

Article

Structure and Stability of Iron Fluoride at High Pressure–Temperature and Implication for a New Reservoir of Fluorine in the Deep Earth

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Abstract: Fluorine (F) is the most abundant halogen in the bulk silicate Earth. F plays an important role in geochemical and biological systems, but its abundance and distribution in the terrestrial mantle are still unclear. Recent studies suggested that F reservoirs in the deep mantle are potentially hosted in terrestrial oxide minerals, especially in aluminous bridgmanite. However, the knowledge about the formation and stability field of fluoride in the Earth's interior is rare. In this study, we combine in situ laser-heated diamond anvil cell, synchrotron X-ray diffraction, and first-principles structure search to show that a new tetragonal structure of FeF₃ is stable at pressures of 78–130 GPa and temperatures up to ~1900 K. Simulation predicted the tetragonal phase takes a much denser structure due to the rotation of FeF₆ octahedral units. The equations of states of tetragonal FeF₃ are determined by experiment and verified by simulation. Our results indicate that FeF₃ can be a potential key phase for storing F in the Earth's lower mantle and may explain some mantle-derived magma with high F concentration.

Keywords: iron fluoride; Earth's lower mantle; laser-heated diamond anvil cell; high pressure-temperature

1. Introduction

The halogens (F, Cl, Br, and I) are important volatile elements but their behaviors are much less constrained than their peer volatile elements, such as carbon and hydrogen. Here, we focus on fluorine (F), which is the most abundant and lightest halogen in the Earth's mantle [1]. Previous studies show that the presence of F in the mantle can significantly affect the chemical and physical properties of the silicate melt, including magma viscosity [2] and melting temperature [3].

The F content in the bulk silicate Earth (BSE) is estimated to be ~25 ppm based on concentrations measured in natural basalts and peridotites, which is lower than that in carbonaceous chondrites (~60 ppm) [1]. Previous experiments observed that some mantle-derived magma (e.g., oceanic island basalt (OIB)) contains ~55 ppm of F [4,5], which is higher than the value of BSE. Moreover, Koga and Rose-Koga [6] summarized that F has surprisingly lower content than its neighboring elements like C, N, O, Ne, and Na based on atomic numbers and the depletion of F might be stored in deep Earth. All the above statements indicate that the distribution of F might be heterogeneous and there could be a potential F reservoir in the deep mantle. This hypothesis was implied by some terrestrial oxide mineral samples, which can contain several thousand ppm F [7–12]. There are also known F-containing minerals in the lower mantle, such as oskarssonite AlF₃, parascandolaite KMgF₃, and a new, unnamed



fluoride, (Ba,Sr)AlF₃ [13,14]. Du et al. [15] predicted that the heavy halogens (Cl, Br, and I) can stay in the Earth's core through the formation of Fe-halides based on the first-principles structural searches up to 360 GPa. Since a large Fe³⁺ excess in the lower mantle can be created by the disproportionation of Fe²⁺ combined with removal of Fe⁰ droplets to the core [16,17], FeF₃ could become an important F carrier in the deep mantle. Although an X-ray diffraction study on FeF₃ up to 23 GPa at room temperature has been performed [18], there is no study about FeF₃ at the Earth's lower-mantle conditions. In this study, we investigate the structure and stability of FeF₃ at high pressure and high temperature corresponding to the conditions from the middle to lowermost mantle (78–130 GPa and up to ~1900 K).

2. Methods

Global structural optimization was performed using the CALYPSO (Crystal structure AnaLYsis by Particle Swarm Optimization) code [19,20] with the particle swarm optimization algorithm, which has successfully predicted structures of various systems ranging from elements to binary and ternary compounds [21–23]. The total energy calculations were carried out using density functional theory within the Perdew–Burke–Ernzerhof (PBE) generalized gradient approximation [24] as implemented in the VASP (Vienna Ab initio Simulation Package) code [25]. We used projector augmented waves [26] with $3s^23p^63d^64s^2$ and $2s^22p^5$ as valence electrons for Fe and F atoms, respectively. The use of a plane-wave kinetic energy cutoff of 550 eV and adoption of dense k-point sampling were shown to give excellent convergence of total energies. To determine the dynamical stability of the studied structures, we performed phonon calculations by using the finite displacement approach, as implemented in the PHONOPY code [27].

