

## Article

# A System Size Analysis of the Fireball Produced in Heavy-Ion Collisions <sup>†</sup>

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**Abstract:** One of the main interests of high-energy physics is the study of the phase diagram and the localization of phase transitions from hadronic to quark–gluonic matter. There are different techniques to study the hot matter. One of them is femtoscopy, which uses two-particle correlations to extract spatiotemporal characteristics of the emission source. Another approach involves obtaining thermodynamic parameters from the momentum distributions of produced particles based on various theoretical models. In this research, we perform a comparative analysis of femtoscopic volumes and volumes obtained using the Tsallis statistical fit. This analysis allows us to estimate system size at the time of kinetic freeze-out and its dependence on collision centrality and energy. We observe that at high energies, the volume values estimated taking the two approaches diverge significantly, while at low energies, they are more consistent. In the future, these results can help to combine these two different methods and provide a more comprehensive picture of the fireball produced in heavy-ion collisions.

**Keywords:** quark–gluon matter; phase transition; heavy ion; fireball evolution



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## 1. Introduction

Recently, a large amount of experimental data have been collected in high-energy physics for studying the properties of nuclear matter. One of the main interests is the study of phase transitions from regular hadronic matter to a deconfined state called quark–gluon plasma (QGP) [1–4]. This transition can occur at high temperatures and densities and can be shown on the Quantum ChromoDynamics (QCD) phase diagram in terms of temperature and baryonic chemical potential.

At low baryonic chemical potential, lattice calculations indicate the presence of a smooth transition—a crossover [5,6]. However, at higher chemical potential, lattice calculations are not possible, and effective models have to be used instead. These models suggest the presence of a first-order phase transition and the critical point separating it from the crossover.

QGP can be formed in ultrarelativistic heavy-ion collisions. After its formation, the hot matter expands and cools, which leads to the transition into the hadronic phase and reaches a state of statistical equilibrium until chemical freeze-out, when inelastic interactions between particles cease. Kinetic freeze-out occurs when hadrons stop interacting elastically, and their kinematic properties remain unchanged.

## 2. Materials and Methods

In our research, we use the published data from the STAR experiment obtained during the first stage of the Beam Energy Scan program at RHIC for Au+Au collisions at  $\sqrt{s_{NN}} = 7.7, 11.5, 19.6, 27$  and  $39$  GeV, aimed at localizing the phase transition boundary [7,8].

The main goal of this research is to estimate the size of the system formed in heavy-ion collisions using different methods and to compare the results obtained through different observables.

The first method is based on the application of a statistical model using Tsallis statistics [9–11]. The fireball is considered a thermodynamic system in the grand canonical ensemble (GCE), and thermodynamic characteristics, such as temperature and system volume, are extracted from the analysis of transverse momentum spectra.

The second method relies on femtoscopy. The interferometry method was initially proposed by Hanbury Brown and Twiss [12] in astrophysics and was later adapted for use in particle physics [13–15]. It has been applied in studies across a wide range of energies [16–19] to extract spatiotemporal properties of the homogeneity region at kinetic freeze-out by analyzing two-particle correlations of final particles.

### 2.1. Tsallis-3 Statistics

Tsallis entropy is defined as [9]

$$S = \sum_i \frac{p_i^q - p_i}{1 - q}, \quad \sum_i p_i = 1, \tag{1}$$

where  $p_i$  is the probability of  $i$ -th microscopic state of the system and  $q \in [0, \infty]$  is the entropic parameter. In the Gibbs limit ( $q \rightarrow 1$ ), the Tsallis entropy (1) recovers the Boltzmann–Gibbs entropy [9]:

$$S = \sum_i p_i \ln p_i. \tag{2}$$

We can see that the Tsallis entropy generalizes the Boltzmann–Gibbs entropy, with the parameter  $q$  indicating the degree of deviation of a given generalized probability distribution from the Boltzmann–Gibbs one. There are a few variants of the Tsallis statistics [10].

In the Tsallis-3 statistics (Tsallis statistics with escort probabilities) [10,11] in the grand canonical ensemble, the thermodynamic potential  $\Omega$  of the system can be written as

$$\Omega = \langle H \rangle - TS - \mu \langle N \rangle, \tag{3}$$

where  $\langle H \rangle = \frac{1}{\theta} \sum_i p_i^q E_i$  is the mean energy of the system,  $\langle N \rangle = \frac{1}{\theta} \sum_i p_i^q N_i$  is the mean number of particles, and  $\theta = \sum_i p_i^q$ .

Let us consider a relativistic ideal gas in the Tsallis-3 statistics in the grand canonical ensemble with the Maxwell–Boltzmann statistics of particles. According to the principle of thermodynamic equilibrium for an open system, similar to the principle of maximum entropy for an isolated system, normalization expressions for the parameters  $\Lambda = -\theta TS + \langle H \rangle - \mu \langle N \rangle$  and  $\theta$  can be expressed as [11]

$$\begin{cases} 1 = \sum_{n=0}^{n_0} \frac{\omega^n}{n! \Gamma(\frac{1}{q-1})} \int_0^\infty t^{\frac{2-q}{q-1}-n} e^{-t+\beta'(\Lambda+\mu n)} (\mathbf{K}_2(\beta' m))^n dt, \\ \theta = \sum_{n=0}^{n_0} \frac{\omega^n}{n! \Gamma(\frac{q}{q-1})} \int_0^\infty t^{\frac{1}{q-1}-n} e^{-t+\beta'(\Lambda+\mu n)} (\mathbf{K}_2(\beta' m))^n dt, \end{cases} \quad (4)$$

where  $\omega = \frac{gV}{2\pi^2} \frac{m^2 T \theta^2}{q-1}$ ,  $\beta' = \frac{-t(1-q)}{T\theta^2}$  and  $n_0$  is the number of terms of the series to be taken into account, starting from zero. The values of  $\Lambda$  and  $\theta$  are found by numerically solving the system of Equation (4). The expression for transverse momentum distribution in rapidity range  $y \in [y_{min}, y_{max}]$  takes the following form [11]:

$$\left. \frac{d^2 N}{p_T dp_T dy} \right|_{y_{min}}^{y_{max}} = \frac{gV}{(2\pi)^2} m_T \int_{y_{min}}^{y_{max}} dy \cosh y \frac{1}{\theta} \sum_{n=0}^{n_0} \frac{\omega^n}{n! \Gamma(\frac{q}{q-1})} \int_0^\infty t^{\frac{1}{q-1}-n} e^{-t+\beta'(\Lambda-m_T \cosh y + \mu(n+1))} (\mathbf{K}_2(\beta' m))^n dt. \quad (5)$$

