Laser Heating Method for an Alkali Metal Atomic Cell with Heat Transfer Enhancement

Yang Li, Guoqing Zhou, Shencheng Tian, Xuejing Liu *, Xiangmei Dong and Xiumin Gao

School of Optical-Electrical and Computer Engineering, University of Shanghai for Science and Technology, Shanghai 200093, China
* Correspondence: liuxuejing@usst.edu.cn

Abstract: Alkali metal atomic cells are crucial components of atomic instruments, such as atomic magnetometers, atomic gyroscopes, and atomic clocks. A highly uniform and stable heating structure can ensure the stability of the alkali metal atom density. The vapor cell of an atomic magnetometer that uses laser heating has no magnetic field interference and ease of miniaturization, making it superior to hot air heating and AC electric heating. However, the current laser heating structure suffers from low heating efficiency and uneven temperature distribution inside the vapor cell. In this paper, we designed a non-magnetic heating structure based on the laser heating principle. We studied the temperature distribution of the heating structure using the finite element method (FEM) and analyzed the conversion and transfer of laser energy. We found that the heat conduction between the vapor cell and the heating chips (colored filters) is poor, resulting in uneven temperature distribution and low heating efficiency in the vapor cell. Therefore, the addition of graphite film to the four surfaces of the vapor cell was an important improvement. This addition helped to balance the temperature distribution and improve the conduction efficiency of the heating structure. It was measured that the power of the heating laser remained unchanged. After the addition of the graphite film, the temperature difference coefficient (CVT) used to evaluate the internal temperature uniformity of the vapor cell was reduced from 0.1308 to 0.0426. This research paper is crucial for improving the heating efficiency of the non-magnetic heating structure and the temperature uniformity of the vapor cell.

Keywords: SERF magnetometer; non-magnetic laser heating; graphite film
additional magnetic noise is introduced, its heating efficiency is low, and the overall structure is too large to be integrated. Electric heating can use a lock-in amplifier to filter out additional magnetic noise in the signal, but the working frequency must remove the heating current frequency, which limits the measurement range of the atomic magnetometer. More importantly, this method cannot completely filter out residual magnetic noise. Although intermittent electric heating can completely avoid additional magnetic noise, the atomic magnetometer does not work during the heating process, and it cannot measure the continuous changes in the magnetic field, directly leading to a decrease in the sensitivity of the atomic magnetometer. In contrast, laser heating is a good option compared to other heating technologies because it does not generate additional magnetic noise. Furthermore, as atomic magnetometers continue to miniaturize and integrate, the energy required to heat chip-level atomic cells to the target temperature is very low. The continuous optimization and updating of the laser heating structure also promote laser heating as the mainstream heating method for alkali metal cells in the future. Therefore, the non-magnetic heating method chosen in this study is laser heating. Preusser et al. [23] used a 915 nm semiconductor laser to directly heat the vapor cell, which was heated to 97 °C with a required laser power of 200 mW. The laser energy cannot be completely absorbed by the vapor cell, and the heating efficiency is limited. Mhaskar et al. [24] and Sheng [25] installed colored filters on both sides of the vapor cell to improve the laser absorption efficiency, but there is still residual energy that has not been utilized. In Sheng’s study, two 1550 nm laser beams were used simultaneously to heat the vapor cell, which made the heating structure more complex. Therefore, improving the laser heating efficiency and the spatial uniformity of the temperature inside the vapor cell while simplifying the heating structure is also the key to improving laser heating technology.

In this paper, we study a heating structure for a cell that utilizes laser heating technology, which is to improve the heating efficiency and temperature uniformity of the vapor cell. First, a finite element (FEM) model is established to study the internal temperature of the vapor cell and analyze the laser energy conversion and transfer. To solve the problems of low heating efficiency and uneven temperature distribution inside the vapor cell caused by the poor thermal conductivity of the vapor cell and the heating chips (colored filters), we added a graphite film to the heating structure. To validate the proposed method, the operating power of the heating laser is kept constant, and graphite films of different sizes are added to the heating structure. This is performed to measure the temperature variation inside the vapor cell. The CVT and the temperature error distribution plots of the vapor cell are used to characterize the uniformity of the temperature distribution inside the vapor cell.