A mixture of FeF₃ (amorphous, 1 μ m grain size, purity 99.999%) and Au powder was ground in a ball mill for 2 h. The mixture of the amorphous FeF₃ and Au powder was cold pressed into a thin foil about $\sim 30 \times 30 \ \mu\text{m}^2$ in size and $\sim 10 \ \mu\text{m}$ in thickness. The sample foil was loaded in a 55 μ m diameter hole in a rhenium gasket indented by diamond anvils with 100 µm culet diameter in a symmetric Mao-Bell diamond anvil cell [28]. Neon gas was used as pressure medium in our experiments. Samples were heated in a double-sided laser-heated diamond anvil cell (DAC) [29] and examined in situ at high pressures and temperatures by synchrotron X-ray powder diffraction at beamline 16-IDB of the Advanced Photon Source, Argonne National Laboratory. The X-ray beam size was $4.2 \times 5.3 \ \mu\text{m}^2$ with a wavelength of 0.4066 Å. For laser heating, the diameter of a laser heating spot was up to \sim 40 μ m at ~1900 K in the flat top area created with two focused ytterbium fiber lasers from both sides of the sample that minimized both radial and axial temperature gradients. Temperatures were determined by fitting the thermal radiation from the central portion of the heated sample to the Planck radiation function [30]. About 2 wt.% Au powder was added to the mixture for infrared laser absorption in the laser-heating experiment. Pressures were determined before and after heating by the self-consistent equation of states (EOS) of the doped Au [31], and the error bar of calibrated pressure is estimated to be of the order 1 to 2 GPa.

3. Results

We first performed structure predictions containing up to four formula units (f.u.) in the simulation cell over a wide range of pressures (0, 100, and 200 GPa). At ambient pressure, we found that the most stable structure was the *R*-3*c* phase, in agreement with the previous experiments [18]. At both 100 and 200 GPa, we uncovered a group of new structures, indicating potential new phases induced by pressure. Among them, the tetragonal $P4_22_12$ structure was found to be the most stable at 200 GPa. The calculated enthalpy curves (relative to the *R*-3*c* structure, Figure 1a) illustrate the relative thermodynamic stabilities of the structures. Our results show that, up to 136 GPa, the ambient *R*-3*c* structure is the most stable one. Beyond 136 GPa and up to 200 GPa, the predicted $P4_22_12$ structure becomes more stable. The phase transition from the ambient *R*-3*c* to the tetragonal $P4_22_12$ involves rotation of FeF₆ octahedrons. This enables the structure to pack closely and, consequently, results in larger volume reduction across the phase transition at high pressure. To further understand the effect

of pressure, we also plotted the enthalpy components (internal energy term U and pressure-volume, pV, terms) of the $P4_22_12$ structure relative to the R-3c structure in Figure 1a. The R-3c structure has the lowest internal energy over the entire pressure range calculated in this work. However, beyond 136 GPa, the enthalpy of the $P4_22_12$ structure becomes lower than that of the R-3c structure. The main contribution to the enthalpy at high pressure is the pV term. At 136 GPa, the relative volume reduction of the $P4_22_12$ structure is 1.6% per formula unit, becoming large enough to compensate for the internal energy difference. The dynamic stabilities of the newly predicted $P4_22_12$ structure was examined by calculating the phonon spectra using the supercell method. No imaginary phonon frequencies were found in the whole Brillouin zone over the studied pressure range, establishing the dynamical stability of the $P4_22_12$ structure (Figure 1b).



Figure 1. Enthalpy curves and phonon dispersion of the $P4_2212$ structure. (**a**) The difference in enthalpy, internal energy *U*, and the *pV* term between the $P4_22_12$ structure and the *R*-3*c* structure per formula unit as a function of pressure. (**b**) Calculated phonon dispersion of the $P4_22_12$ structure at 140 GPa.

We then conducted high-pressure experiments in search of the simulated structure. The mixture of the amorphous FeF₃ and Au powder was first compressed to 78 GPa at room temperature with Au as an internal pressure standard [31]. The sample was then heated using a double-sided laser system and held at target temperatures for ~10 min at ~1900 K. At 78 GPa and 1900 K, we observed the formation of the predicted tetragonal $P4_22_12$ phase. The sample was subsequently compressed to 130 GPa at room temperature in nine pressure steps (Table 1). At each pressure step, the sample was annealed for 10 min at ~1800 K, and in situ X-ray diffraction (XRD) spectra were acquired before and after the annealing. The diffraction data show that the tetragonal $P4_22_12$ phase was stable in the pressure range from 78 to 130 GPa and up to 1800 K. It is worth mentioning that our experimental temperatures were lower than an average geotherm of the lowermost mantle, although they were closer to that of the cold slab subduction regions [32]. In Figure 2, we show an XRD pattern collected at 119 GPa after laser heating to 1800 K. All peaks indexed to the calculated tetragonal unit cell are assigned in the XRD pattern with a few minor peaks indexed to a pyrite-type FeF₂ (Figure 2), which is also a stable structure based on our calculation. The minor pyrite-FeF₂ phase could be formed by releasing F from FeF₃ under high temperature or due to the lack of F in the starting composition. Unit cell parameters of the tetragonal P4₂2₁2 phase acquired at 300 K from 78–130 GPa are summarized in Table 1, and the PV data are plotted in Figure 3. We fit the PV data to the third-order Birch–Murnaghan equation of state:

$$P(V) = \frac{3K_0}{2} \left[\left(\frac{V_0}{V} \right)^{\frac{2}{3}} - \left(\frac{V_0}{V} \right)^{\frac{5}{3}} \right] \left\{ 1 + \frac{3}{4} \left(K'_0 - 4 \right) \left[\left(\frac{V_0}{V} \right)^{\frac{2}{3}} - 1 \right] \right\}$$
(1)

and determined $K_0 = 109(10)$ GPa and $K_{0'} = 3.9(2)$. The tetragonal $P4_22_12$ phase of FeF₃ is more compressible than iron and iron hydrides [33] but less than FeH₅ [34]. We also calculated the equation of state of the tetragonal FeF₃, and listed the result in Table 2 from the first-principles simulation and plotted in Figure 3 together with the experimentally measured data for comparison.