The fitting of the final particle spectra was performed using the ROOT analysis framework with MINUIT2 numerical minimization software library [20]. The technique of simultaneous fitting of all particle types by  $\chi^2$ -minimization method was applied, where parameters of the medium, such as temperature  $T$  and the parameter  $q$ , are common for all types of particles. The fitting procedure involves substituting some initial guess fit parameters (the state variables of GCE) into the system of Equation (4), which is then solved numerically to determine the values of the thermodynamic normalization functions  $\Lambda$  and  $\theta$ . Calculated values of the normalization functions, along with the initial parameters, are then used in the final Formula (5) for the transverse momentum distribution for  $\chi^2$  minimization. The procedure can then be applied iteratively to reach the global minimum of fit parameters. For simultaneous fitting, we minimized the sum of  $\chi^2$  for all particle species. Figure 1 shows the Tsallis-3 model fits for  $K^\pm$ ,  $\pi^\pm$ , at  $\sqrt{s_{NN}} = 7.7, 19.6, 39$  GeV and centrality classes: 0–5%, 30–40%, 70–80%. The parameter values for all energies and centralities are presented in Table A1. The Tsallis-3 statistics are sensitive to the particle type; therefore, during the fitting process, protons were excluded, and only mesons were included in the fit. Since the next method uses combined data from  $\pi^+$  and  $\pi^-$ , their spectra are also combined here. The low  $p_T$  part of the pion spectra up to 0.5 GeV/c is affected by resonance decays and was excluded from the calculations. The fit ranges were [0.25, 2.00] GeV/c for kaons and [0.50, 2.00] GeV/c for pions. The volume values were obtained as a free fit parameter.

### 2.2. Femtoscopic Approach

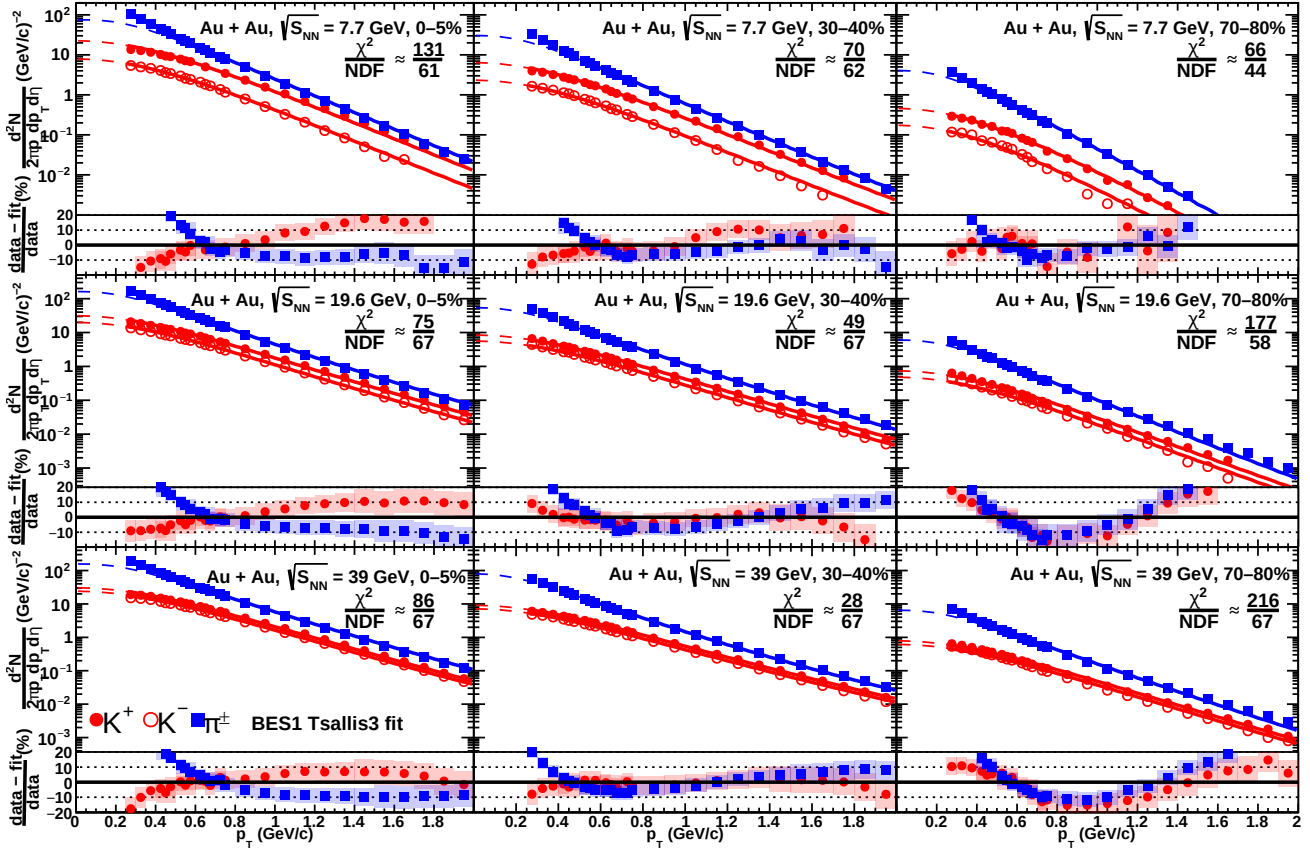
The experimental correlation function is constructed using distributions of the relative pair momentum  $\vec{q} = \vec{p}_1 - \vec{p}_2$ :

