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Multiple Alternatives of Offset Boosting in a Symmetric Hyperchaotic Map

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Abstract: The offset as the average value of a variable plays an important role in signal processing and system design. Offset boosting can be realized by a non-bifurcation parameter or an initial condition. In this work, symmetric coexisting attractors with opposite polarity and a 2D hyperchaotic map with multiple modes of offset boosting are proposed, where the offset can be controlled both by the initial condition and system parameter, and as a result, multiple alternatives of offset boosting and offset competition show up. Consequently, the final offset is determined eventually by the balance of two factors. The theoretical findings are verified through the hardware experiment based on the STM32. Finally, a pseudo-random number generator (PRNG) is constructed based on the newly proposed hyperchaotic map, demonstrating its high performance in engineering applications.

Keywords: hyperchaotic map; symmetric coexisting attractors; offset boosting; offset boosting regime



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1. Introduction

Chaos control, including amplitude control [1–3] and offset boosting [4–7], is a hot topic in the nonlinear community, which becomes more important for a multistable system since it hosts many coexisting attractors [8–12]. Offset boosting of a chaotic attractor brings great convenience for chaotic signal conditioning and thus brings increasing interest. The introduction of a periodic function drives attractor self-reproduction, where any of the coexisting attractors locates at a definite position in phase space and can be extracted out by an initial condition [13]. A periodically forced megastable chaotic system may give birth to countless embedded coexisting attractors [14]. In Chua's memristive multi-double-scroll system, the internal parameter of the memristor can be employed to construct the desired number of multi-double-scroll chaotic attractors [15]. All the phenomena resort to the technology of offset boosting. However, the associated research on offset control in the field of the discrete map is relatively poor, which leaves a significant margin [16–18]. It has been proved that offset boosting plays an important role in the discrete map for the desired multistability [19,20]. Zhou et al. proposed homogenous multistability in a 3D map, in which the chaotic attractor structure and the distance between any two petals can be controlled by initial conditions [21]. Offset boosting can be applied to construct a discrete multi-cavity [22]. A self-reproduction hyperchaotic map may exhibit compound lattice dynamics associated with offset boosting [23]. Offset boosting shows its specific function in discrete systems for attractor control.

Offset boosting can be realized by a non-bifurcation parameter or an initial condition. Therefore, sometimes it is difficult to determine how exactly offset boosting happens since the final derived offset may be the balanced result from the competition of a parameter and an initial start. In this work, a 2D hyperchaotic map with symmetric coexisting attractors is designed, where hyperchaotic sequences with opposite polarities are obtained, and multiple

modes of offset boosting controlled both by the initial condition and system parameter are found. As a result, multiple alternatives of offset boosting and offset competition show up. The main contributions are summarized as follows:

- (1) A new control knob is found, which can control the interval and direction of attractor self-reproduction.
- (2) Various alternatives of offset boosting are systematically explored, which include initially controlled offset boosting, parameter-oriented offset boosting and competitive offset boosting.
- (3) Competitive offset boosting is firstly discussed, where the newly introduced constants show the function of offset boosting determining the direction of the shifted attractor in the phase space, and as a result, based on the constant, the bipolar and unipolar chaotic signals can be switched accordingly.
- (4) Periodic windows are caught by the plot of the bifurcation diagram, which poses a great threat to engineering applications. However, competitive offset boosting can be easily applied to cross this interval and reach hyperchaos.
- (5) An STM32-based circuit implementation is constructed to prove two different offset boosting regimes. PRNG is employed to demonstrate the practical application of the proposed map.

The rest of this paper is organized as follows. A basic analysis of the 2D hyperchaotic map, including stability of fixed points and bifurcation are provided in Section 2. In Section 3, various mechanisms of offset boosting are exhaustively explored. In Section 4, STM32-based circuit implementation is set up for physical verification, and PRNG is derived to explore its chaos-based applications. Discussions and conclusions are presented in Section 5.

2. A 2D Hyperchaotic Map

2.1. Model Description

By introducing a periodic nonlinearity [23], a novel 2-D discrete map is proposed,

$$\begin{cases} x_{n+1} = x_n + \sin(x_n + y_n), \\ y_{n+1} = ax_n, \end{cases} \quad (1)$$

where n represents a natural number, x_n and y_n respectively denote the n th states. System parameters are represented by the symbol a . When $x_{n+1} \rightarrow -x_{n+1}$, $y_{n+1} \rightarrow -y_{n+1}$, $x_n \rightarrow -x_n$, $y_n \rightarrow -y_n$, map (1) keeps its polarity balance indicating that map (1) is of inversion symmetry.

2.2. Bifurcation Analysis

The Lyapunov exponent is a direct theoretical description and widely used criteria indicating the existence of chaos, which is widely accepted in nonlinearity community. The Lyapunov exponent of a differential equation $x_{i+1} = f(x_i)$, denoted as $\lambda_f(x)$, is mathematically defined as

$$\lambda_{f(x)} = \lim_{n \rightarrow \infty} \left\{ \frac{1}{n} \ln \left| \frac{f^n(x_0 + \varepsilon)}{\varepsilon} \right| \right\}, \quad (2)$$

where ε is a very small positive value nearing zero. One positive value indicates chaos, and more positive values show hyperchaos.

When $(x_0, y_0) = (1, 0)$ and a varies in $(-3.4, 2.5)$, several kinds of evolution including period, chaos, and hyperchaos are found in map (1) [24–26], as depicted in Figure 1. When a is in $(-2.7, -2.248)$, hyperchaos shows up; when a varies in $(-2.248, -2.03)$, map (1) is chaotic; when a increases in $(-2.03, -2.01)$, a small range periodic window is captured; when a increases in $(1.758, 1.882)$, quasi-periodic oscillation behavior is exhibited; when the parameter a varies in $(1.892, 2.5)$, map (1) is chaotic, and two separate periodic windows $(1.99, 2.115)$ and $(2.333, 2.365)$ are embedded. Typical dynamical behaviors of map (1) are summarized in Table 1, corresponding phase orbits are shown in Figure 2. When $x_0 \rightarrow -x_0$,

$y_0 \rightarrow -y_0$, the polarity in map (1) is switched and its symmetrical attractors are produced, as shown in Figure 2d–f, their symmetrical waveforms are displayed in Figure 3.

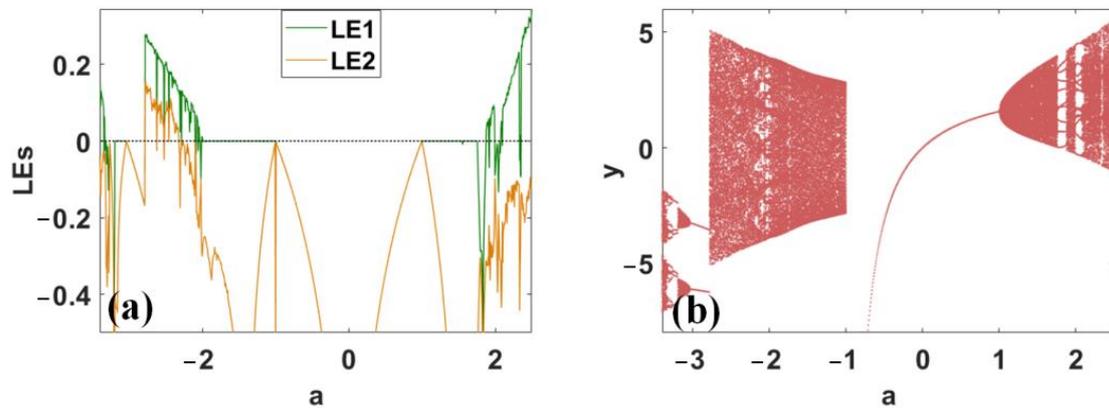


Figure 1. Dynamical evolution of hyperchaotic map (1) under the initial condition $(x_0, y_0) = (1, 0)$ when a varies in $(-3.4, 2.5)$: (a) Lyapunov exponents, (b) bifurcation diagram.