P (GPa)	a (Å)	c (Å)	<i>V</i> /(Å ³ /f.u.)	Density (g/cm ³)
78	4.010(2)	6.992(2)	28.11(2)	6.501(5)
81	3.993(2)	6.980(2)	27.83(3)	6.567(7)
90	3.983(2)	6.939(4)	27.52(2)	6.639(5)
98	3.949(2)	6.896(3)	26.88(2)	6.799(5)
106	3.919(1)	6.836(4)	26.25(2)	6.962(5)
110	3.910(1)	6.806(4)	26.01(2)	7.026(5)
118	3.905(1)	6.778(3)	25.84(2)	7.072(6)
122	3.884(1)	6.765(4)	25.69(2)	7.113(6)
126	3.878(1)	6.756(4)	25.40(2)	7.193(6)
130	3.854(1)	6.725(4)	24.98(2)	7.316(6)

Table 1. Lattice parameters of tetragonal FeF₃ up to 130 GPa and 1900 K.



Figure 2. X-ray diffraction patterns at 119 GPa (Au pressure), 300 K after laser heating. At the center of the heated area, the majority of diffraction peaks (other than those from Au and Re) belong to the tetragonal $P4_22_12$ phase of FeF₃. A few minor peaks are consistent with the pyrite-type FeF₂. Data were acquired using an X-ray wavelength of 0.4066 Å.



Figure 3. Measured FeF₃ unit cell volumes at high pressure and 300 K. Open purple circles from this study show pressures based on the Au pressure calibration (the error bar of ± 2 GPa). The solid curve at relatively high pressures from 78 to 130 GPa shows the equation of state for FeF₃ with a tetragonal unit-cell structure. Solid red circles are calculated by first-principles simulation. The equation of state of high-pressure tetragonal FeF₃ is compared with the low-pressure trigonal phase [18].

a (Å)	c (Å)	V (Å ³ /f.u.)	Density (g/cm ³)
4.435	7.867	38.68	4.853
4.290	7.425	34.17	5.493
4.205	7.104	31.40	5.978
4.138	6.901	29.54	6.354
4.083	6.747	28.12	6.675
4.037	6.625	26.99	6.955
3.996	6.527	26.06	7.203
3.961	6.442	25.26	7.431
	<i>a</i> (Å) 4.435 4.290 4.205 4.138 4.083 4.037 3.996 3.961	a (Å) c (Å) 4.435 7.867 4.290 7.425 4.205 7.104 4.138 6.901 4.083 6.747 4.037 6.625 3.996 6.527 3.961 6.442	a (Å) c (Å) V (Å ³ /f.u.)4.4357.86738.684.2907.42534.174.2057.10431.404.1386.90129.544.0836.74728.124.0376.62526.993.9966.52726.063.9616.44225.26

Table 2. Lattice parameters of tetragonal FeF₃ from first-principles simulation.

4. Discussion

Recent synthesis experiments showed that Al-bearing bridgmanite under the conditions of the uppermost lower mantle can contain up to ~1.3 wt.% F [12], which is much higher than in the other oxide minerals, such as ~5100 ppm in olivine [7–9], ~660 ppm in pyroxenes [10,11], and ~1100 ppm in pyrope [8] and ~2110 ppm in wadsleyite [11]. The concentrations of F incorporated in oxide minerals can be affected by water contents due to the similar charge and comparable ionic radius of F^- and OH⁻ [35,36]. That is, the presence of F may play a similar role as water in impacting the mineralogy of the deep mantle and the distribution of water in the Earth's interior [37–42]. Our experimental results show that the dense FeF₃ (~6.5–7.3 g/cm³, see Table 1) with a tetragonal unit-cell structure is gravitationally stable at the bottom of the lower mantle and indicate that it can be a potential F reservoir at such depths. This further supports the speculation of an F-enriched source region in the Earth's interior [4,5,12] and that may explain the phenomenon of the content of F being much lower than other elements with similar atomic numbers [6]. To better constrain distribution of F and other halogens in the BSE, however, required more dedicated high pressure-temperature experiments along average mantle geotherm conditions, such as F partitioning between FeF₃ and oxide minerals, its crystal chemistry, and the stabilities of other halogenides [15].

Author Contributions: Y.L. and Q.H. designed the research; Y.L., Q.H., and Y.M. performed experiments; Q.H. and Y.L. analyzed data; L.Z. did the model calculation; and Y.L. and Q.H. wrote the paper with input from all authors. All authors have read and agreed to the published version of the manuscript.

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Conflicts of Interest: The authors declare no conflict of interest.

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