$$C(\vec{q}) = \frac{N(\vec{q})}{D(\vec{q})}, \quad (6)$$

where the numerator  $N(\vec{q})$  is constructed with pairs of particles from the same event and, thus, contains information on quantum statistics and effects of particle final state interactions [21], while the denominator  $D(\vec{q})$  uses uncorrelated pairs obtained from mix-events. The relative pair momentum  $\vec{q}$ , expressed in the longitudinal co-moving system (LCMS), is projected onto the Bertsch–Pratt coordinate system [22,23]. The  $q_{out}$  component is directed along the pair transverse momentum  $\vec{k}_T = (\vec{p}_{T,1} + \vec{p}_{T,2})/2$ , the  $q_{long}$  component is directed along the beam axis, and the  $q_{side}$  component is perpendicular to the other two directions. Further, the Bowler–Sinyukov Gaussian parametrization [24,25] is applied:

$$C(\vec{q}) \propto (1 - \lambda) + K_{Coul}(q_{inv}) \lambda \exp(-q_0^2 R_0^2 - q_s^2 R_s^2 - q_l^2 R_l^2 - 2q_o q_s R_{os}^2 - 2q_o q_l R_{ol}^2), \quad (7)$$

where  $K_{\text{Coul}}(q_{\text{inv}})$  is the Coulomb correction factor,  $\lambda$  is the correlation strength related to secondary particles from resonance decays,  $R_{\text{out,side,long}}$  are the correlation radii, and  $R_{\text{out-side,out-long}}$  are the cross terms. The particle emission source correlation radii are extracted by fitting. The radius  $R_{\text{side}}$  depends on the geometrical size of the source,  $R_{\text{out}}$  is influenced by both the geometrical size and the emission duration, and  $R_{\text{long}}$  depends on the system's lifetime.



**Figure 1.** Tsallis-3 model fits of  $K^\pm$ ,  $\pi^\pm$  at  $\sqrt{s_{\text{NN}}} = 7.7, 19.6, 39$  GeV and the centrality classes 0–5%, 30–40%, 70–80%. Data points are taken from [7].

We use the data measured by the STAR collaboration during the BES-I program [8]. The dependence of the correlation radii on transverse mass,  $m_T$ , is considered below. The homogeneity region of the system is smaller for pairs with larger  $m_T$  [26]. Therefore, to assess the entire volume of the system, we extrapolate the  $m_T$ -dependence of femtoscopic radii to the minimum possible transverse mass,  $m_T = m$ . We employed several parameterizations for this extrapolation: the first is grounded in basic hydrodynamic assumptions (known as  $m_T$ -scaling [27])

$$R_i \propto \frac{1}{\sqrt{m_T}}, \tag{8}$$

the second uses a simple power-law model

$$R_i \propto m_T^{-\alpha}, \tag{9}$$

and the third applies a model based on the Blast-Wave expansion scenario [28]

$$R_{\text{side}} = \frac{R_0}{\sqrt{1 + \rho_0^2(m/T + 1/2)}}, \quad R_{\text{long}} = \tau \sqrt{\frac{T}{m_T} \frac{K_2(m_T/T)}{K_2(m_T/T)}}, \tag{10}$$

where  $i = \text{out, side, long}$ ,  $\rho_0$  is the maximum flow rapidity, and  $T$  is the kinetic freeze-out temperature. The values of these parameters are obtained by fitting  $\pi^\pm, K^\pm, p, \bar{p}$  transverse momentum spectra using the Blast-Wave model [29] and are fixed for further analysis. These parameter values are listed in Table A2. We use linear transverse flow rapidity profile  $\rho = \rho_0(r/R)$  to remain consistent with previous HBT studies [18,30].  $R_0$  and  $\tau$  are free normalization parameters that are not of interest in this study. Figure 2 shows the extrapolation of the femtosopic  $\pi^\pm$  radii  $R_{\text{side}}$  and  $R_{\text{long}}$  to zero transverse momentum. The extrapolated values for all energies and centralities are presented in Table A3.

The radii obtained from femtoscopy follow Gaussian distributions and thus characterize only the average size of the system, whereas the Tsallis statistics provide the total volume of the system. To compare the femtosopic radii with the statistical model, it is necessary to recalculate them by equating the variances of the Gaussian and hard sphere distributions:

$$\sigma_{\text{hard}}^2 = \langle x^2 \rangle = \frac{\int_V x^2 dV}{\int_V dV} = \frac{R_{\text{hard}}^2}{5}, \quad \sigma_{\text{Gauss}}^2 = R_{\text{Gauss}}^2, \quad (11)$$

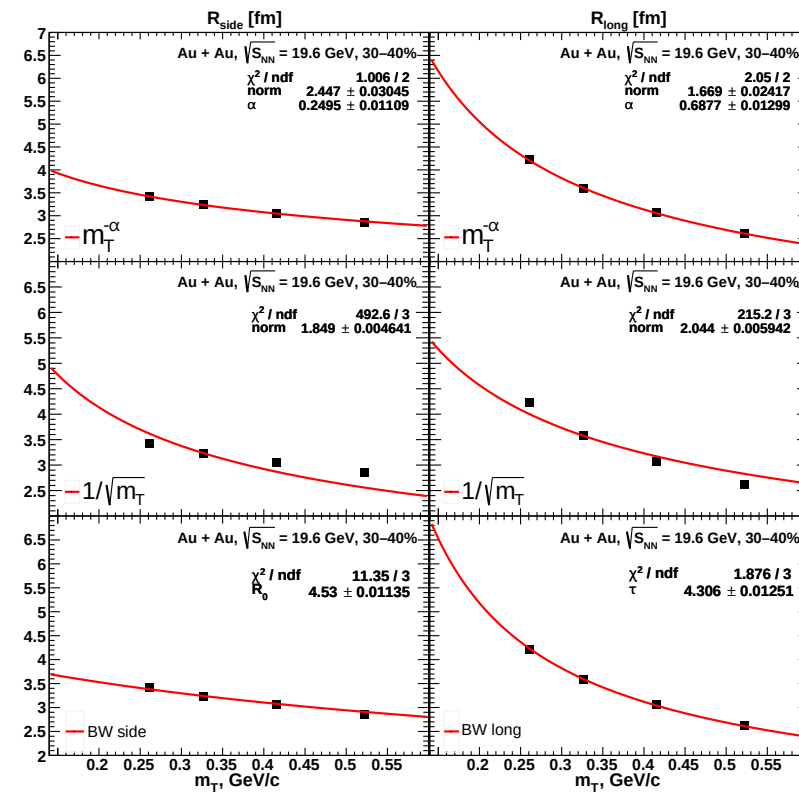
where  $i = \text{out, side, long}$ , and  $dV = r^2 \sin \theta dr d\phi d\theta$  is the volume element in spherical coordinates. Then, the relation of the radii is

$$R_{\text{hard}} = \sqrt{5} R_{\text{Gauss}}. \quad (12)$$

Subsequently, the system volume is calculated as

$$V = \frac{4}{3} \pi (5)^{\frac{3}{2}} R_{\text{side}}^2 R_{\text{long}}, \quad (13)$$

where we use  $R_{\text{side}}$  instead of  $R_{\text{out}}$  because  $R_{\text{out}}$  depends on the particle emission time [31] and does not reflect the exact geometric size of the system.