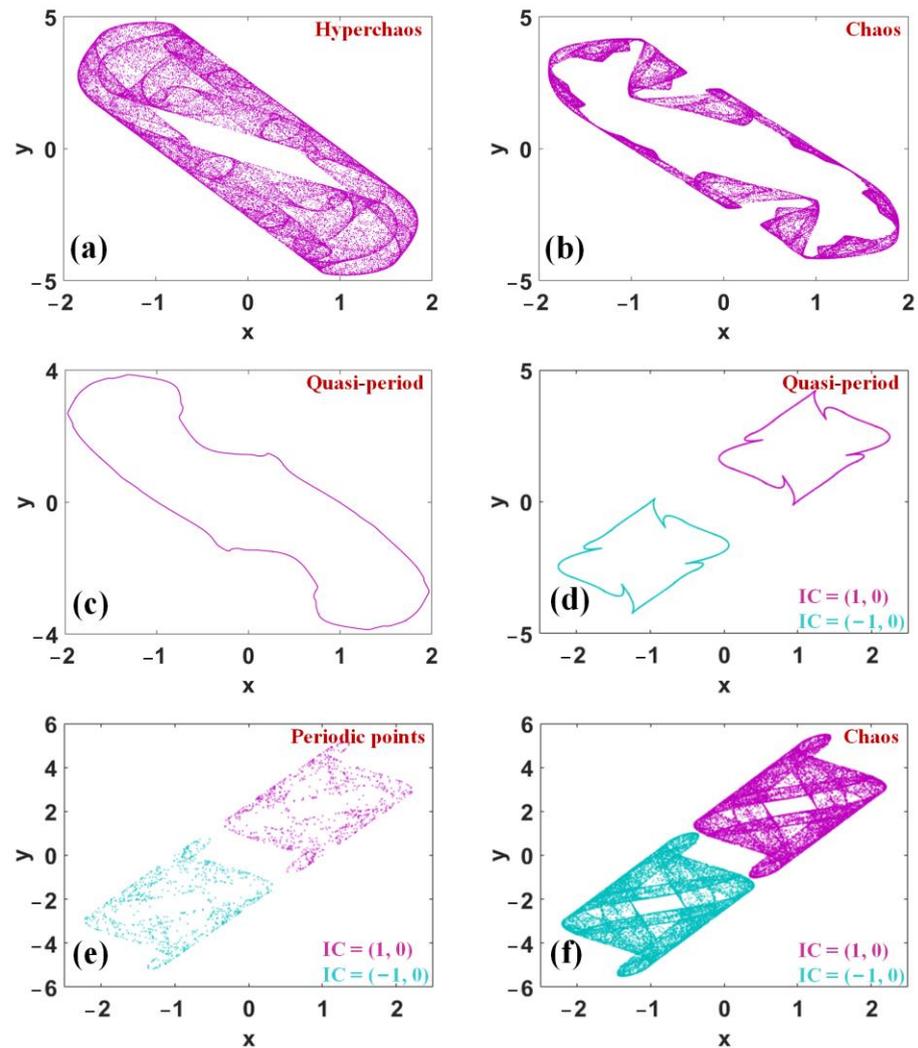


Figure 2. Typical phase trajectories of map (1): (a) $a = -2.6$, (b) $a = -2.2$, (c) $a = -1.96$, (d) $a = 1.887$, (e) $a = 2.336$, (f) $a = 2.5$, where (a–c) under the initial condition $(x_0, y_0) = (1, 0)$.

Table 1. Typical dynamics of map (1) when $(x_0, y_0) = (1, 0)$.

a	Attractor Type	Lyapunov Exponents
-2.6	Hyperchaos	0.2256, 0.09648
-2.2	Chaos	0.1034, -0.0136
-1.96	Quasi-period	0, -0.2493
1.877	Quasi-period	0, -0.3204
2.336	Periodic points	-0.1134, -0.2626
2.5	Chaos	0.3369, -0.1136

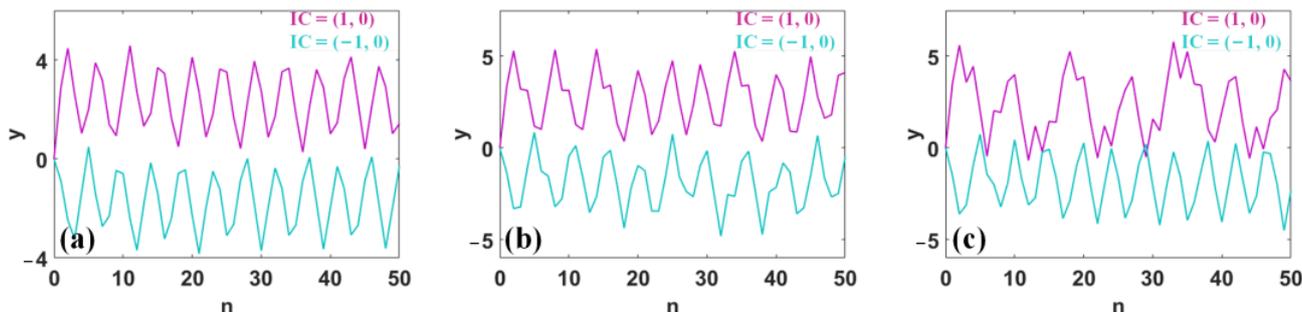


Figure 3. Symmetrical waveforms of map (1): (a) $a = 1.887$, (b) $a = 2.336$, (c) $a = 2.5$.

3. Multiple Alternatives of Offset Boosting

3.1. Initially Controlled Offset Boosting

Discrete map (1) is a dynamic system with self-reproduction. According to the iterative relationship, suppose the initial data of x_n receives the offset boosting, and the properties of the attractor are not disturbed by the value of offset boosting. For the periodicity of the sine function, the attractor is shifted in phase space with inconsistent steps, and the coexisting attractors are arranged by the polarity of x_0 and a .

Case I: $x_0 > 0, a < 0$: when $x_0 > 0, a < 0$, attractor self-reproducing is heading in the negative direction of y , as shown in Figure 4a,b.

Case II: $x_0 > 0, a > 0$: when $x_0 > 0, a > 0$, coexisting attractors are arranged in the positive direction of y , as plotted in Figure 4c,d.

Case III: $x_0 < 0, a < 0$: when $x_0 < 0, a < 0$, the direction of attractor self-reproduction is in the negative direction of y , which is the same as that of case I.

Case IV: $x_0 < 0, a > 0$: when $x_0 < 0, a > 0$, self-reproducing attractors are extracted in the phase space in the positive direction of y , as in case II.

The initial condition x_0 controls the position of coexisting attractors in two-dimensional space. As depicted in Figure 5a, the plot of Lyapunov exponents with many periodic windows indicates the nonsmooth switch from one state to another when $a = -2.6$ and x_0 varies in $[0, 15]$. When $a = -2.6$ and x_0 varies on the negative side of the coordinate axis, in pace with the change in x_0 in the positive direction, the dynamic behavior is also accompanied by mutations. Correspondingly, when $a = 2.5$ and x_0 varies in $[0, 8]$, Figure 5b shows that the dynamic evolution is relatively smooth without a change in state. Moreover, when $a = 2.5$ and $x_0 > 0$, a similar smooth dynamic evolution will also appear. The basins of attraction of the coexisting attractors plotted in Figure 6 also prove this phenomenon. When $a = -2.6$, the separate basins are nonsmooth, while when $a = 2.5$, the boundary of the basins is smooth. The rule of attractor reproduction agrees with the evolution of the square plot of bifurcation.