2. Analysis and Simulation Methods

2.1. Heat Transfer Analysis

The heating process involves three types of heat transfer: thermal radiation, thermal convection, and thermal conduction. The heating chips convert laser energy to thermal energy and transfer it evenly to the outer wall of the vapor cell through heat conduction. The outer wall of the vapor cell then radiates heat to the gas inside the vapor cell. Simultaneously, the colored filters and the vapor cell emit thermal energy to the surrounding environment through thermal radiation. Eventually, the entire system reaches a dynamic equilibrium state, where heat conduction can be expressed by Fourier’s law of heat conduction as follows [26]:

\[
\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} = \frac{\rho \cdot C_p}{k} \cdot \frac{\partial T}{\partial t}.
\] (1)

In Equation (1), \(T\) is temperature, \(t\) is time, \(\rho\) is the density of the object, \(C_p\) is constant pressure heat capacity, \(k\) is the thermal conductivity. The left side of the equation shows the
rate of heat transfer, so heat transfer is only affected by the material properties of the object. The heat radiation can be expressed by the Stefan–Boltzmann equation as follows:

\[ q_1 = \varepsilon \sigma A \left( T_{\text{amd}}^4 - T^4 \right), \]  

where \( q_1 \) is the heat flow rate, \( \varepsilon \) is the rate of radiation, \( \sigma \) is Stefan–Boltzmann constant, \( A \) is the area of radiation surface, \( T_{\text{amd}} \) is temperature of the vapor cell’s wall, and \( T \) is the temperature of the alkali metal gas in the vapor cell. The heat convection can be expressed by Newton cooling equation as follows:

\[ q_2 = h A_1 (T_{\text{Solid}} - T_{\text{Fluid}}), \]  

where \( h \) is the heat transfer coefficient, \( A_1 \) is the area of the radiating surface, \( T_{\text{Solid}} \) is the temperature of solid surface, and \( T_{\text{Fluid}} \) is the fluid temperature around the solid. When the system finally reaches dynamic stability, the thermal balance equation is obtained as follows:

\[ Q = \rho C_p \frac{\partial T}{\partial t} + \rho C_p v \cdot \nabla T + \nabla \cdot (-k \cdot \nabla T), \]  

where \( Q \) is the total heat input; \( v \) is the velocity field in the fluid; \( \nabla \) is the Nabla operator, and it is defined as a differential operator; it also has vector form, so \( \nabla T \) is the changes of temperature in all directions in space; and \( \rho \) is the density of the object. The heating chips and the vapor cell are closely bonded and placed in the insulating layer. When the system reaches a stable state, the primary heat transfer mode is conduction, and convection and radiation can be disregarded. In this study, both the vapor cell and the heating chips are made of quartz glass, which has low thermal conductivity and high thermal capacity at constant pressure. Therefore, the thermal conductivity of the system is poor. To improve the thermal conductivity of the system, the graphene thin layers are added to the heating structure. As a material with high thermal conductivity, graphite sheets have a lamellar structure with anisotropy [27–29]. The thermal conductivity is high along the lamellar direction and low in the vertical direction.

### 2.2. Finite Element Analysis

Based on the COMSOL Multiphysics simulation software, the temperature distribution of a cell heated by a pair of heating chips on the X-axis is studied using the finite element analysis method. The required physical field is heat transfer in solid. The simulation process consists of four steps. First, the heating structure model is built, including heating chips, the vapor cell, and graphite films. Second, material properties are set, and physical fields are added for solid and fluid heat transfer. The heating chip heats the vapor cell by absorbing the energy of the laser, so the boundary of the heating chip is set as the heat source in the way of heat flux. The energy of the heating laser is a continuous laser with a Gaussian spatial distribution (GSD). The laser source is defined as follows [30]:

\[ Q = GSD \times \rho_{\text{laser}}. \]  

\( Q \) is the heat produced by the laser; \( \rho_{\text{laser}} \) is the energy density of the laser, expressed as \( 2 \cdot P_{\text{laser}} / D \); \( P_{\text{laser}} \) is the working power of the laser; and \( D \) is the area of the laser spot. Gaussian spatial distribution (GSD) is represented by Equation (6):

\[ GSD = e^{-2 \cdot r_{\text{focus}}^2 / r_{\text{spot}}^2}, \]  

\( r_{\text{spot}} \) is the radius of the laser spot. \( r_{\text{focus}} \) is the distance to the center of the laser spot, which is expressed as \( r_{\text{focus}} = \sqrt{(x - x_{\text{focus}})^2 + (y - y_{\text{focus}})^2 + (z - z_{\text{focus}})^2}; \) \( (x_{\text{focus}}, y_{\text{focus}}, z_{\text{focus}}) \) represents the coordinates of the center point of the laser source; the radius of the laser spot is set to 4 mm; and the initial temperature of the environment
is set to 20 °C. In order to improve the calculation precision, the whole heating model is partitioned by grid. The sizes of the grid units of the heating chips and the vapor cell are, respectively, set to fine and extremely fine. Finally, the transient temperature of the heating structure is analyzed, and the temperature distribution when temperature of the vapor cell reaches dynamic stability is analyzed.