**Figure 2.** Femtosopic  $\pi^\pm$  radii fits at  $\sqrt{s_{\text{NN}}} = 19.6$  GeV and the centrality class 30–40% using Equations (8)–(10). Data points are taken from [8].

### 3. Results and Discussion

Figure 3 shows the volume values obtained by femtoscopy and by fitting the particle momenta distributions using the Tsallis-3 statistics. The volume values for all energies and centralities are presented in Table A4. The temperature values decrease from central to peripheral events, which is consistent with a previous study [32] using the same data and a similar model based on q-dual statistics [33]. However, the behavior of the parameter  $q$  and the values of temperature and  $q$  differ significantly. This is due to the fact that in this study, we used the exact Tsallis-3 statistics, taking into account the higher-order terms, while the previous study used a zeroth term approximation. As demonstrated in [11], including higher-order terms in the expansion significantly impacts the parameter values, increasing the temperature and decreasing the  $q$  parameter.

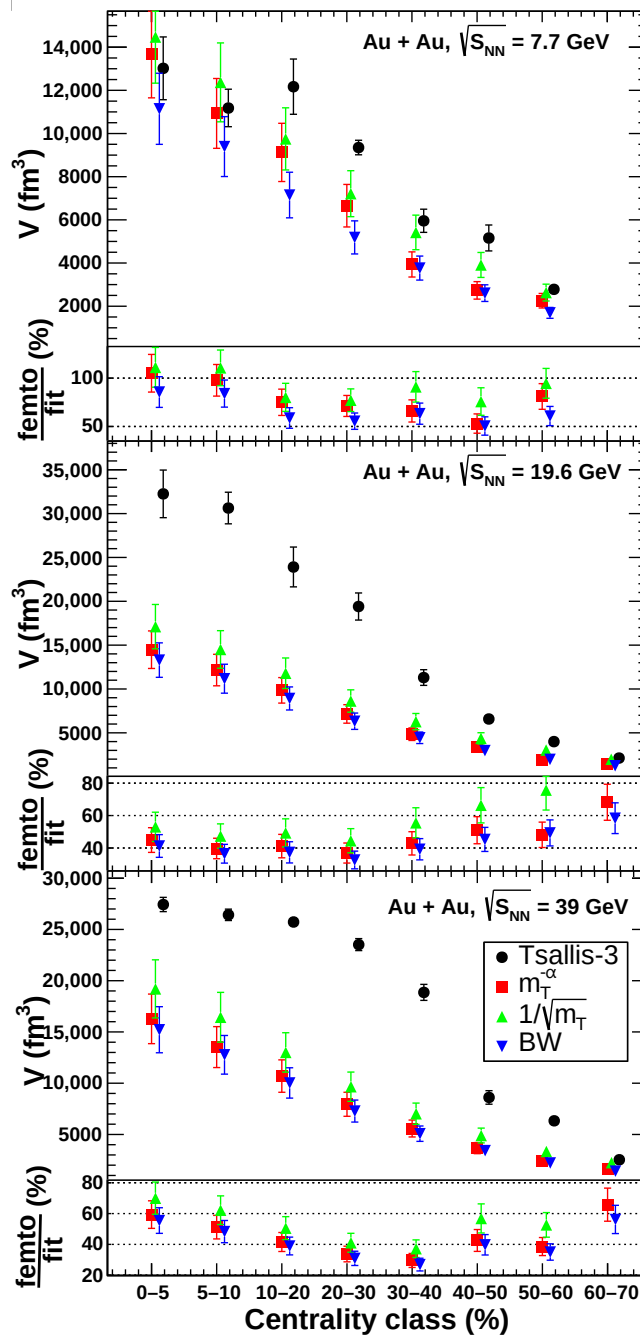


Figure 3. Comparison of volumes at  $\sqrt{s_{NN}} = 7.7, 19.6, 39$  GeV and the centrality classes from 0–5% to 60–70% obtained using the Tsallis-3 statistics and calculated using Equations (8)–(10).

It was established that the  $m_T$ -scaling (8) provides an unsatisfactory description of the femtoscopic radii, significantly overestimating  $R_{\text{side}}$  and underestimating  $R_{\text{long}}$ . The power-law dependence (9) demonstrates that the parameter  $\alpha$  deviates substantially from the value equal to 0.5, which may indicate contributions from flow and collective effects. The Blast-Wave model (10) accounts for these factors and provides a satisfactory description of the correlation radii.

At low energies, the volumes agree with each other; however, with increasing energy, the volumes begin to diverge, with femtoscopic volumes appearing smaller. This can be explained by several factors:

- Femtoscopic correlations measure only the region of homogeneity [34], while in the statistical approach using the Tsallis distribution, we obtain the volume of the whole system. If the collective expansion is strong, the region of homogeneity can be smaller than the entire source volume [21].
- The correlation function may deviate from a Gaussian shape due to exponential tails caused by resonance decay contributions [35]. This can lead to an underestimation of the femtoscopic radii.
- In the Tsallis statistics, the particle interactions and collective effects, including rescattering, are modeled phenomenologically by deformation of entropy using the nonextensive entropic parameter  $q$ .

#### 4. Conclusions

In this study, we compared two different methods for the determination of the size of the system formed in ultrarelativistic heavy-ion collisions at kinetic freeze-out using femtoscopy and a statistical approach based on the Tsallis distribution. A procedure for recalculating correlation radii was presented, ensuring proper comparison between the methods. The results show that collective effects, such as flow, strongly influence the femtoscopic radii. Simple models like  $m_T$  scaling fail to describe these radii accurately, while the Blast-Wave model provides a better description.