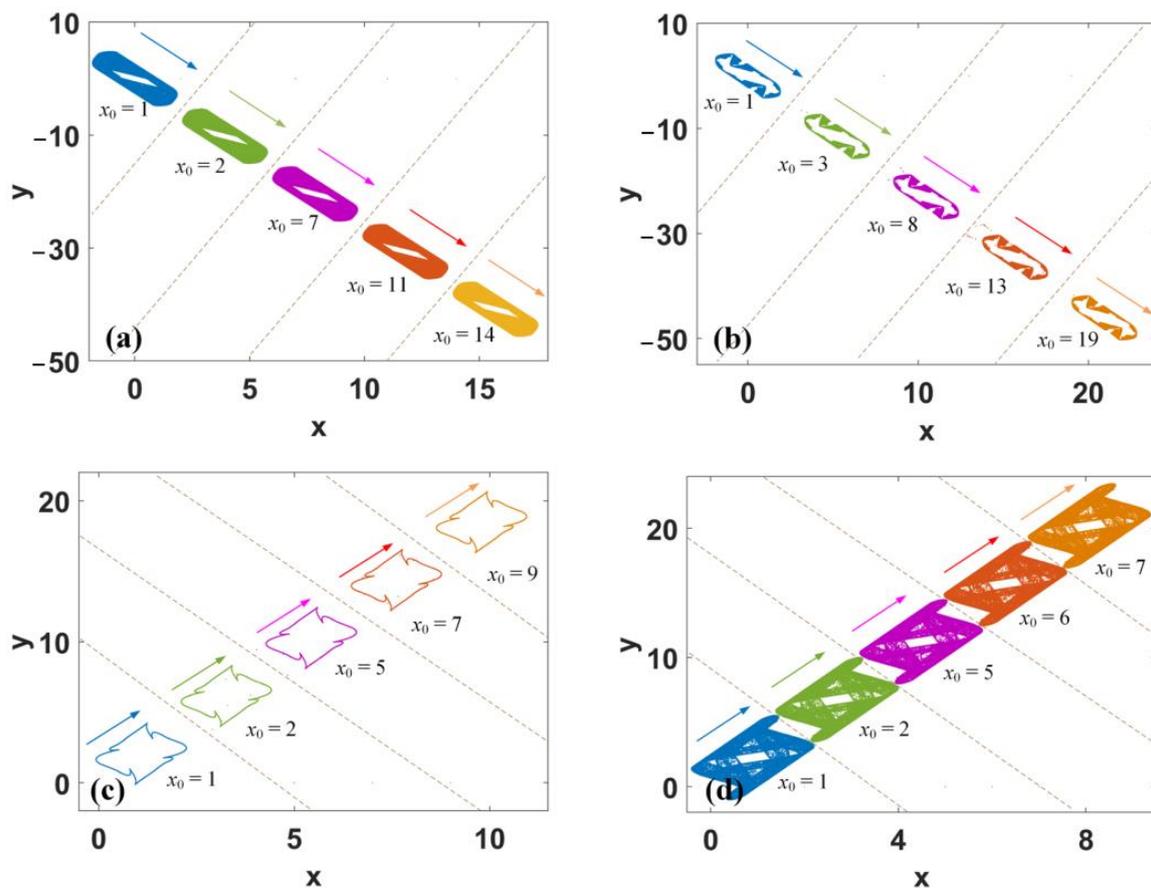


Figure 4. Coexisting phase trajectories of map (1) under various initial conditions: (a) $a = -2.6$, (b) $a = -2.2$, (c) $a = 1.887$, (d) $a = 2.5$.

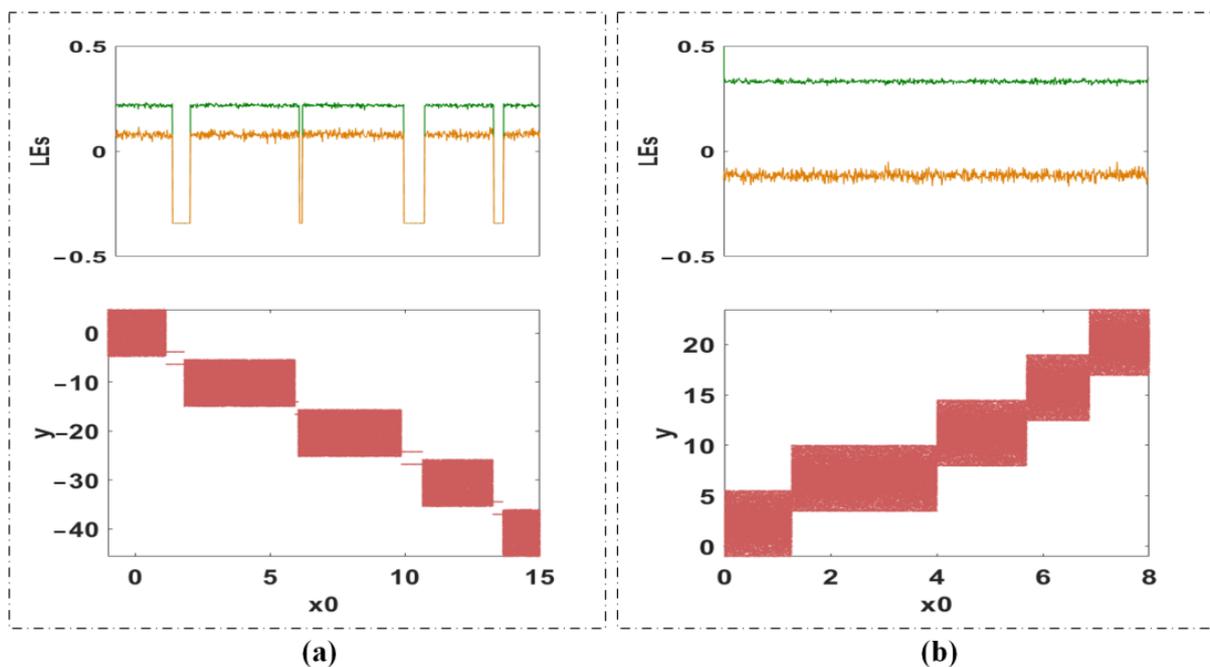


Figure 5. Bifurcation of map (1) under the initial value x_0 : (a) $a = -2.6$, (b) $a = 2.5$.

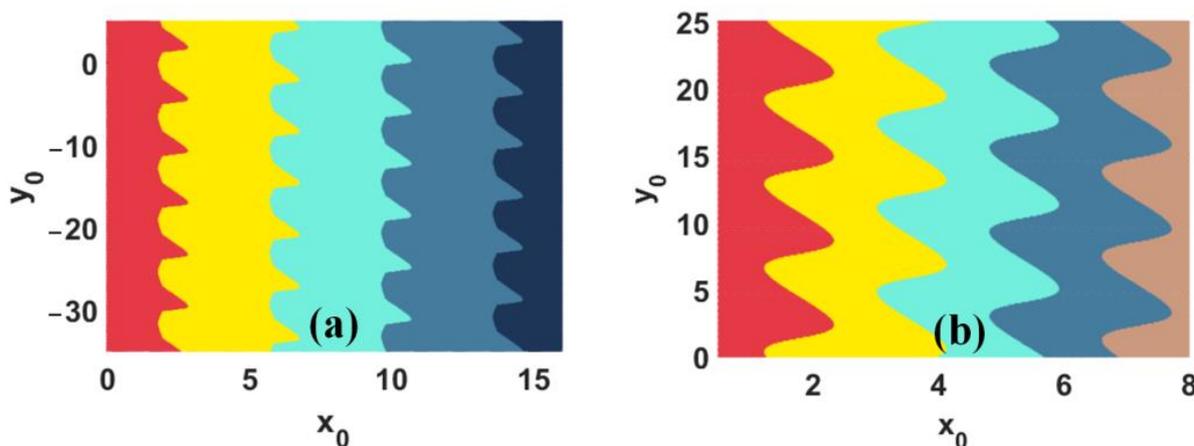


Figure 6. Basins of attraction of the coexisting attractors: (a) $a = -2.6$, (b) $a = 2.5$.