3. Experimental Setup

The rubidium SERF magnetometer comprise a cell, a pump laser, a detection laser, a magnetically shielded barrel, and a non-magnetic heating structure. The non-magnetic heating structure is responsible for heating the rubidium vapor to operational temperature without creating any magnetic noise during the heating process. In this paper, we selected a laser to heat the rubidium gas cell. As illustrated in Figure 1, the heating structure consists of the following components: a 1550 nm heating laser, two heating chips (colored filters), a temperature sensor (PT100), a cell, an insulation layer, a graphite film, and a thermal imaging camera with a real-time 160 × 120-pixel thermal map.

In this experiment, colored filters from SCHOTT were utilized as heating chips to absorb laser energy in the 1550 nm wavelength range. The incident and exit window filters for the heating laser were 1 mm UG-5 and 3 mm RG-9, respectively; the incident optical power absorbed by the two filters was 56.3% and 100%, allowing each filter to absorb approximately equal amounts of heating laser energy; and the converted heat was then used to heat the vapor cell. The temperature and temperature distribution of the vapor cell were measured via a temperature sensor (PT100) and thermal imager, respectively, the rubidium atomic vapor and buffer gas are enclosed in a 20 mm × 20 mm × 20 mm cubic alkali metal cell. An insulating layer is incorporated to reduce heat loss and enhance heat transfer efficiency; the cube-shaped insulating layer with a side length of 40 mm is made of polyimide foam, which has a thermal conductivity of 0.035 W/(m·K). Two cylindrical holes with a diameter of 10 mm intersect with the cube, allowing for the passage of lasers that irradiate the vapor cell.

After conducting the experiment, it was found that the thermal conductivity of the vapor cell and heating was poor. Therefore, a graphite film is added to the heating structure as a high thermal conductive layer to balance the temperature distribution and improve the conduction efficiency. The thickness of the graphite film is 0.2 mm. The thermal conductivity along the fiber axis is 800 W/(m·K), and the thermal conductivity perpendicular to the fiber axis is 10 W/(m·K). The graphene thin layer, as shown in Figure 2, is a square sheet of 20 mm × 80 mm in size. This sheet contains two circular holes of 10 mm diameter each for light passage and one circular hole of 6 mm diameter for passage through the vapor cell shank.
The ambient temperature is 20 °C, the laser power is set to 2.4 W, and the diameter of the heated laser spot is 8 mm. Heating time can be reduced via a pair of heating chips to heat the vapor cell simultaneously. The temperature gradient of the vapor cell is smaller when a heater is placed in each of the two windows of the heating laser compared to placing a heater in only one window. Using the finite element method to simulate the heat transfer of the heating structure with added graphite film, Figure 4a shows the temperature distribution of the two heating chips when heating the vapor cell.
When graphite film is added to the four sides of the vapor cell, the operating power of the 1550 nm heating laser is 2.4 W. It can be seen from Figure 5 that when the power of the heating laser is 2.4 W, the simulated interior temperature of the vapor cell is 120 °C, while the experimental temperature of the vapor cell is 118 °C. This is because the thermal insulation structure cannot completely eliminate the vapor cell and the outside thermal radiation, resulting in heat loss of the vapor cell, so there is a deviation between theoretical simulation and experimental data.

4.2. Experiments

To study the uniformity of the vapor cell under laser heating, we measured the internal temperature using a temperature sensor (PT100) and a thermal imager. The temperature sensor (PT100) was attached to the upper center of the vapor cell, and the measurement results are shown in Figure 5.