One key finding is that femtoscopic volumes tend to be smaller than the volumes obtained in the Tsallis model, particularly at higher collision energies. This difference points to the need for further investigation, as it suggests potential underlying physics not fully captured by either method in their current forms.

Future work could refine this approach by calculating the system volume as the volume of a cylinder at kinetic freeze-out instead of relying on a simple product of radii. This might give more precise estimates and make the comparison between the two methods more reliable. Adding Tsallis statistics to the Blast-Wave framework could also help to improve the analysis.

Combining these methods can provide a more complete understanding of the system's size and behavior. This will help to study possible phase transitions and better understand the properties of hot and dense matter in heavy-ion collisions.

**Author Contributions:** Conceptualization, A.A.; Methodology, A.P. and V.B.L.; Formal analysis, E.N.; Writing—original draft, E.N.; Writing—review and editing, A.A., A.P. and V.B.L.; Visualization, E.N.; Supervision, A.A.; Funding acquisition, A.A. All authors have read and agreed to the published version of the manuscript

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**Data Availability Statement:** The data presented in this study are available on request from the corresponding author.



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## Appendix A

**Table A1.** Tsallis-3 model fit of  $\pi^\pm, K^\pm$   $p_T$ -spectra at  $\sqrt{s_{NN}} = 7.7, 11.5, 19.6, 27, 39$  GeV and centralities from 0–5%, to 70–80%.  $p_T$  ranges: [0.5, 2], [0.25, 2] (GeV/c) are taken for  $\pi^\pm, K^\pm$ , correspondingly.

Centrality	$T$ (MeV)	$q$	$\chi^2/\text{NDF}$
7.7 GeV			
0–5%	$187.5 \pm 1.5$	$1.0221 \pm 0.0022$	131/61
5–10%	$184.9 \pm 1.7$	$1.0227 \pm 0.0031$	93/63
10–20%	$186.9 \pm 1.1$	$1.0346 \pm 0.0034$	67/63
20–30%	$178.2 \pm 1.0$	$1.0331 \pm 0.0029$	73/63
30–40%	$170.5 \pm 0.8$	$1.0277 \pm 0.0030$	70/62
40–50%	$161.7 \pm 0.9$	$1.0326 \pm 0.0029$	65/57
50–60%	$151.6 \pm 0.9$	$1.0206 \pm 0.0031$	81/57
60–70%	$143.8 \pm 1.0$	$1.0157 \pm 0.0032$	73/53
70–80%	$136.8 \pm 0.9$	$1.0082 \pm 0.0041$	66/44
11.5 GeV			
0–5%	$201.3 \pm 1.7$	$1.0330 \pm 0.0029$	107/63
5–10%	$197.6 \pm 0.8$	$1.0316 \pm 0.0007$	75/65
10–20%	$195.6 \pm 1.1$	$1.0384 \pm 0.0031$	56/65
20–30%	$190.8 \pm 0.8$	$1.0412 \pm 0.0028$	41/65
30–40%	$182.8 \pm 1.0$	$1.0447 \pm 0.0024$	55/64
40–50%	$173.0 \pm 0.8$	$1.0363 \pm 0.0020$	62/64
50–60%	$164.5 \pm 0.9$	$1.0323 \pm 0.0019$	60/63
60–70%	$154.9 \pm 0.8$	$1.0228 \pm 0.0023$	82/56
70–80%	$147.1 \pm 0.9$	$1.0163 \pm 0.0025$	100/51
19.6 GeV			
0–5%	$218.9 \pm 3.5$	$1.0479 \pm 0.0010$	75/67
5–10%	$216.3 \pm 2.3$	$1.0520 \pm 0.0005$	70/67
10–20%	$210.1 \pm 0.9$	$1.0502 \pm 0.0025$	57/67
20–30%	$201.9 \pm 1.0$	$1.0527 \pm 0.0017$	27/67
30–40%	$191.6 \pm 0.9$	$1.0443 \pm 0.0019$	49/67
40–50%	$182.1 \pm 0.8$	$1.0377 \pm 0.0017$	92/65
50–60%	$174.0 \pm 0.8$	$1.0343 \pm 0.0016$	159/65
60–70%	$165.2 \pm 0.8$	$1.0279 \pm 0.0017$	229/61
70–80%	$154.6 \pm 1.1$	$1.0219 \pm 0.0020$	177/58



**Table A1.** *Cont.*

Centrality	$T$ (MeV)	$q$	$\chi^2/\text{NDF}$
27 GeV			
0–5%	$221.8 \pm 1.9$	$1.0421 \pm 0.0031$	103/66
5–10%	$221.2 \pm 1.4$	$1.0465 \pm 0.0004$	76/67
10–20%	$211.4 \pm 0.1$	$1.0490 \pm 0.0001$	117/67
20–30%	$209.5 \pm 1.2$	$1.0570 \pm 0.0022$	32/67
30–40%	$201.3 \pm 0.0$	$1.0528 \pm 0.0009$	28/67
40–50%	$191.3 \pm 1.1$	$1.0486 \pm 0.0016$	49/67
50–60%	$182.1 \pm 0.9$	$1.0424 \pm 0.0014$	94/67
60–70%	$172.5 \pm 0.9$	$1.0351 \pm 0.0014$	180/67
70–80%	$163.8 \pm 0.9$	$1.0284 \pm 0.0013$	255/67
39 GeV			
0–5%	$232.6 \pm 1.2$	$1.0454 \pm 0.0007$	86/67
5–10%	$231.6 \pm 1.3$	$1.0499 \pm 0.0007$	64/67
10–20%	$229.1 \pm 1.0$	$1.0568 \pm 0.0007$	42/67
20–30%	$218.6 \pm 0.3$	$1.0614 \pm 0.0008$	19/67
30–40%	$205.0 \pm 0.2$	$1.0612 \pm 0.0012$	28/67
40–50%	$199.9 \pm 1.2$	$1.0521 \pm 0.0016$	56/67
50–60%	$185.0 \pm 0.1$	$1.0493 \pm 0.0010$	134/67
60–70%	$182.2 \pm 1.1$	$1.0422 \pm 0.0014$	143/67
70–80%	$174.9 \pm 1.1$	$1.0354 \pm 0.0015$	216/67

**Table A2.** Blast-Wave model fit results for  $\pi^\pm, K^\pm, p, \bar{p}$   $p_T$ -spectra at  $\sqrt{s_{\text{NN}}} = 7.7, 11.5, 19.6, 27, 39$  GeV and centralities from 0–5%, to 70–80%.  $p_T$  ranges: [0.5, 1.50], [0.25, 1.5], [0.4, 1.5] (GeV/ $c$ ) are taken for  $\pi^\pm, K^\pm, p(\bar{p})$ , correspondingly.