3.2. Parameter-Oriented Offset Boosting

Suppose the sequence x_n has the offset shift l while sequence y_n obtains offset boosting with $-l$, map (1) turns out to be,

$$\begin{cases} x_{n+1} + l = x_n + l + \sin(x_n + l + y_n - l), \\ y_{n+1} - l = a(x_n + l), \end{cases} \tag{3}$$

set $d = al + l = (a + 1)l$, Equation (3) turns out to be,

$$\begin{cases} x_{n+1} = x_n + \sin(x_n + y_n), \\ y_{n+1} = ax_n + d, \end{cases} \tag{4}$$

It means that the constant d in Equation (4) indicates the offset boosting $l = d/(a + 1)$ in the opposite direction along the axis of x and y . Shifted attractors of map (4) are arranged in phase space along a specific direction defined by the parameters a and d .

Case I: $a < -1, d < 0$ or $a > -1, d > 0$: when $a < -1, d < 0$ or $a > -1, d > 0$, then $l > 0$, so the sequence x_n is given a positive offset boosting (PX), and the sequence y_n is offered a negative offset boosting (NY).

Case II: $a > -1, d < 0$ or $a < -1, d > 0$: when $a > -1, d > 0$ or $a < -1, d > 0$, then $l < 0$, and then the sequence x_n is assigned with negative offset boosting (NX), and the sequence y_n is equipped with positive offset boosting (PY).

As shown in Figure 7, the attractor is constantly shifted on the y -axis as well as moved in a certain interval on the x -axis. Four different modes of offset boosting are revealed in Table 2.

Table 2. Various modes of offset boosting in the map (4).

Parameters	$a < -1$	$a > -1$
$d < 0$	$l > 0$	$l > 0$
	PX	NX
	NY	PY
$d > 0$	$l < 0$	$l > 0$
	NX	PX
	PY	NY

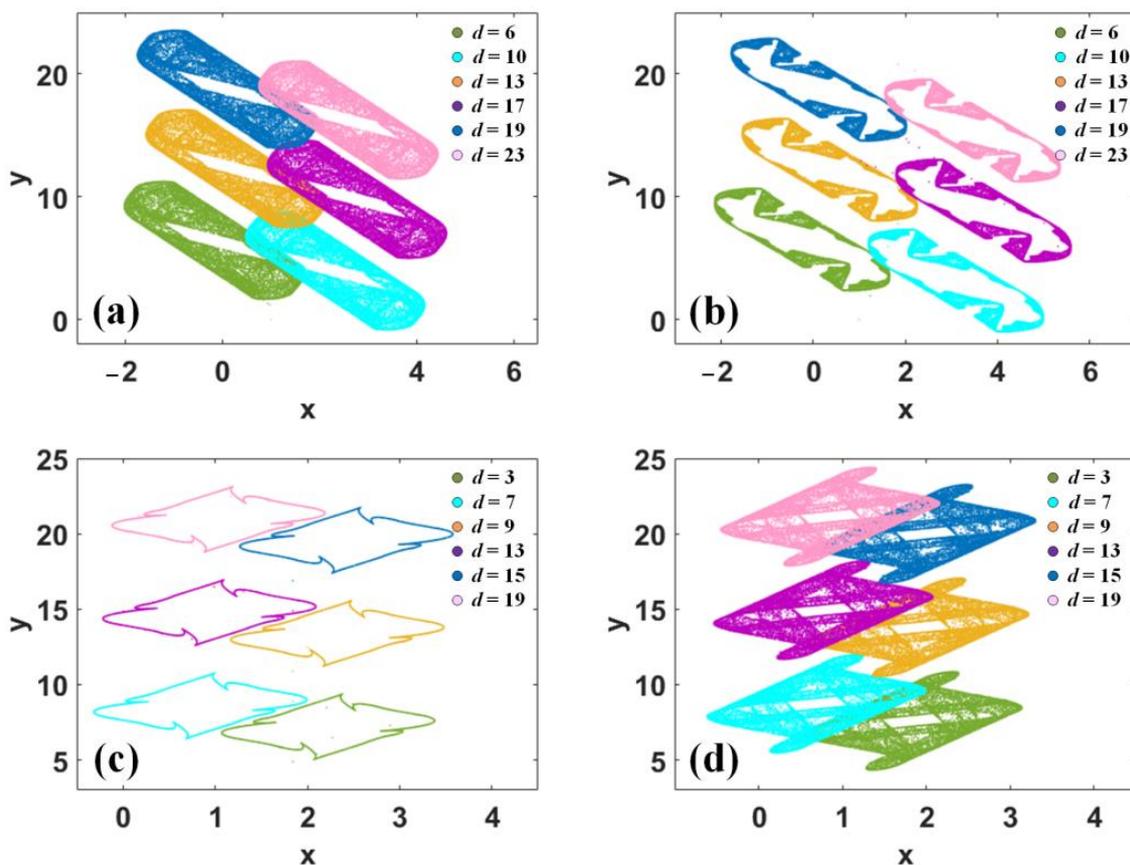


Figure 7. Phase trajectories of map (4) under $(x_0, y_0) = (1, 0)$: (a) $a = -2.6$, (b) $a = -2.2$, (c) $a = 1.887$, (d) $a = 2.5$.

When $a < -1$, $d > 0$, and $l < 0$, the chaotic signal x_n is shifted in the positive direction, and the signal y_n is moved in the negative direction. The dynamical evolution of d is accompanied by many periodic windows, indicating the threat to engineering can be effectively diagnosed by competitive offset boosting. Different from the above-mentioned offset boosting, when $a > -1$, $d > 0$, and $l > 0$, the attractor moves in the negative direction of x and the positive direction of y . The dynamical behavior of d is non-bifurcation without revising Lyapunov exponents, as plotted in Figure 8b. From Figure 8, it is obvious that the evolution of parameter d makes the offset boosting of the signals x_n and y_n different according to the period of 2π . Further explanation will be discussed later.

Suppose the sequence of x_n receives offset boosting with p while the offset of y_n is revised by q , map (1) turns out to be,

$$\begin{cases} x_{n+1} + p = x_n + p + \sin(x_n + p + y_n + q), \\ y_{n+1} + q = a(x_n + p), \end{cases} \tag{5}$$

When $q = ap$, and let $p + q = h$, Equation (5) turns out to be,

$$\begin{cases} x_{n+1} = x_n + \sin(x_n + y_n + h), \\ y_{n+1} = ax_n, \end{cases} \tag{6}$$

It indicates that the constant h in Equation (6) implies the offset boosting $p = h/(a + 1)$ in various combinations of movement in the axis of x and y when a and h are located in different regions, respectively. When $a < -1$ and $h > 0$, then $p < 0$ and $q > 0$, the sequence x_n is given a negative offset boosting, and the sequence y_n is offered a positive offset boosting. Conversely, when $a > 0$ and $h > 0$, then $p > 0$ and $q > 0$, the sequence x_n and y_n are assigned with positive offset boosting. Six different modes of competitive offset boosting are revealed

in Table 3. As shown in Figure 9, the attractors are arranged in a certain interval on the x -axis and y -axis.

