interface. Contrary to AgCO<sub>2</sub>-species, which are easily decomposed via CO<sub>2</sub> detachment, PdCO<sub>2</sub> moieties are prone to withstand decomposition due to significant CO<sub>2</sub> desorption energies of 0.7–1.2 eV. These PdCO<sub>2</sub> moieties anchored to ceria particles could be experimentally detected by the red shifts of the CO<sub>2</sub> asymmetric stretching frequency (vs. that of free CO<sub>2</sub> molecule) by 828 cm<sup>-1</sup> (one O vacancy nearby Pd) and 735 cm<sup>-1</sup> (two O vacancies nearby Pd).

Unlike pristine ceria, carbonate structures at ceria-supported Pd and Ag atoms are hardly formed before CO<sub>2</sub> desorption due to the high barriers of CO<sub>2</sub> transformation to  $CO_3^{2-}$  (up to 1.8 eV for PdCO<sub>3</sub> moieties) and weak CO<sub>2</sub> binding (below ~0.2 eV for AgCO<sub>3</sub> moieties). Detailed analysis of the vibrational spectra of MCO<sub>3</sub>/NP complexes has shown that the two highest v(CO<sub>3</sub>) stretching frequencies lie in the well-resolved intervals 1755–1690 and 1170–1120 cm<sup>-1</sup> for the CO<sub>3</sub> moiety coordinated in a bidentate fashion and 1575–1430 and 1340–1220 cm<sup>-1</sup> for the carbonate groups in the tridentate coordination. These frequency ranges are similar to those for the M-free CO<sub>3</sub>/NP sites. Thus, discrimination between the M-containing and bare ceria particles solely using vibrational spectroscopy data seems hardly possible.

In summary, such characteristics of the studied models as moderately strong CO adsorption, an exothermic CO oxidation process, sufficiently low barriers of MCO to  $MCO_2$  transformations, and ease of  $CO_2$  desorption render CO oxidation by lattice ceria oxygen atoms more favourable at the sites with Ag than with Pd.

**Supplementary Materials:** The following data are available online at https://www.mdpi.com/ article/10.3390/ma14226888/s1: Coordinates of the atomic positions of all considered Ce<sub>21</sub>O<sub>42- $\delta$ </sub>, CO<sub>x</sub>/Ce<sub>21</sub>O<sub>42- $\delta$ </sub>, AgCO<sub>x</sub>/Ce<sub>21</sub>O<sub>42- $\delta$ </sub> and PdCO<sub>x</sub>/Ce<sub>21</sub>O<sub>42- $\delta$ </sub> complexes and selected transition state structures connecting these complexes along with their total energies. Author Contributions: Conceptualisation, V.A.N., E.A.I.-S. and K.M.N.; methodology, V.A.N. and E.A.I.-S.; software, V.A.N. and S.S.L.; investigation, V.A.N. and S.S.L.; validation, V.A.N. and E.A.I.-S.; formal analysis, V.A.N., A.M.S. and E.A.I.-S.; resources, V.A.N. and E.A.I.-S.; data curation, V.A.N.; writing—original draft preparation, V.A.N., E.A.I.-S. and K.M.N.; writing—review and editing, V.A.N., E.A.I.-S., A.M.S., S.S.L. and K.M.N.; visualisation, V.A.N. and A.M.S.; supervision, V.A.N., E.A.I.-S. and K.M.N. All authors have read and agreed to the published version of the manuscript.

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