Centrality	$\rho_0$	$T$ (MeV)	$\chi^2/\text{NDF}$
7.7 GeV			
0–5%	$0.743 \pm 0.016$	$112 \pm 3$	76/100
5–10%	$0.710 \pm 0.017$	$115 \pm 3$	64/102
10–20%	$0.649 \pm 0.018$	$122 \pm 3$	48/106
20–30%	$0.605 \pm 0.021$	$123 \pm 3$	79/104
30–40%	$0.544 \pm 0.023$	$129 \pm 3$	81/104
40–50%	$0.470 \pm 0.028$	$133 \pm 4$	88/99
50–60%	$0.391 \pm 0.035$	$136 \pm 5$	101/97
60–70%	$0.340 \pm 0.034$	$136 \pm 4$	80/95
70–80%	$0.182 \pm 0.058$	$143 \pm 4$	68/81
11.5 GeV			
0–5%	$0.748 \pm 0.014$	$115 \pm 2$	37/103
5–10%	$0.721 \pm 0.016$	$118 \pm 3$	33/106
10–20%	$0.684 \pm 0.016$	$121 \pm 3$	29/106
20–30%	$0.616 \pm 0.018$	$129 \pm 3$	29/107
30–40%	$0.566 \pm 0.022$	$134 \pm 3$	30/107
40–50%	$0.421 \pm 0.026$	$147 \pm 4$	58/104
50–60%	$0.359 \pm 0.029$	$151 \pm 4$	52/103
60–70%	$0.290 \pm 0.035$	$151 \pm 4$	66/99
70–80%	$0.248 \pm 0.036$	$149 \pm 4$	96/96

**Table A2.** *Cont.*

Centrality	$\rho_0$	$T$ (MeV)	$\chi^2/\text{NDF}$
19.6 GeV			
0–5%	$0.765 \pm 0.015$	$117 \pm 3$	25/105
5–10%	$0.740 \pm 0.015$	$120 \pm 3$	24/101
10–20%	$0.710 \pm 0.017$	$123 \pm 3$	24/101
20–30%	$0.649 \pm 0.019$	$129 \pm 3$	28/101
30–40%	$0.573 \pm 0.020$	$136 \pm 3$	35/102
40–50%	$0.507 \pm 0.022$	$141 \pm 4$	46/100
50–60%	$0.414 \pm 0.026$	$149 \pm 4$	59/100
60–70%	$0.350 \pm 0.026$	$152 \pm 3$	89/99
70–80%	$0.250 \pm 0.033$	$157 \pm 3$	133/99
27 GeV			
0–5%	$0.796 \pm 0.015$	$115 \pm 3$	42/98
5–10%	$0.771 \pm 0.016$	$118 \pm 3$	31/99
10–20%	$0.738 \pm 0.017$	$122 \pm 3$	33/99
20–30%	$0.683 \pm 0.018$	$128 \pm 3$	27/99
30–40%	$0.618 \pm 0.021$	$136 \pm 3$	23/99
40–50%	$0.534 \pm 0.024$	$144 \pm 4$	30/99
50–60%	$0.451 \pm 0.026$	$151 \pm 4$	37/99
60–70%	$0.352 \pm 0.030$	$158 \pm 4$	59/99
70–80%	$0.251 \pm 0.035$	$163 \pm 4$	114/99
39 GeV			
0–5%	$0.824 \pm 0.015$	$116 \pm 3$	23/99
5–10%	$0.785 \pm 0.015$	$121 \pm 3$	22/99
10–20%	$0.758 \pm 0.017$	$125 \pm 3$	20/99
20–30%	$0.709 \pm 0.018$	$130 \pm 3$	22/99
30–40%	$0.632 \pm 0.015$	$140 \pm 2$	17/99
40–50%	$0.567 \pm 0.024$	$145 \pm 4$	29/99
50–60%	$0.476 \pm 0.024$	$153 \pm 4$	43/99
60–70%	$0.379 \pm 0.031$	$164 \pm 4$	55/99
70–80%	$0.280 \pm 0.035$	$172 \pm 4$	97/99

**Table A3.** Femtoscopic  $\pi^\pm$  radii extrapolation to  $k_T = 0$  at  $\sqrt{s_{\text{NN}}} = 7.7, 11.5, 19.6, 27, 39$  GeV and the centrality classes from 0–5% to 60–70% using Equations (8)–(10). The radii are given in fm.

Centrality	$R_{\text{side}} (m_T^{-\alpha})$	$R_{\text{long}} (m_T^{-\alpha})$	$R_{\text{side}} (1/\sqrt{m_T})$	$R_{\text{long}} (1/\sqrt{m_T})$	$R_{\text{side}} (\text{BW})$	$R_{\text{long}} (\text{BW})$
7.7 GeV						
0–5%	$6.22 \pm 0.16$	$7.54 \pm 0.28$	$6.89 \pm 0.20$	$6.51 \pm 0.21$	$5.43 \pm 0.13$	$8.06 \pm 0.32$
5–10%	$5.58 \pm 0.13$	$7.50 \pm 0.28$	$6.54 \pm 0.18$	$6.18 \pm 0.19$	$5.11 \pm 0.11$	$7.69 \pm 0.29$
10–20%	$5.17 \pm 0.11$	$7.28 \pm 0.26$	$5.99 \pm 0.15$	$5.80 \pm 0.17$	$4.58 \pm 0.09$	$7.27 \pm 0.26$
20–30%	$4.61 \pm 0.09$	$6.68 \pm 0.22$	$5.38 \pm 0.12$	$5.32 \pm 0.14$	$4.08 \pm 0.07$	$6.66 \pm 0.22$
30–40%	$3.84 \pm 0.06$	$5.69 \pm 0.16$	$4.88 \pm 0.10$	$4.85 \pm 0.12$	$3.63 \pm 0.06$	$6.11 \pm 0.18$
40–50%	$3.31 \pm 0.05$	$5.35 \pm 0.14$	$4.37 \pm 0.08$	$4.37 \pm 0.09$	$3.17 \pm 0.04$	$5.53 \pm 0.15$
50–60%	$2.87 \pm 0.04$	$5.85 \pm 0.17$	$3.85 \pm 0.06$	$3.79 \pm 0.07$	$2.75 \pm 0.03$	$4.79 \pm 0.11$