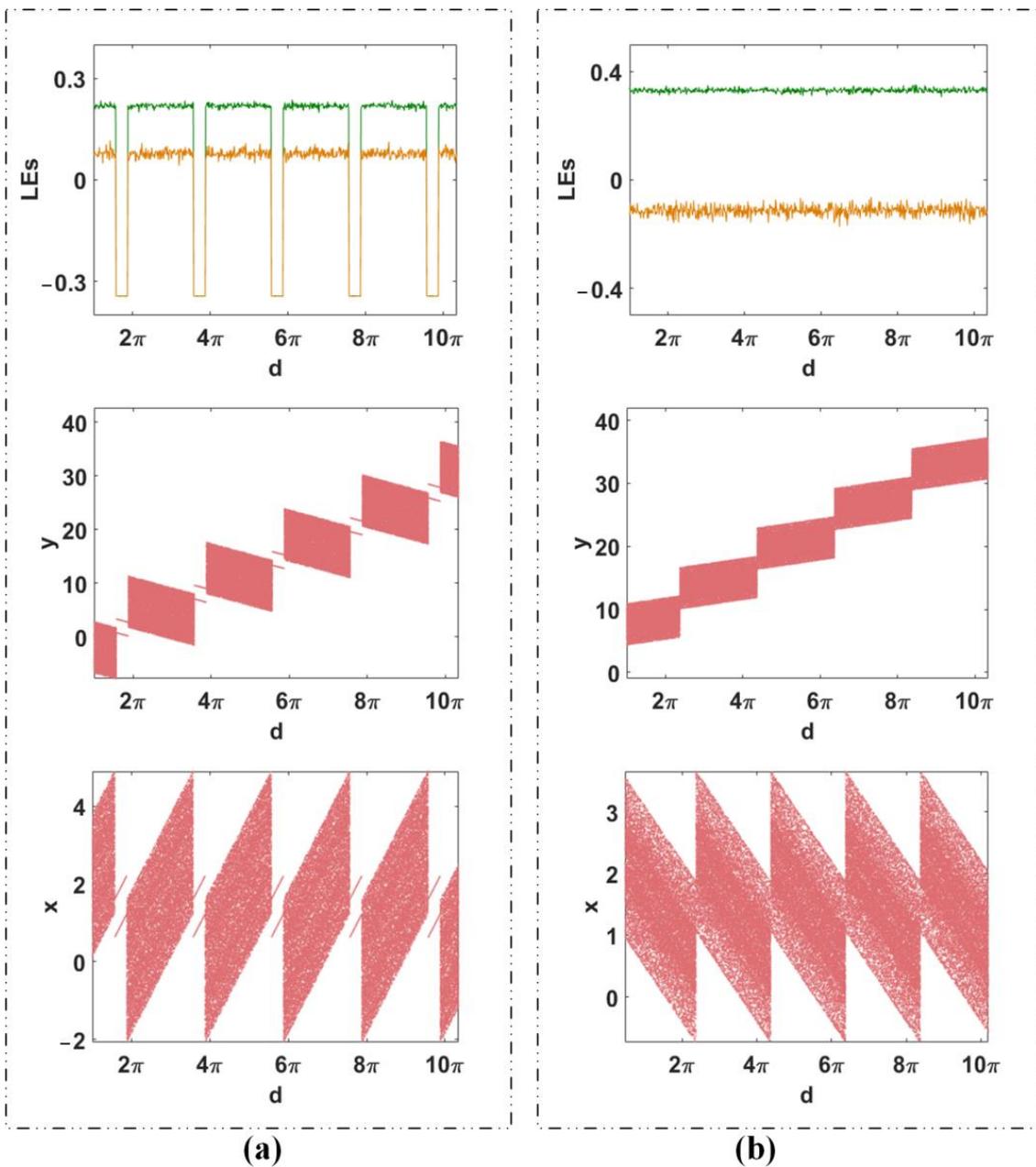


Figure 8. Bifurcation of map (4) under $(x_0, y_0) = (1, 0)$: (a) $a = -2.6$, (b) $a = 2.5$.

Table 3. Regulated directions of offset boosting in the map (6).

Parameters	$a < -1$	$-1 < a < 0$	$a > 0$
$h < 0$	$p > 0, q < 0$	$p < 0, q > 0$	$p < 0, q < 0$
	PX	NX	NX
	NY	PY	NY
$h > 0$	$p < 0, q > 0$	$p > 0, q < 0$	$p > 0, q > 0$
	NX	PX	PX
	PY	NY	PY

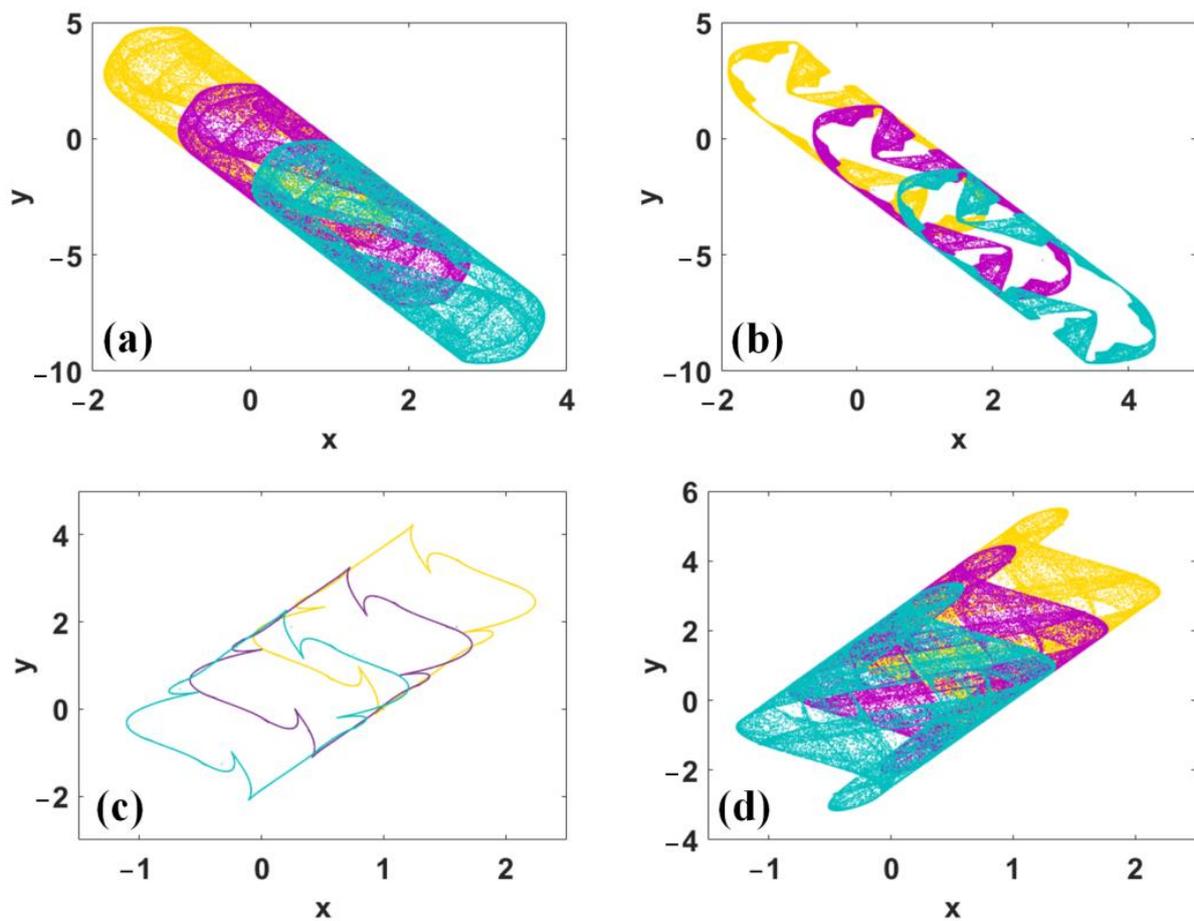


Figure 9. Phase trajectories of map (6) with a different offset constant h under $(x_0, y_0) = (1, 0)$, where gold: $h = 0$, violet red: $h = 1.5$, turquoise: $h = 3$: (a) $a = -2.6$, (b) $a = -2.2$, (c) $a = 1.887$, (d) $a = 2.5$.

The starting position of the chaotic signal can be regulated by the offset constant h , as shown in Figure 10a. When $a < -1$, $h > 0$, $p < 0$, and $q > 0$, the dynamical evolution of h is accompanied by many periodic windows, and signal x_n and y_n have migrations in the opposite direction, as shown in Figure 10b–d.

Because the periodic function appears in the iterative equation of map (1) in the first dimension, offset boosting of map (1) becomes more flexible and shows multiple patterns. According to parameter d in Equation (4), the offset boosting of the signal x_n oscillates periodically with the interval of 2π , but that of the signal y_n climbs periodically, as shown in Figure 8. According to parameter h in Equation (6), the offset boosting with signals x_n and y_n oscillates periodically with a period of 2π , as shown in Figure 10. The multiple alternatives of offset boosting controlled by parameters are summarized in Table 4. Average values of chaotic sequences x_n and y_n are shown in Figure 11.