We define the center point of the vapor cell as A, and the space coordinates of point A is (0,0,10), and two heating chips are distributed at 10 mm and −10 mm of the x axis. Figure 4b shows the temperature distribution of the heating structure on the XAY, YAZ, and XAZ planes. L1, L2, and L3 represent three straight lines passing through the x, y, and z axes of the central point inside the vapor cell. In contrast to the temperature graph in Figure 4c, the temperature gradient in the direction of the X-axis is larger. This is because the perforations of the insulating layer are oriented in the direction of both the x-axis and the y-axis, and the two heating chips are also located in the direction of the x-axis. Furthermore, the heating chip located in the incident window of the heating laser absorbed 12.6% more laser energy than the heating chip located in the outgoing window of the heating laser. This is the main reason for the large temperature gradient along the X-axis. The smaller temperature gradient in the Y-axis direction is due to the smaller aperture and less thermal exchange between the vapor cell and the external environment. In the Z-axis direction, the vapor cell’s temperature gradient is the smallest, and the outer wall temperature of the vapor cell is close to the interior temperature. Therefore, to reduce the temperature measurement error, the temperature sensor (PT100) can be placed in the upper shell center of the vapor cell to measure its temperature.

Figure 4. (a) The temperature distribution diagram of the heating structure when the temperature in the vapor cell is around 120 °C; (b) Temperature distribution on the XAY, YAZ, and XAZ planes; (c) Taking the center of the vapor cell as the coordinate origin, the temperature curve along the three axes.
Figure 5. Simulated (red) and experimental (black) temperature curves of the vapor cell were measured.

Figure 6 is a thermal image of the internal temperature radiation of the vapor cell with the Y-axis facing. This is when the graphite film is not added to the surface of the vapor cell, and the graphite film with different sizes is added.

Figure 6. Temperature radiography was performed on the vapor cell with different sizes of graphite films added. (a) the vapor cell without any graphite films; (b) the vapor cell with graphite films on two surfaces; (c) the vapor cell with graphite films on three surfaces; (d) the vapor cell with graphite films on four surfaces.

From Figure 6, it is evident that the temperature near the laser incident window in the vapor cell is the highest. This is because the heating chip located in the exit window absorbed too much laser energy, resulting in the two colored filters absorbing different amounts of laser energy. Therefore, one can select either a thin heating chip or a heating chip with a low absorption rate to the heating laser. Figure 6a shows the temperature radiography of the vapor cell without the addition of a graphite film. The material of the vapor cell and the heating chips is quartz glass, which has a poor thermal conductivity at 1.4 W/(m K). This leads to a significant accumulation of heat in the incident window of the laser. The thermal energy inside the vapor cell comes from the radiant heat of the vapor cell wall. However, there is a shank at the bottom of the vapor cell, which increases the radiant distance at the bottom of the vapor cell. As a result, the temperature of the bottom side is lower than the upper side. In the experiment, when the laser power reached...
600 mW, the UG-5 heating chip broke due to the large difference in temperature between the center and the surrounding areas. In Figure 6b, the graphite film was pasted on the two surfaces of the vapor cell that were in contact with the heating chips, which obviously improved the uniformity of the temperature distribution of the incident window. However, the temperature of the exit window remained low, which was due to the poor thermal conductivity of the vapor cell. In Figure 6c, we have pasted graphite films on the upper surface of the vapor cell to increase heat transfer along the X-axis. As can be seen from the figure, the heat from the incident window is transferred to the exit window by the graphite film, improving the temperature uniformity of the upper part of the vapor cell.

In order to improve the temperature uniformity of the vapor cell, the graphite film covers the upper and lower top surfaces of the vapor cell, the incident light window, and the outgoing light window at the same time. Under the same heating condition, the thermal radiation Figure 6d is obtained. By comparing the temperature radiation in Figure 6, it can be concluded that the graphite film improves the thermal conductivity of the system and makes the temperature distribution of the vapor cell more uniform. However, the thermal radiation between the vapor cell and the outside also increased, and the overall temperature of the vapor cell decreased.

Converting the temperature radiography of the vapor cell in Figure 6 to the temperature error analysis diagram shown in Figure 7 in order to visualize the temperature uniformity of the vapor cell.

![Temperature error profiles](image)

**Figure 7.** Temperature error profiles were measured for the vapor cell with different graphite films being added. (a) The vapor cell without graphite films; (b) the vapor cell with graphite film added to two surfaces; (c) the vapor cell with graphite films added to three surfaces; (d) the vapor cell with graphite films added to four surfaces.