**Table A3.** *Cont.*

Centrality	$R_{\text{side}} (m_T^{-\alpha})$	$R_{\text{long}} (m_T^{-\alpha})$	$R_{\text{side}} (1/\sqrt{m_T})$	$R_{\text{long}} (1/\sqrt{m_T})$	$R_{\text{side}} (\text{BW})$	$R_{\text{long}} (\text{BW})$
11.5 GeV						
0–5%	$5.83 \pm 0.14$	$8.37 \pm 0.34$	$6.89 \pm 0.20$	$7.01 \pm 0.24$	$5.42 \pm 0.12$	$8.73 \pm 0.37$
5–10%	$5.52 \pm 0.13$	$8.13 \pm 0.32$	$6.48 \pm 0.18$	$6.73 \pm 0.22$	$5.06 \pm 0.11$	$8.41 \pm 0.35$
10–20%	$5.03 \pm 0.11$	$7.54 \pm 0.28$	$6.03 \pm 0.15$	$6.32 \pm 0.20$	$4.65 \pm 0.09$	$7.92 \pm 0.31$
20–30%	$4.59 \pm 0.09$	$6.86 \pm 0.23$	$5.43 \pm 0.12$	$5.72 \pm 0.16$	$4.09 \pm 0.07$	$7.22 \pm 0.26$
30–40%	$4.01 \pm 0.07$	$6.26 \pm 0.19$	$4.90 \pm 0.10$	$5.17 \pm 0.13$	$3.63 \pm 0.06$	$6.55 \pm 0.21$
40–50%	$3.51 \pm 0.05$	$5.81 \pm 0.17$	$4.32 \pm 0.08$	$4.65 \pm 0.11$	$3.06 \pm 0.04$	$5.97 \pm 0.17$
50–60%	$2.98 \pm 0.04$	$5.17 \pm 0.13$	$3.77 \pm 0.06$	$4.18 \pm 0.09$	$2.62 \pm 0.03$	$5.38 \pm 0.14$
19.6 GeV						
0–5%	$5.88 \pm 0.15$	$8.94 \pm 0.39$	$6.96 \pm 0.21$	$7.53 \pm 0.28$	$5.50 \pm 0.13$	$9.38 \pm 0.43$
5–10%	$5.52 \pm 0.13$	$8.51 \pm 0.36$	$6.57 \pm 0.18$	$7.19 \pm 0.25$	$5.16 \pm 0.11$	$8.97 \pm 0.39$
10–20%	$5.13 \pm 0.11$	$7.99 \pm 0.31$	$6.12 \pm 0.16$	$6.72 \pm 0.22$	$4.75 \pm 0.10$	$8.44 \pm 0.35$
20–30%	$4.57 \pm 0.09$	$7.33 \pm 0.26$	$5.52 \pm 0.13$	$6.05 \pm 0.18$	$4.20 \pm 0.08$	$7.65 \pm 0.29$
30–40%	$4.00 \pm 0.07$	$6.47 \pm 0.21$	$4.95 \pm 0.10$	$5.47 \pm 0.15$	$3.70 \pm 0.06$	$6.93 \pm 0.24$
40–50%	$3.47 \pm 0.05$	$5.96 \pm 0.17$	$4.36 \pm 0.08$	$4.90 \pm 0.12$	$3.20 \pm 0.04$	$6.23 \pm 0.19$
50–60%	$2.87 \pm 0.04$	$5.00 \pm 0.12$	$3.87 \pm 0.06$	$4.34 \pm 0.09$	$2.76 \pm 0.03$	$5.54 \pm 0.15$
60–70%	$2.62 \pm 0.03$	$4.52 \pm 0.10$	$3.34 \pm 0.05$	$3.81 \pm 0.07$	$2.33 \pm 0.02$	$4.87 \pm 0.12$
27 GeV						
0–5%	$5.94 \pm 0.15$	$9.17 \pm 0.41$	$7.04 \pm 0.21$	$7.77 \pm 0.30$	$5.61 \pm 0.13$	$9.67 \pm 0.46$
5–10%	$5.65 \pm 0.14$	$8.79 \pm 0.38$	$6.66 \pm 0.19$	$7.42 \pm 0.27$	$5.27 \pm 0.12$	$9.26 \pm 0.42$
10–20%	$5.14 \pm 0.11$	$8.14 \pm 0.33$	$6.23 \pm 0.16$	$6.90 \pm 0.23$	$4.86 \pm 0.10$	$8.65 \pm 0.37$
20–30%	$4.63 \pm 0.09$	$7.46 \pm 0.27$	$5.60 \pm 0.13$	$6.26 \pm 0.19$	$4.31 \pm 0.08$	$7.88 \pm 0.30$
30–40%	$4.07 \pm 0.07$	$6.79 \pm 0.23$	$5.02 \pm 0.11$	$5.63 \pm 0.16$	$3.79 \pm 0.06$	$7.14 \pm 0.25$
40–50%	$3.55 \pm 0.05$	$6.10 \pm 0.18$	$4.47 \pm 0.08$	$5.03 \pm 0.12$	$3.29 \pm 0.05$	$6.41 \pm 0.20$
50–60%	$2.99 \pm 0.04$	$5.27 \pm 0.14$	$3.91 \pm 0.07$	$4.46 \pm 0.10$	$2.81 \pm 0.03$	$5.72 \pm 0.16$
60–70%	$2.61 \pm 0.03$	$4.63 \pm 0.11$	$3.40 \pm 0.05$	$3.88 \pm 0.07$	$2.37 \pm 0.02$	$4.99 \pm 0.12$
39 GeV						
0–5%	$6.11 \pm 0.16$	$9.32 \pm 0.43$	$7.18 \pm 0.22$	$7.96 \pm 0.31$	$5.71 \pm 0.14$	$9.96 \pm 0.49$
5–10%	$5.74 \pm 0.14$	$8.76 \pm 0.38$	$6.81 \pm 0.20$	$7.56 \pm 0.28$	$5.36 \pm 0.12$	$9.50 \pm 0.44$
10–20%	$5.18 \pm 0.11$	$8.50 \pm 0.35$	$6.28 \pm 0.17$	$7.04 \pm 0.24$	$4.91 \pm 0.10$	$8.88 \pm 0.39$
20–30%	$4.68 \pm 0.09$	$7.75 \pm 0.30$	$5.69 \pm 0.14$	$6.38 \pm 0.20$	$4.39 \pm 0.08$	$8.07 \pm 0.32$
30–40%	$4.14 \pm 0.07$	$6.97 \pm 0.24$	$5.10 \pm 0.11$	$5.76 \pm 0.16$	$3.84 \pm 0.06$	$7.35 \pm 0.26$
40–50%	$3.61 \pm 0.06$	$6.03 \pm 0.18$	$4.53 \pm 0.09$	$5.10 \pm 0.13$	$3.35 \pm 0.05$	$6.52 \pm 0.21$
50–60%	$3.11 \pm 0.04$	$5.39 \pm 0.14$	$3.97 \pm 0.07$	$4.52 \pm 0.10$	$2.86 \pm 0.04$	$5.80 \pm 0.17$
60–70%	$2.73 \pm 0.03$	$4.78 \pm 0.11$	$3.45 \pm 0.05$	$4.03 \pm 0.08$	$2.42 \pm 0.03$	$5.21 \pm 0.13$