Table 4. Multiple parameter-oriented patterns of offset boosting.

	x	y
d	Periodic Offset Boosting	Oscillatory Offset Boosting
h	Periodic Offset Boosting	Periodic Offset Boosting

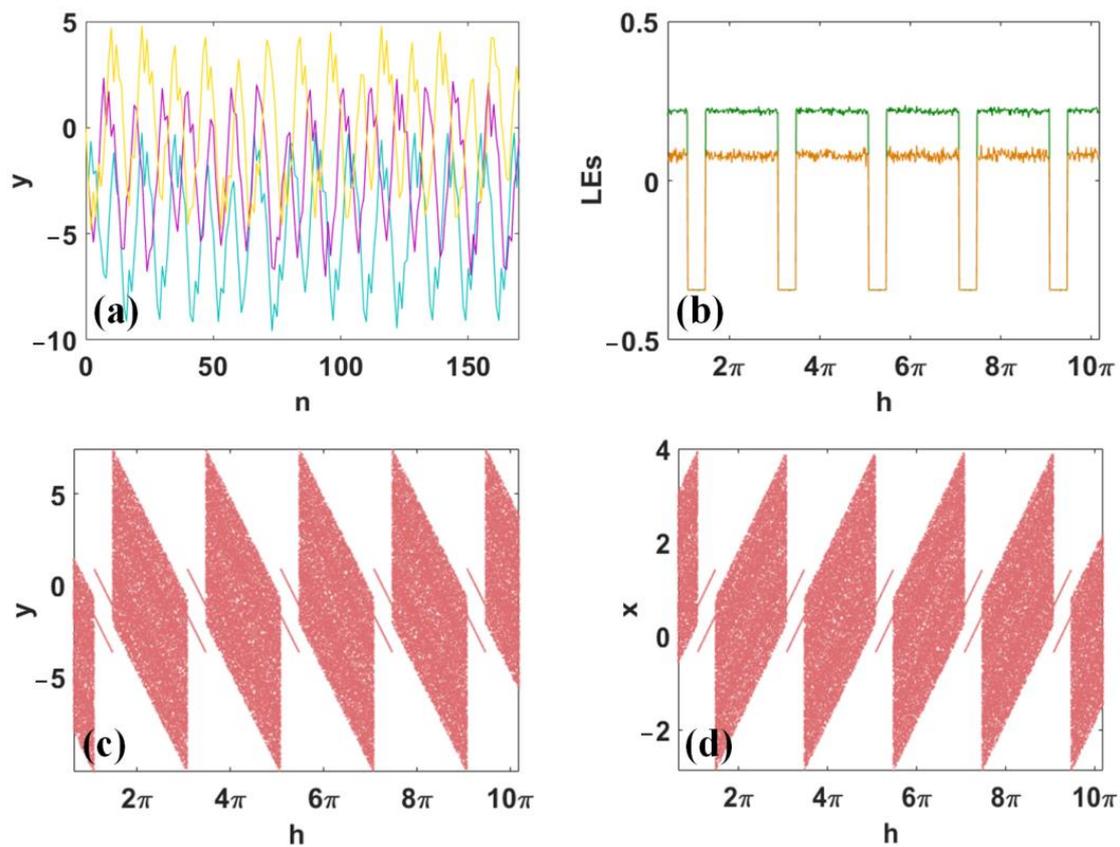


Figure 10. Feature of competitive offset boosting of map (6) with $a = -2.6$ under $(x_0, y_0) = (1, 0)$: (a) discrete sequences, (b) Lyapunov exponents, (c,d) bifurcation diagram.

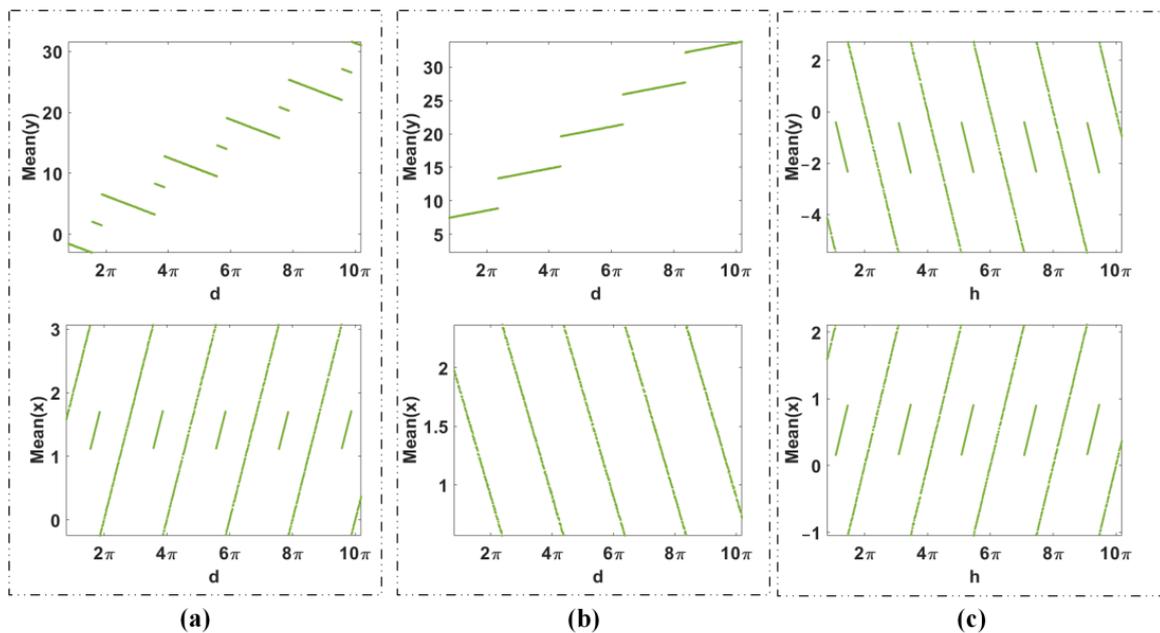


Figure 11. Average values of x_n and y_n : (a) $a = -2.6$ and d varies in $[0, 10\pi]$, (b) $a = 2.5$ and d varies in $[0, 10\pi]$, (c) $a = -2.6$ and h varies in $[0, 10\pi]$.

As shown in Figures 8a and 11a, according to parameter d , the offset of signal y_n climbs globally but falls locally with a period of 2π , the offset of signal x_n oscillates periodically with a period of 2π , as plotted in Figures 8b and 11b. That is to say, there are two different

offset evolutions in a large scope of initial conditions. Specifically, when $a < -1$, the offset of x_n increases in the period but totally remains at the same level according to parameter d in Equation (4). The offset of y_n decreases in the period but grows globally, as shown in Figures 8a and 11a. The evolution of offset boosting according to h in Equation (6) seems to be similar. Here, the offset of y_n decreases periodically while the offset of x_n increases periodically, as shown in Figures 10c,d and 11c.

3.3. Competitive Offset Boosting

Suppose map (1) is given constants d and h at the same time, map (1) turns out to be,

$$\begin{cases} x_{n+1} = x_n + \sin(x_n + y_n + h), \\ y_{n+1} = ax_n + d, \end{cases} \tag{7}$$

As proved above, the offset boosting in the dimensions x and y can be defined by the newly introduced constants d and h along with the initial conditions. This can be called offset competition. In the map (7), due to the existence of the periodic function, the parameters d and h and the initial value of x_0 influence the offset of coexisting attractors, as shown in Figure 12, which means that to obtain a desired offset, the parameter and the initial condition should match each other. From Figure 12, we also know that in map (7) the power of offset boosting of parameter d seems stronger than parameter h , which is because parameter h is in the sinusoidal function. More demonstrations based on Lyapunov exponents can verify this phenomenon further. As shown in Figure 13, when $a < -1$ and x_0 varies with constants d and h , respectively, in different ranges, the largest 2D Lyapunov exponent under the mixture control of a parameter and initial condition is given, showing the strength of offset boosting under the evolutions of a parameter or an initial condition.