It can be seen from Figure 7 that the graphite film can improve the thermal conductivity of the heating structure and balance the temperature distribution of the vapor cell. However, due to the heating chip of the incident window absorbing too much laser energy, the temperature of the incident window is higher. The reason why the two surfaces in the Y-axis direction of the vapor cell are not chosen to stick graphite films is to capture the temperature radiation pattern of the vapor cell.

Figures 6 and 7 show the influence of graphite film on the temperature uniformity of the vapor cell. Finally, we introduce the temperature difference coefficient to evaluate the temperature uniformity of the vapor cell.
The expression for the CVT is as follows:

\[
CVT = \frac{f}{T_{\text{aver}}},
\]

(7)

where \( T_{\text{aver}} \) is the average temperature, and \( f \) is the unbiased variance of the temperature. \( T_{\text{aver}} \) denotes \( \sum T_i / S \), where \( T_i \) is the temperature of each plane point, and \( S \) is the surface area; \( f \) is denoted by \( \sqrt{\sum (T_i - T_{\text{aver}})^2 / S} \). According to Equation (7), the smaller the CVT is, the more uniform the temperature distribution is. The measured temperature radiography was converted into a temperature matrix to obtain the CVT values, and the different CVTs were obtained, as shown in Table 1.

Table 1. The CVT value of the vapor cell changes when graphite films of different sizes are applied to the outer wall of the vapor cell.

<table>
<thead>
<tr>
<th>Size</th>
<th>CVT</th>
</tr>
</thead>
<tbody>
<tr>
<td>0</td>
<td>0.1308</td>
</tr>
<tr>
<td>20 mm × 20 mm (two pieces)</td>
<td>0.1038</td>
</tr>
<tr>
<td>20 mm × 60 mm</td>
<td>0.0697</td>
</tr>
<tr>
<td>20 mm × 80 mm</td>
<td>0.0426</td>
</tr>
</tbody>
</table>

From Table 1, as the size of the graphite film increases, the calculated CVT of the vapor cell decreases, and the temperature uniformity of the vapor cell improves. Therefore, the temperature uniformity of the heating structure is greatly improved by adding graphite films.

5. Conclusions

In this paper, a vapor cell heating structure based on laser heating technology was constructed, and the spatial distribution of temperature was optimized via graphite films. By constructing the finite element model, the numerical analysis of the thermal transfer in the heating system shows that the uniformity of the temperature in the vapor cell is limited because the heat transfer is poor between the vapor cell and the heating chips. The laser power is constant, and the inner temperature CVT of the vapor cell decreased by 67% after adding graphite film on the four sides of the vapor cell, so the structure of graphite film can efficiently improve the temperature distribution homogeneity in the vapor cell. In the next step, we will study the action mechanism between the temperature homogeneity and the magnetometer output signal.

Author Contributions: Conceptualization, Y.L. and G.Z.; methodology, G.Z.; software, Y.L.; validation, Y.L., G.Z. and S.T.; formal analysis, Y.L.; investigation, G.Z.; resources, X.G.; data curation, Y.L.; writing—original draft preparation, X.L.; writing—review and editing, G.Z.; visualization, S.T.; supervision, S.T.; project administration, X.D.; funding acquisition, X.G. All authors have read and agreed to the published version of the manuscript.

Funding: This research was funded by the National Natural Science Foundation of China (NSFC), grant number 62005167.

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

Data Availability Statement: Not applicable.

Conflicts of Interest: The authors declare no conflict of interest.
References


8. Perry, A.R.; Bulatowicz, M.D.; Larsen, M.; Walker, T.G.; Wyllie, R. All-optical intrinsic atomic gradiometer with sub-20 fT/cm/√Hz sensitivity in a 22 mT earth-scale magnetic field. *Opt. Express* 2020, 28, 36696–36705. [CrossRef]


15. Lu, J.X.; Wang, J.; Yang, K.; Zhao, J.P.; Quan, W.; Han, B.C.; Ding, M. In-Situ Measurement of Electrical-Heating-Induced Magnetic Field for an Atomic Magnetometer. *Sensors* 2020, 20, 1826. [CrossRef]


**Disclaimer/Publisher’s Note:** The statements, opinions and data contained in all publications are solely those of the individual author(s) and contributor(s) and not of MDPI and/or the editor(s). MDPI and/or the editor(s) disclaim responsibility for any injury to people or property resulting from any ideas, methods, instructions or products referred to in the content.