**Table A4.** Comparison of volumes at  $\sqrt{s_{\text{NN}}} = 7.7, 11.5, 19.6, 27, 39$  GeV and the centrality classes from 0–5% to 60–70% obtained using Tsallis-3 statistics and calculated using Equations (8)–(10). The volumes are given in  $\text{fm}^3$ .

Centrality	Tsallis-3	$m_T^{-\alpha}$	$1/\sqrt{m_T}$	BW
7.7 GeV				
0–5%	$13,020 \pm 1454$	$13,674 \pm 2022$	$14,470 \pm 2137$	$11,145 \pm 1646$
5–10%	$11,184 \pm 869$	$10,931 \pm 1618$	$12,370 \pm 1826$	$9394 \pm 1387$
10–20%	$12,171 \pm 1279$	$9125 \pm 1349$	$9750 \pm 1440$	$7150 \pm 1056$
20–30%	$9350 \pm 338$	$6657 \pm 985$	$7213 \pm 1065$	$5187 \pm 766$
30–40%	$5956 \pm 537$	$3934 \pm 582$	$5418 \pm 800$	$3766 \pm 556$
40–50%	$5161 \pm 600$	$2734 \pm 405$	$3911 \pm 578$	$2605 \pm 385$
50–60%	$2782 \pm 189$	$2254 \pm 336$	$2632 \pm 389$	$1691 \pm 251$

**Table A4.** *Cont.*

Centrality	Tsallis-3	$m_T^{-\alpha}$	$1/\sqrt{m_T}$	BW
11.5 GeV				
0–5%	19,335 ± 1589	13,329 ± 1969	15,570 ± 2299	12,001 ± 1772
5–10%	15,199 ± 505	11,611 ± 1716	13,229 ± 1953	10,079 ± 1488
10–20%	14,691 ± 1562	8951 ± 1322	10,766 ± 1590	8022 ± 1185
20–30%	11,556 ± 1170	6775 ± 1001	7883 ± 1164	5654 ± 835
30–40%	9780 ± 1035	4717 ± 697	5800 ± 856	4051 ± 598
40–50%	5656 ± 450	3346 ± 494	4072 ± 601	2616 ± 386
50–60%	3410 ± 267	2152 ± 319	2781 ± 411	1730 ± 256
19.6 GeV				
0–5%	32,249 ± 2714	14,479 ± 2138	17,107 ± 2526	13,294 ± 1963
5–10%	30,633 ± 1803	12,157 ± 1796	14,506 ± 2142	11,171 ± 1649
10–20%	23,912 ± 2270	9844 ± 1454	11,794 ± 1741	8920 ± 1317
20–30%	19,404 ± 1545	7155 ± 1057	8628 ± 1274	6323 ± 934
30–40%	11,305 ± 894	4844 ± 715	6277 ± 927	4434 ± 655
40–50%	6582 ± 470	3355 ± 496	4369 ± 645	2981 ± 440
50–60%	4010 ± 273	1929 ± 285	3037 ± 448	1977 ± 292
60–70%	2127 ± 144	1451 ± 215	1987 ± 293	1242 ± 184
27 GeV				
0–5%	27,576 ± 2472	15,125 ± 2233	18,024 ± 2661	14,229 ± 2101
5–10%	26,184 ± 551	13,118 ± 1937	15,403 ± 2274	12,016 ± 1774
10–20%	26,783 ± 401	10,061 ± 1486	12,529 ± 1850	9586 ± 1415
20–30%	22,091 ± 2252	7482 ± 1105	9195 ± 1358	6870 ± 1014
30–40%	13,673 ± 486	5258 ± 776	6642 ± 981	4797 ± 708
40–50%	8441 ± 636	3591 ± 530	4695 ± 693	3251 ± 480
50–60%	4618 ± 302	2201 ± 325	3193 ± 472	2112 ± 312
60–70%	2381 ± 155	1476 ± 218	2095 ± 309	1314 ± 194
39 GeV				
0–5%	27,434 ± 696	16,276 ± 2418	19,200 ± 2835	15,215 ± 2247
5–10%	26,424 ± 556	13,523 ± 1998	16,435 ± 2427	12,771 ± 1886
10–20%	25,733 ± 225	10,695 ± 1580	13,003 ± 1920	10,024 ± 1480
20–30%	23,522 ± 585	7944 ± 1173	9659 ± 1426	7281 ± 1075
30–40%	18,861 ± 781	5581 ± 825	7023 ± 1037	5076 ± 750
40–50%	8612 ± 660	3666 ± 542	4896 ± 723	3424 ± 506
50–60%	6334 ± 213	2445 ± 362	3335 ± 493	2221 ± 328
60–70%	2537 ± 173	1668 ± 247	2248 ± 332	1425 ± 211

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