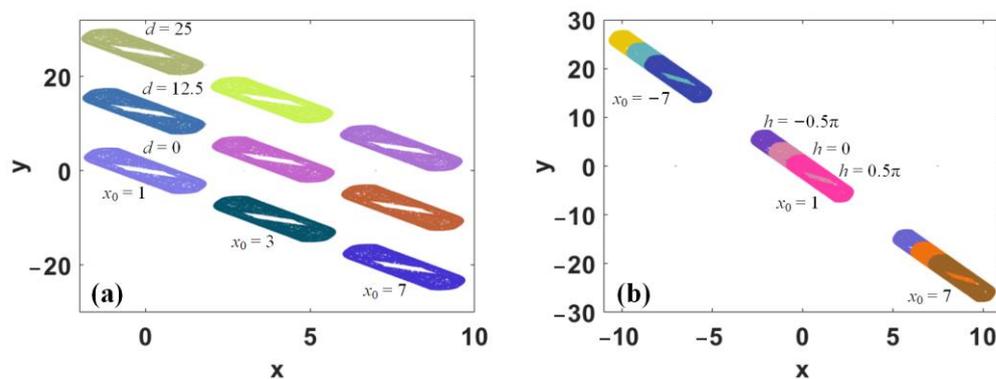


Figure 12. Competitive offset boosting of map (7) when $a = -2.6$ and $y_0 = 0$: (a) $d = 0, 12.5$ and $25, x_0 = 1, 3$ and 7 , (b) $h = -0.5\pi, 0$ and $0.5\pi, x_0 = -7, 1$ and 7 .

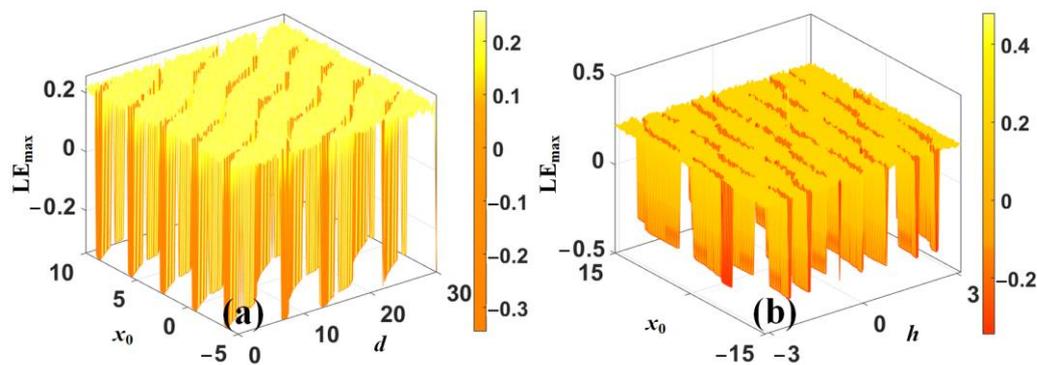


Figure 13. The largest Lyapunov exponent of map (7) when $a = -2.6$ and $y_0 = 0$: (a) x_0 varies in $[-5, 10]$ and d varies in $[0, 30]$, (b) x_0 varies in $[-15, 15]$ and h varies in $[-\pi, \pi]$.

4. Hardware Implementation and Application

4.1. Hardware Implementation Based on STM32

In this section, the implementation of chaotic signal acquisition and transmission using the STM32F103 microcontroller as the core and the TLV5618 module for digital-to-analog conversion are obtained. The resulting waveforms are displayed on an oscilloscope of SDS1102X [27,28]. The TLV5618 is a 12-bit voltage output digital-to-analog converter. The hardware implementation for map (1) is presented in Figure 14. The waveform of map (1) under the initial condition of $(x_0, y_0) = (1, 0)$ is shown in Figure 15. Coexisting phase orbits of map (1) observed from the oscilloscope under various initial conditions are presented in Figure 16. Phase trajectories of map (1) observed from the oscilloscope with competitive offset boosting under $(x_0, y_0) = (1, 0)$ are illustrated in Figure 17. Phase orbits of map (1) with different offset constant h under $(x_0, y_0) = (1, 0)$ are shown in Figure 18. All figures from the hardware implementation based on the STM32 microcontroller and TLV5618 module are consistent with the numerical simulation.

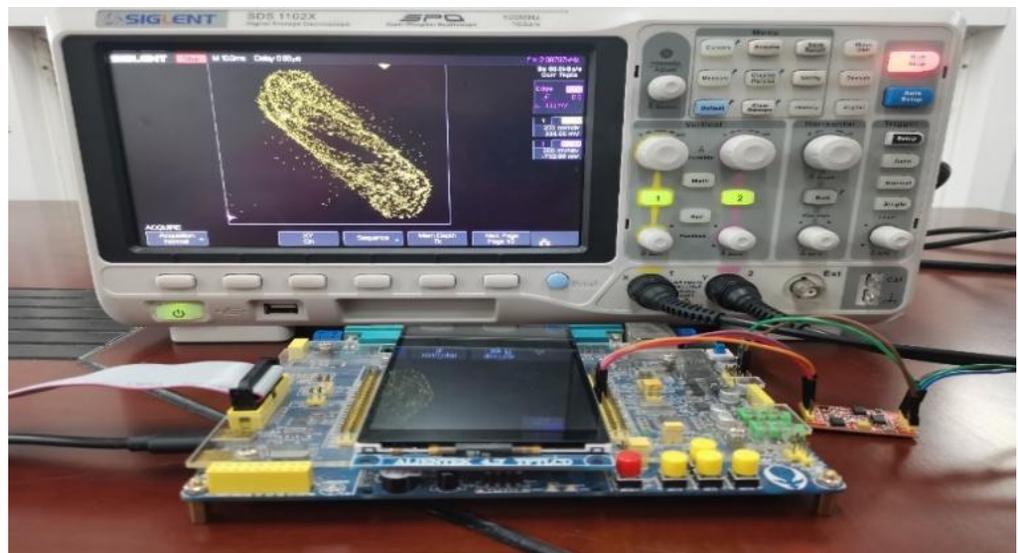


Figure 14. The hardware implementation based on STM32 for map (1).



Figure 15. The hardware-implemented waveform of in map (1) under $(x_0, y_0) = (1, 0)$: (a) hyperchaos: $a = -2.6$, (b) chaos: $a = 2.5$.



Figure 16. Coexisting phase trajectories of map (1) observed from the oscilloscope under different initial conditions: (a) $a = -2.6$, (b) $a = -2.2$, (c) $a = 1.887$, (d) $a = 2.5$.

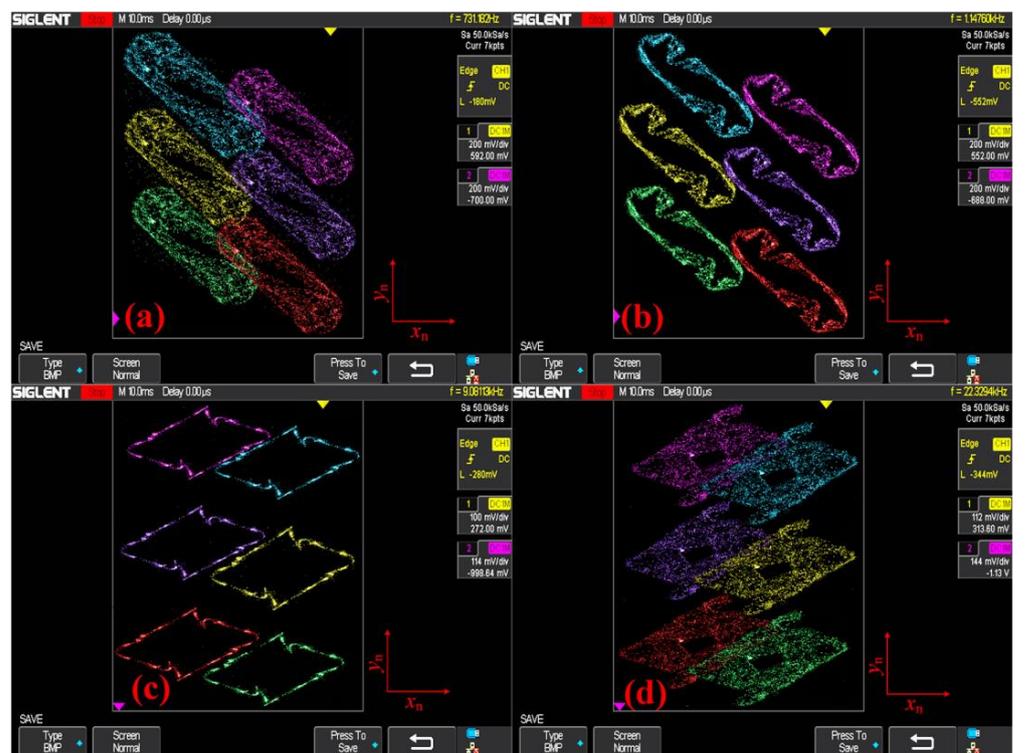


Figure 17. Phase trajectories of map (4) from the oscilloscope under $(x_0, y_0) = (1, 0)$: (a) $a = -2.6$, (b) $a = -2.2$, (c) $a = 1.887$, (d) $a = 2.5$.

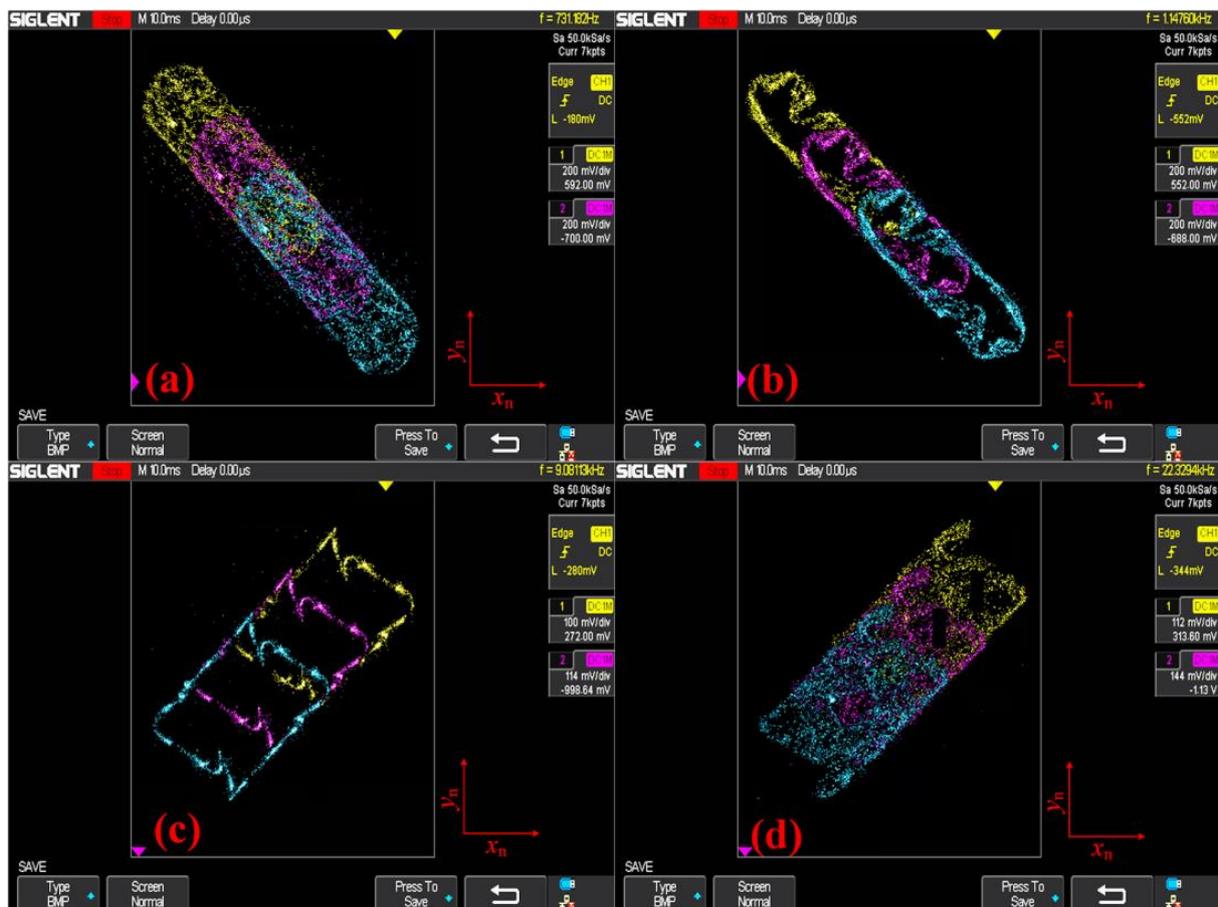


Figure 18. Phase trajectories of map (6) with different offset constant h under $(x_0, y_0) = (1, 0)$, where gold: $h = 0$, violet red: $h = 1.5$, turquoise: $h = 3$: (a) $a = -2.6$, (b) $a = -2.2$, (c) $a = 1.887$, (d) $a = 2.5$.

4.2. Application in PRNG

Assuming that the chaotic sequence generated by map (1) is denoted as $X = \{x_1, x_2, \dots, x_n\}$ or $Y = \{y_1, y_2, \dots, y_n\}$ [29,30]. The quantization function P_i employed in this experiment can be expressed as follows

$$P_i = \lfloor (X_i + |X_{\min}| \cdot K) \rfloor \bmod N, \quad (8)$$

let $K = 10^7$ and $N = 256$ in this paper.

In this experiment, we utilized the hyperchaotic sequence generated by map (1) with parameters $a = -2.6$ and initial condition $(x_0, y_0) = (1, 0)$, the National Institute of Standards and Technology (NIST) test suite is used to measure its performance. To ensure high precision in detection, each sequence set should consist of no less than 1×10^6 bits. By setting the number of test groups as $m = 128$, the proportion (Prop) of passing test groups can indicate the pseudo-randomness of the sequence, thereby demonstrating the potential value for engineering applications of map (1). The results of the NIST statistical test of PRNG are illustrated in Table 5.

Table 5. NIST statistical test of the proposed PRNG.

No.	Statistical Test Terms	PRNG Generated by x		PRNG Generated by y	
		Prop	p -Value	Prop	p -Value
01	Frequency	0.992	0.534146	0.984	0.689019
02	Block frequency	0.984	0.804337	0.976	0.654467
03	Cumulative sums	0.992	0.585209	0.984	0.204076
04	Runs	0.992	0.392456	1	0.178278
05	Longest run	1	0.057146	1	0.437274
06	Rank	0.992	0.723129	0.992	0.242986
07	FFT	0.984	0.253551	0.992	0.090936
08	Non-overlapping template	1	0.991468	1	0.980885
09	Overlapping template	1	0.134686	0.984	0.551026
10	Universal	0.984	0.170294	1	0.324108
11	Approximate entropy	0.968	0.204076	0.992	0.848588
12	Random excursions	1	0.162606	1	0.602458
13	Random excursions variant	1	0.275709	1	0.213309
14	Serial	0.992	0.739918	0.992	0.500934
15	Linear complexity	0.992	0.452799	1	0.324108

5. Discussion and Conclusions

In this paper, a new symmetric hyperchaotic map with multiple regimes of offset boosting is proposed and exhaustively explored. Multiple alternatives of offset boosting are proposed by two parameters and initial conditions. Offset competition is disclosed, where offset is determined by two parameters and initial conditions. In this case, infinitely many coexisting phase orbits are born, which are arranged by the coefficient-matched initial conditions. Therefore, the location of the coexisting hyperchaotic attractor is flexibly selected by the offset parameter and a matched initial condition. This elegant hyperchaotic map with multiple alternatives of offset boosting definitely shows greater potential in other relevant research fields and chaos-based applications, including Brownian movement, neural networks, and image encryption [31–35], which could be our future work.

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